

# Gravitational Anomalies, Dark Matter and Leptogenesis

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## References:

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Tie together  
Bary(Lepto)ogenesis (BAU),  
dark matter and neutrino  
mass from a string-inspired  
perspective

**Our recipe:**

**D** Dark matter  
( gravitational axion)

**A** Anomalies  
(gravitational + gauge)

**M** ordinary matter  
(SM)

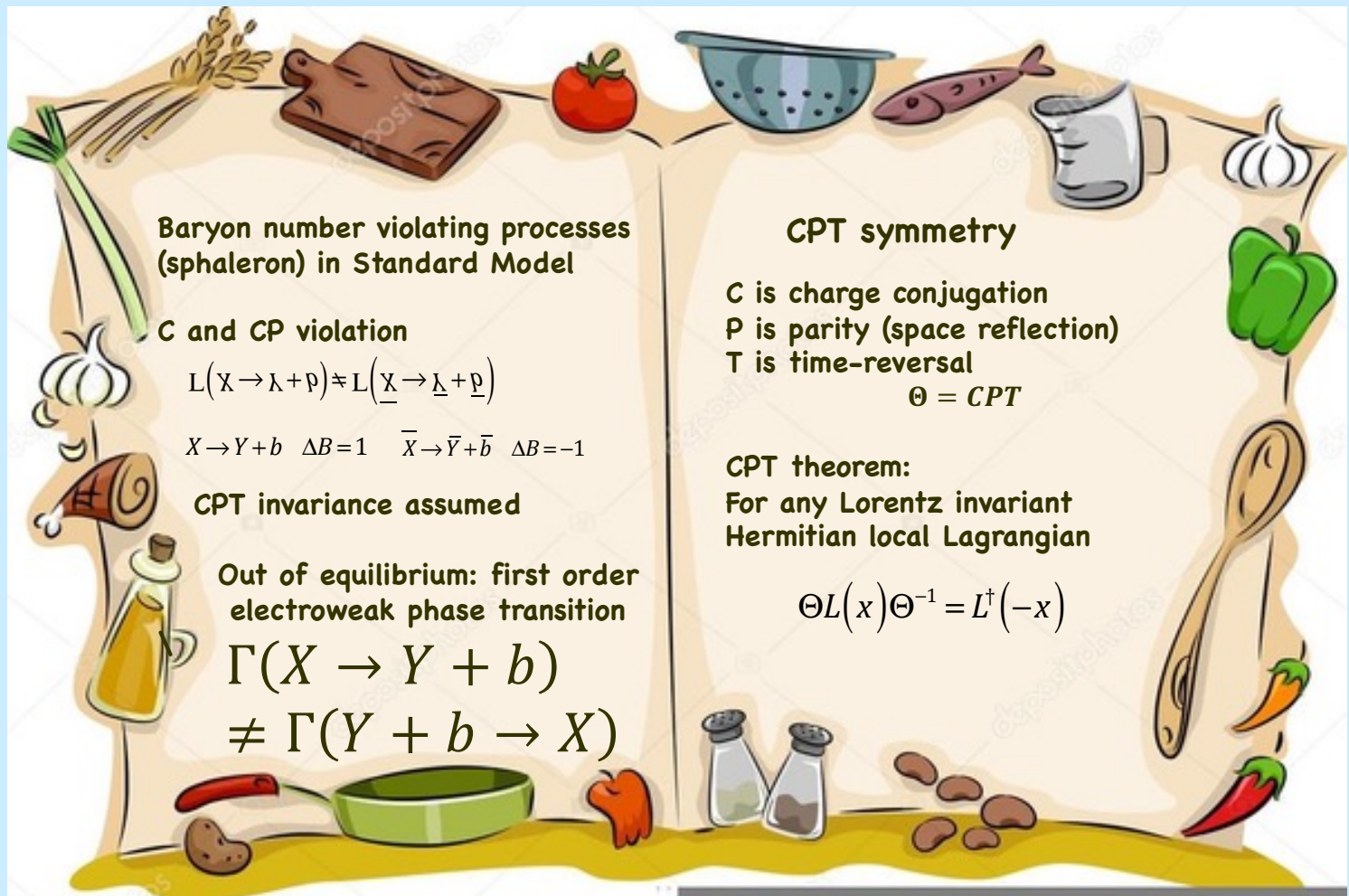
**S** strings  
(gravity with torsion)



