Zero-dead-time Differential Spectroscopy Beyond the Laser Coherence Limit

WILL MCGREW NIST-BOULDER



Image credit: Nate Phillips

Outline

I - Clocks and tests of fundamental physics

II - Differential spectroscopy

III - Zero-dead-time differential spectroscopy

Part I – Clocks and tests of fundamental physics

Clocks

		C 2008 Encyclopæda Britanica, Inc.	Cs 55 J32,00 Cesium	
Ticking rate (Hz)	10 ⁰	107	10 ¹⁰	10 ¹⁵
Instability (s/day)	10 ⁰	10 ⁻⁷	10^{-10}	10 ⁻¹³

Do we need better?



Tests of ultralight dark matter



Arvanitaki, A., Huang, J., & Van Tilburg, K. (2015). Searching for dilaton dark matter with atomic clocks. *Physical Review D*, *91*(1), 015015. <u>https://doi.org/10.1103/PhysRevD.91.015015</u>

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Tests of ultralight dark matter



Highly local

Extraordinary precision (better than 10⁻¹⁶ in fractional frequency)

Bandwidth of ~Hz and below (~10⁻¹⁵ eV and below)

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Constraints on ultralight dark matter



BACON Collaboration. (2021). Frequency ratio measurements at 18-digit accuracy using an optical clock network. *Nature, 591,* 564–569. <u>https://doi.org/10.1038/s41586-021-03253-4</u>

Clock instability

Dick effect: technical biasing phenomenon related to the OLO noise properties

Other sources of noise:

$$\sigma_y(\tau) = \left(\frac{s}{N}\right)^{-1} n_{meas}^{-1/2}$$

Contributions: technical detection noise, photon shot noise, quantum projection noise (QPN)

QPN is fundamental (for uncorrelated quantum systems). Other sources are subject to technical reduction.

Quantum projection noise



 $(\phi,\psi)\in \mathfrak{R} imes \mathfrak{R}$

 $p \in \mathfrak{B}$

$$\left(\frac{S}{N}\right)^{-1} = \frac{\Delta f}{f_0} = \frac{\left(\frac{df}{dp}\right)\Delta p}{f_0}$$
, but due to projection noise, $\Delta p \approx \frac{1}{2}$. For Ramsey spectroscopy, $\frac{df}{dp} \approx \frac{1}{\pi T_i}$.
 $\left(\frac{S}{N}\right)^{-1} = \frac{1}{2\pi T_i f_0}$ for a single quantum oscillator. For N oscillators, $\left(\frac{S}{N}\right)^{-1} = \frac{1}{2\pi T_i f_0 \sqrt{N}}$.

Quantum projection noise

$$\sigma_y(\tau) = \left(\frac{S}{N}\right)^{-1} n_{meas}^{-1/2}$$

The number of measurements we have is $n_{meas} = \frac{\tau}{T_c}$, where τ is the measurement time and T_c is the cycle time.

$$\sigma_y(\tau) = \frac{1}{2\pi T_i f_0} \sqrt{\frac{T_c}{N\tau}} \approx \frac{1}{2\pi f_0 \sqrt{NT_i \tau}}$$

$$\sigma_y(\tau) \approx \frac{1}{2\pi f_0 \sqrt{NT_i \tau}}$$

Increase **7**

Increase f_0

Increase N

Increase *T_i*







$$\sigma_y(\tau) \approx \frac{1}{2\pi f_0 \sqrt{NT_i \tau}}$$

Increase au

Increase f_0

Increase N

Increase *T_i*

Can we increase N?

Two major types of optical clocks: ion clocks and neutral atom clocks

- $\,^{\rm o}\,$ Lattice clocks can operate with $N \approx 10^4$
- Coupling between ions have so far constrained ion clocks to N = 1

However, exemplary systematic uncertainties have been attained by ion clocks

Also, some ion clocks have high sensitivity for tests of fundamental physics

TABLE I. Fractional frequency shifts $(\Delta \nu / \nu)$ and associated systematic uncertainties for the ²⁷Al⁺ quantum-logic clock.

Effect	Shift (10^{-19})	Uncertainty (10 ⁻¹⁹)
Excess micromotion	-45.8	5.9
Blackbody radiation	-30.5	4.2
Quadratic Zeeman	-9241.8	3.7
Secular motion	-17.3	2.9
Background gas collisions	-0.6	2.4
First-order Doppler	0	2.2
Clock laser Stark	0	2.0
AOM phase chirp	0	<1
Electric quadrupole	0	<1
Total	-9336.0	9.4

Brewer, S. M., Chen, J. S., Hankin, A. M., Clements, E. R., Chou, C. W., Wineland, D. J., ... Leibrandt, D. R. (2019). Al+ 27 Quantum-Logic Clock with a Systematic Uncertainty below 10-18. *Physical Review Letters*, *123*(3), 033201.

