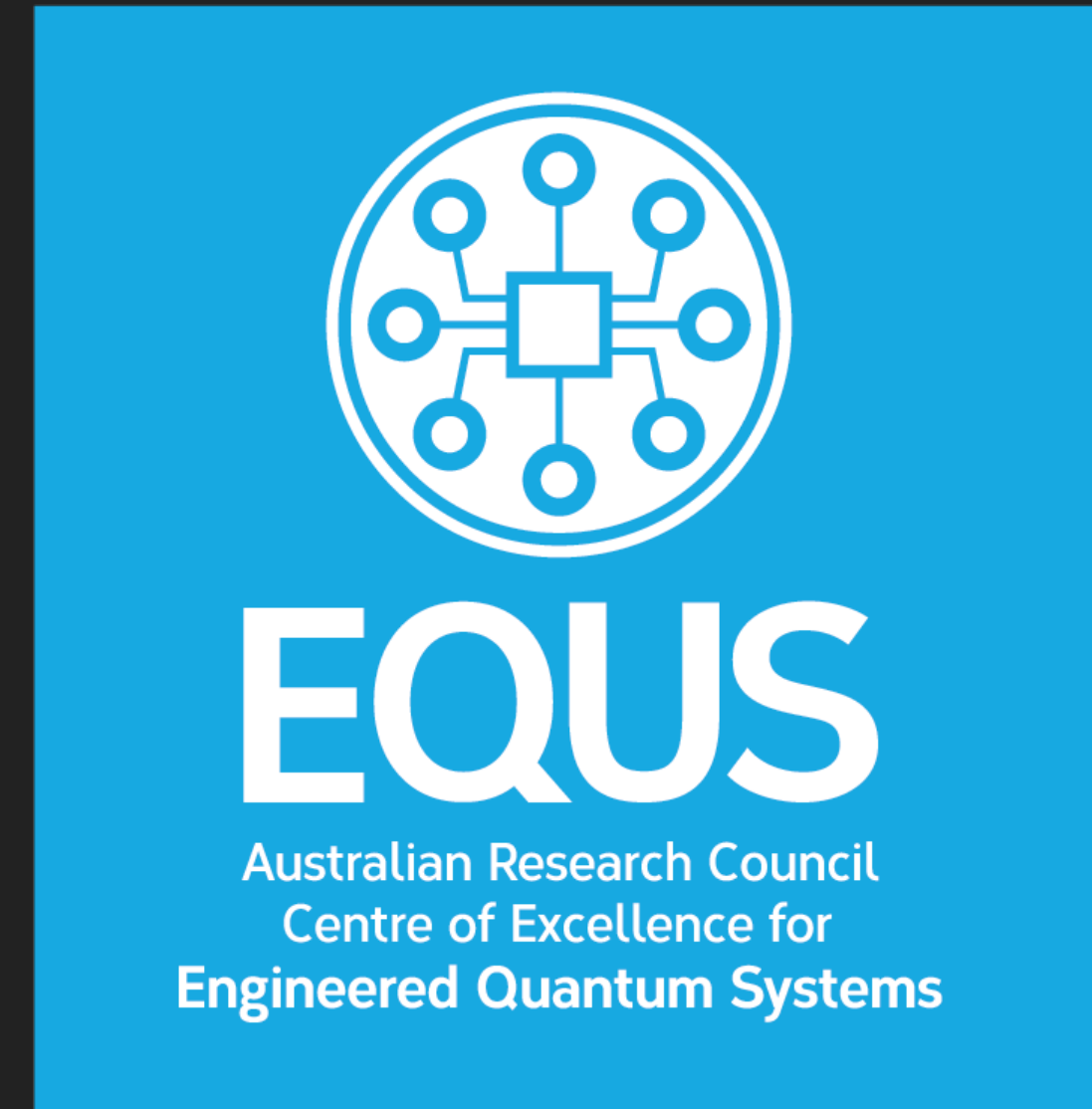


Testing Fundamental Physics and Searching for Dark Matter using Precision Resonators and Oscillators

Prof. Michael Tobar



THE UNIVERSITY OF
**WESTERN
AUSTRALIA**



MG16 2021





EQUS

Australian Research Council
Centre of Excellence for
Engineered Quantum Systems

The QDM Lab: <https://www.qdmlab.com/>
QUANTUM TECHNOLOGIES AND DARK
MATTER RESEARCH LAB