# Components of the Universe's energy density:

- $\rho_0$  component energy density
- $\rho_c$  critical energy density
- $\Omega = \frac{\rho_0}{\rho_c}$ ,  $\Omega_{matter} = 0.27 \pm 0.04$ ,  $\Omega_B = 0.044 \pm 0.004$
- $\Omega_B$  is the ratio for baryonic matter
- $n_B$  baryon number density,  $n_\gamma$  photon number density
- $\frac{n_B}{n_\gamma} = 6.1 \pm 0.3 \times 10^{-10}$  at  $T \sim 1 \text{ GeV}$
- If Universe had been matter-antimatter symmetric  $\frac{n_B}{n_\gamma} = \frac{n_{\bar{B}}}{n_\gamma} \approx 10^{-18}$
- **Leptogenesis: creation of primordial lepton-antilepton asymmetry** (M Fukugita and T Yanagida, Phys. Lett. B174 45 (1986))

# Sakharov's Baryogenesis recipe (A D Sakharov, JETP Lett. 5 24 (1967))



Anomalies here there and everywhere:  
( literally the spice of life)

- **Electroweak interactions chiral**


- $J_\mu^B = \frac{1}{3} \sum_f (\bar{q}_{fL} \gamma_\mu q_{fL} + \bar{u}_{fR} \gamma_\mu u_{fR} + \bar{d}_{fR} \gamma_\mu d_{fR})$

$$J_\mu^L = \sum_f (\bar{l}_{fL} \gamma_\mu l_{fL} + \bar{e}_{fR} \gamma_\mu e_{fR})$$

- **Triangle (ABJ) anomaly** ( S L Adler, Phys. Rev 177 2426 (1969))

- $\partial^\mu J_\mu^L = \partial^\mu J_\mu^B = \frac{N_f}{32\pi^2} (-g^2 W_{\mu\nu} \tilde{W}^{\mu\nu} + g'^2 B_{\mu\nu} \tilde{B}^{\mu\nu})$

where the  $W_{\mu\nu}$  and  $B_{\mu\nu}$  are the  $SU(2)$  and  $U(1)$  gauge field two forms

- $B - L$  is conserved  Leptogenesis leads to baryogenesis

- Sakharov recipe difficulties: insufficient CP violation in the Standard Model, electroweak phase transition required to be 1<sup>st</sup> order
- DAMS ingredients beyond SM

Bosonic gravitational multiplet of strings consists of a

- graviton,  $g_{\mu\nu}$
- spin 0 scalar field  $\Phi$  the dilaton
- spin 1 antisymmetric gauge field  $B_{\mu\nu}$  with gauge invariant field strength

$$H_{\mu\nu\rho} = \partial_{[\mu} B_{\nu\rho]}$$

$$e^{-2\Phi} \epsilon_{bac}{}^{\mu} H^{abc} = 4\partial^{\mu} b(x) \xleftarrow{\text{axion}} = 4B_{\mu}(x)$$

- $N$  right-handed sterile neutrino with large Majorana mass
- A CPTV background with non-zero  $B_{\mu}(x)$

# Strings Gravity with torsion

- The Kalb-Ramond field tensor in the low energy string gravitational action appears as:

$$S = \int d^4x \sqrt{-g} \left( \frac{1}{2\kappa^2} R - \frac{1}{6} H_{\mu\nu\rho} H^{\mu\nu\rho} \right) = \int d^4x \sqrt{-g} \left( \frac{1}{2\kappa^2} \bar{R} \right)$$

where  $\bar{R}$  is the curvature for the *affine* connection

$$\bar{\Gamma}_{\nu\rho}^{\mu} = \Gamma_{\nu\rho}^{\mu} + \frac{\kappa}{\sqrt{3}} H_{\nu\rho}^{\mu} \neq \bar{\Gamma}_{\rho\nu}^{\mu}$$