Our previous comparisons

Sr lattice clock at JILA

Al⁺ ion clock and Yb lattice clock at NIST

20 days of continuous comparison to reach 1×10^{-18} instability for Al⁺/Yb or Al⁺/Sr



BACON Collaboration. (2021). Frequency ratio measurements at 18-digit accuracy using an optical clock network. *Nature*, *591*, 564–569.

Interspecies comparisons

Comparisons between different optical clocks



$$\sigma_y(\tau) \approx \frac{1}{2\pi f_0 \sqrt{NT_i \tau}}$$

Increase **7**

Increase f_0

Increase N

Increase T_i

The coherence limit to interrogation



Increase the stability of the OLO?



Matei, D. G., Legero, T., Häfner, S., Grebing, C., Weyrich, R., Zhang, W., ... Sterr, U. (2017). 1.5 μm Lasers with Sub-10 mHz Linewidth. *Physical Review Letters*, *118*(26), 263202.

Part II – Differential spectroscopy

Novel interrogation schemes

Use one or several systems to prestabilize another

- Borregaard, J., & Sørensen, A. S. (2013). Efficient atomic clocks operated with several atomic ensembles. *Physical Review Letters*, *111*(9), 090802.
- Schioppo, M., Brown, R. C., McGrew, W. F., Hinkley, N., Fasano, R. J., Beloy, K., ... Ludlow, A. D. (2017). Ultrastable optical clock with two cold-atom ensembles. *Nature Photonics*, *11*, 48–52.

Dynamical decoupling

• Dörscher, S., Al-Masoudi, A., Bober, M., Schwarz, R., Hobson, R., Sterr, U., & Lisdat, C. (2020). Dynamical decoupling of laser phase noise in compound atomic clocks. *Communications Physics*, *3*, 185.

Novel interrogation schemes

Nondestructive measurements to achieve a phase lock

- Kohlhaas, R., Bertoldi, A., Cantin, E., Aspect, A., Landragin, A., & Bouyer, P. (2015). Phase locking a clock oscillator to a coherent atomic ensemble. *Physical Review X*, 5(2), 021011.
- Bowden, W., Vianello, A., Hill, I. R., Schioppo, M., & Hobson, R. (2020). Improving the Q Factor of an Optical Atomic Clock Using Quantum Nondemolition Measurement. *Physical Review X*, *10*(4), 041052.

Correlation spectroscopy

- Chwalla, M., Kim, K., Monz, T., Schindler, P., Riebe, M., Roos, C. F., & Blatt, R. (2007). Precision spectroscopy with two correlated atoms. *Applied Physics B: Lasers and Optics, 89*, 483–488.
- Clements, E. R., Kim, M. E., Cui, K., Hankin, A. M., Brewer, S. M., Valencia, J., ... Hume, D. B. (2020). Lifetime-Limited Interrogation of Two Independent Al+ 27 Clocks Using Correlation Spectroscopy. *Physical Review Letters*, 125(6), 243602.

Differential spectroscopy – a proposal

Phase lock OLOs to each other

Interrogate synchronously

Phase corrections from Yb



Hume, D. B., & Leibrandt, D. R. (2016). Probing beyond the laser coherence time in optical clock comparisons. *Physical Review A*, *93*(3), 032138.

Differential spectroscopy – implementation

Provide phase corrections from a high-(S/n), lower-frequency clock

The high frequency clock can operate outside of its inversion regime

Increase interrogation time by the frequency ratio (2.16 for Al⁺/Yb)



Differential spectroscopy – results with 500 ms spectroscopy



Part III – ZDT differential spectroscopy

Zero-dead-time spectroscopy

Originally demonstrated to suppress the Dick effect

Can be naturally paired with differential spectroscopy



Schioppo, M., Brown, R. C., McGrew, W. F., Hinkley, N., Fasano, R. J., Beloy, K., ... Ludlow, A. D. (2017). Ultrastable optical clock with two cold-atom ensembles. *Nature Photonics*, *11*, 48–52.

Combining differential spectroscopy with zero-dead-time spectroscopy



Differential spectroscopy results



No phase corrections from Yb

Single-clock differential spectroscopy

ZDT differential spectroscopy

Four cycles, 1.7 s spectroscopy time

Lifetime limit for QPN:

 $6 \times 10^{-17} / \sqrt{\tau}$

Previously, 20 days to reach 1×10^{-18}

Now, 7 hours



Image credit: Nick Nardelli

What continues to limit the instability?

Magnetic field fluctuations on both clocks

Phase-noise-cancellation instability

Frequency transfer uncertainty through the comb

All these are technical and subject to reduction

Interspecies clock comparisons

Comparisons between different optical clocks



Perspectives: improvement to DM constraints

Improved interspecies instability affords improved constraint of DM coupling

ZDT differential spectroscopy can improve any interspecies comparison

 $K_{Yb} = 0.31$ $K_{Al^+} = 0.008$

 $K_{Yb^+E3} = -5.95$ (Ekkehard Peik's talk)

HCIs, Th nuclear clock, etc. offer even better improvements (Marianna Safronova's and Steven King's talks)



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