(D J Gross and J H Sloan, Nucl. Phys. B  
291 41 (1987))

contorsion

- Torsion tensor  $T^{\kappa}_{\lambda\mu} = \Gamma^{\kappa}_{\lambda\mu} - \Gamma^{\kappa}_{\mu\lambda}$
- In a generalisation of Einstein's theory: add torsion to get Einstein-Cartan theory
- Torsion is necessary to incorporate fermions

# DAMS Leptogenesis effective Lagrangian

$$\mathcal{L} = \mathcal{L}_{SM} + i \bar{N} \gamma^\mu \partial_\mu N - \frac{m_N}{2} (\bar{N}^c N + \bar{N} N^c) - \bar{N} \gamma^\mu B_\mu \gamma^5 N - \sum_k y_k \bar{L}_k \tilde{\phi} N + h.c.$$

CPTV and LIV background from torsion

- Model leads to acceptable leptogenesis at the tree level, owing to Majorana nature of  $N$  which is taken to be very heavy
- Freeze out of  $N$  at  $T = T_D$ .
- $B$ - $L$  conservation converts  $\Delta L$  to  $\Delta B$
- $B_\mu = \partial_\mu b$
- Non-zero  $B_0$  background
- The *underlying torsion* connects this model to dark matter



## Details of leptogenesis

- Stick to  $k = 1$  for simplicity
- **Background:**

$$B_i = 0, \quad i = 1, 2, 3$$
$$B_0(z) = \Phi f(z)$$

where  $z = \frac{m_N}{T}$ , and  $f(z) = 1$  or  $f(z) = z^{-3}$ .

- **Gravitational background**

$g_{00} = 1, g_{ij} = -a^2(t)\delta_{ij}$ , space Cartesian

- **Generation of lepton asymmetry due to CP and CPTV tree-level decays:**
- $N \rightarrow l^- h^+, \nu h^0$  channel 1
- $N \rightarrow l^+ h^-, \bar{\nu} h^0$  channel 2

where  $l^\pm$  are charged leptons,  $\nu (\bar{\nu})$  are active (anti)neutrinos

- For  $\Phi \neq 0$  the decay rates of  $N$  differ for channels 1 and 2 leading to a lepton asymmetry which freezes out at  $T_D$ .

# Results for leptogenesis

$$\frac{\Delta L^{tot}}{s} \simeq .017 \frac{B_0}{m_N} \text{ at freezeout temperature } T = T_D$$

$$\frac{m_N}{T_D} \simeq 1.6$$

For  $y \sim 10^{-5}$  and  $m_N = O(100)TeV$  and  $B_0 \sim 0.1MeV$  for phenomenologically acceptable leptogenesis to occur at

$$T_D \sim 60TeV$$

# Gravitational axions as Dark Matter I

- Peccei and Quinn solution of strong CP problem

$$\mathcal{L} \supset \vartheta(x) \frac{\alpha_s}{2\pi} \text{Tr}(G \wedge G)$$

where the  $\vartheta$  vacuum parameter becomes a field and  $G$  is the QCD gauge field 2-form (R D Peccei and H R. Quinn, Phys Rev. Lett 38 1440 (1977))

Standard procedure: select a vacuum with  $\langle \vartheta(x) \rangle = 0$  with a fluctuation  $a(x)$

- *Move away from Goldstone axion  $a(x)$  to gravitational axion  $b(x)$  whose couplings are determined by the Planck scale*
- *Classical torsionful spacetime manifold characterised by:*

$$\underline{e}^a = e^a_\mu dx^\mu$$

$$\underline{\omega}^{ab} = \omega^{ab}_\mu dx^\mu$$

$$d\underline{e}^a + \underline{\omega}^a_b \wedge \underline{e}^b = \mathfrak{S}^a,$$

$$d\underline{\omega}^a_c + \underline{\omega}^a_b \wedge \underline{\omega}^b_c = \underline{R}^a_c$$

torsion 2-form

curvature 2-form

where  $\underline{e}^a$  is the vierbein and  $\underline{\omega}^a_b$  is the spin-connection 1-forms

## Gravitational axions as Dark Matter II

The action of fermions of SM on a torsionful manifold is given by

$$S = S_{gravity} - \frac{1}{2} \sum_f \int \left( \bar{\Psi}_f \underline{\gamma} \wedge * \mathcal{D}\Psi_f - \mathcal{D}\bar{\Psi}_f \wedge * \underline{\gamma} \Psi_f \right) - \frac{1}{2} \int \mathcal{F} \wedge * \mathcal{F} - \int Tr[G \wedge * G]$$

where for any fermionic field  $\Psi$

$$\mathcal{D}\Psi = d\Psi + \frac{1}{4} \underline{\omega}^{ab} \gamma_{ab} \Psi + ie \underline{\mathcal{A}} \Psi + ig \underline{\mathcal{B}} \Psi,$$

$$S_{gravity} = \frac{1}{4\kappa^2} \int \varepsilon_{abcd} \underline{\mathcal{R}}^{ab} \wedge \underline{e}^c \wedge \underline{e}^d$$

$$\kappa^2 = 8\pi M_{Pl}^2$$

with  $\underline{\mathcal{A}}$  and  $\underline{\mathcal{B}}$  the gauge connection 1-forms for  $U(1)_{em}$  and  $SU(3)_c$ ,

$$\underline{\gamma} = \gamma_a \underline{e}^a \text{ and } \gamma_{ab} = \frac{1}{2} [\gamma_a, \gamma_b].$$

Equation of motion with respect to  $\underline{\omega}^{ab}$  gives that the torsion satisfies

$$\mathcal{T}_{abc} = -\frac{\kappa^2}{2} \varepsilon_{abcd} \sum_f J_f^{5d}$$

where  $J_f^{5d} = i \bar{\Psi}_f \gamma^d \gamma^5 \Psi_f$ . (See also R Utiyama, Phys Rev 101, 1597 (1956); T W B Kibble, J Math Phys 2 212 (1961); D W Sciama, Monthly Not. Roy. Astr. Soc. 113 34 (1953); F W Hehl, Gen Rel Grav 4 333 (1973); E Cartan, Riemannian Geometry in an Orthogonal Frame (World Scientific, 2001), O Casillo-Feliso; a et al. Phys Rev D 91 085017 (2015) )

# Anomalies

(cf S Basilakos et al. PRD 101 045001 (2020))

M J Duncan, N Kaloper, K A Olive, Nucl. Phys. B387 215 (1992))

Writing  $J_-^5 = \sum_f (J_f^5) \frac{e^a}{a_-}$  as a 1-form the axial anomaly implies

$$d * J_-^5 = -\frac{\alpha_{em} \bar{Q}^2}{\pi} \mathcal{F} \wedge \mathcal{F} - \frac{\alpha_s N_q}{2\pi} \text{Tr}[G \wedge G] - \frac{N_f}{8\pi^2} \tilde{R}^{ab} \wedge \tilde{R}^{ab}$$

which leads to the following *quantised* torsion-field contribution to the action

$$-\frac{\alpha_{em} \bar{Q}^2}{\pi f_b} \int b \mathcal{F} \wedge \mathcal{F} - \frac{1}{2} \int db \wedge * db - \frac{1}{8\pi^2} \int \left( \Theta + \frac{N_f}{f_b} b \right) \tilde{R}^{ab} \wedge \tilde{R}_{ab} - \frac{\alpha_s}{2\pi} \int \left( \theta + \frac{N_q}{f_b} b \right) \text{Tr}[G \wedge G]$$

The parameter  $f_b = \kappa^{-1} \sqrt{8/3} = 4 \times 10^{18}$  GeV.

$\Theta$  and  $\theta$  terms label Pontryagin densities, are allowed in the theory, play the role of counterterms and do not affect the equations of motion.

From the last term see that  $b$  solves the strong CP problem and  $m_b \sim 10^{-12} eV$  and so can be regarded as the QCD axion.

## Concluding remarks

- The gravitational anomaly (GA) may not be cancelled in the early universe in a quantum theory of gravity.
- The gauge anomaly though can be designed to cancel for a suitable GUT theory. Anomaly cancellation is necessary for the theories to be renormalizable.
- GA implies breakdown of diffeomorphism Invariance and hence local Lorentz invariance (LIV)
- LIV allows a non-zero  $\Phi$  required for our leptogenesis model
- The Kalb-Ramond axion could play the role of a QCD axion and so be axionic dark matter (M Lattanzi and S Mercuri, Phys Rev D 81 125015 (2010))