THE UNIVERSITY OF
**WESTERN
AUSTRALIA**



Our Team

HDR/PHD STUDENTS

- Graeme Flower
- Catriona Thomson
- William Campbell
- Aaron Quiskamp
- Elrina Hartman

UNDERGRAD STUDENTS

- Jay Mummery (Masters)
- Robert Crew (BPhil)
- Daniel Tobar (BPhil)
- Michael Hatzon (BPhil)

ACADEMIC

- Michael Tobar
- Eugene Ivanov
- Maxim Goryachev

POSTDOCS

- Ben McAllister
- Cindy Zhao
- Jeremy Bourhill

TECHNICIAN

- Steven Osborne

ADJUNCT

- Alexey Veryaskin (Trinity Labs)

Physics: An experimental science



Computer Scientist
Invented COBOL

White, Harvey Elliott



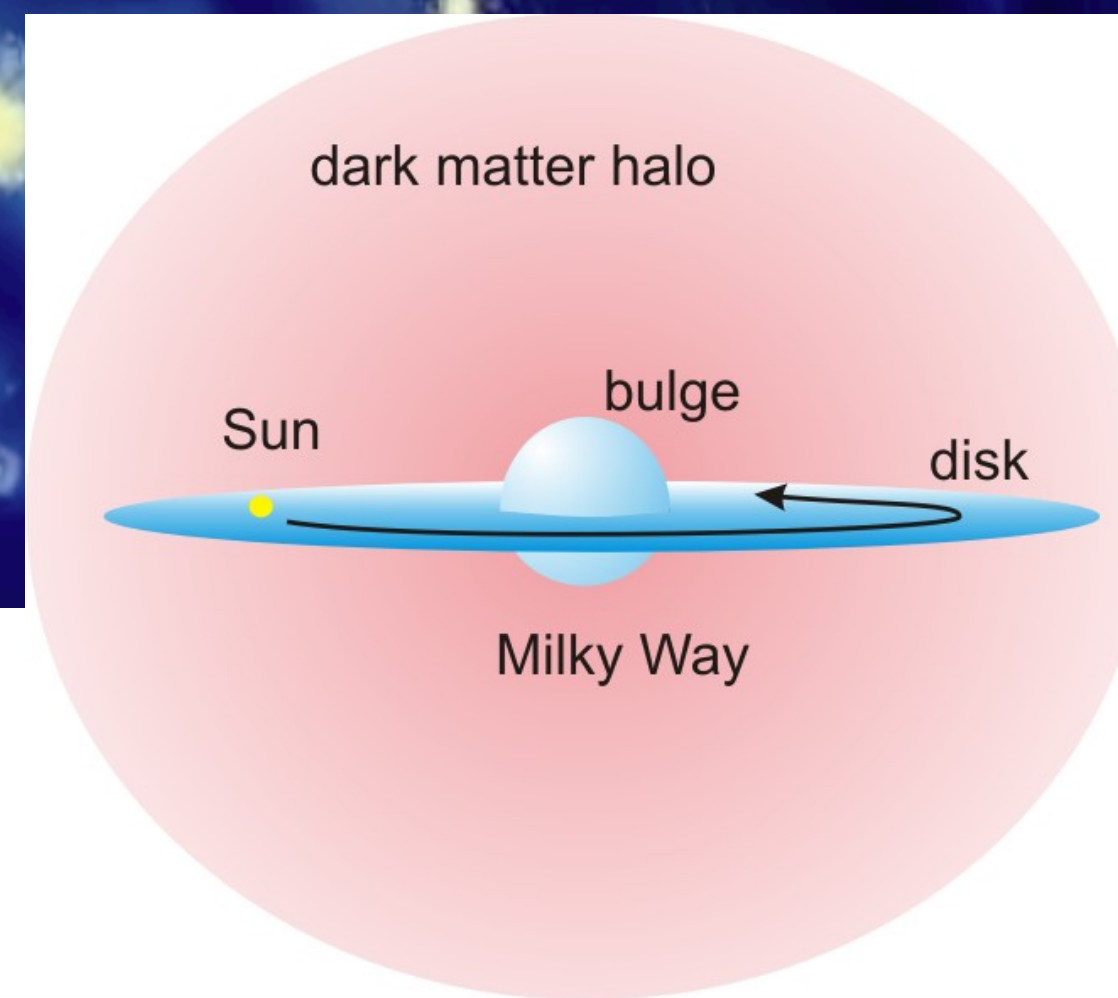
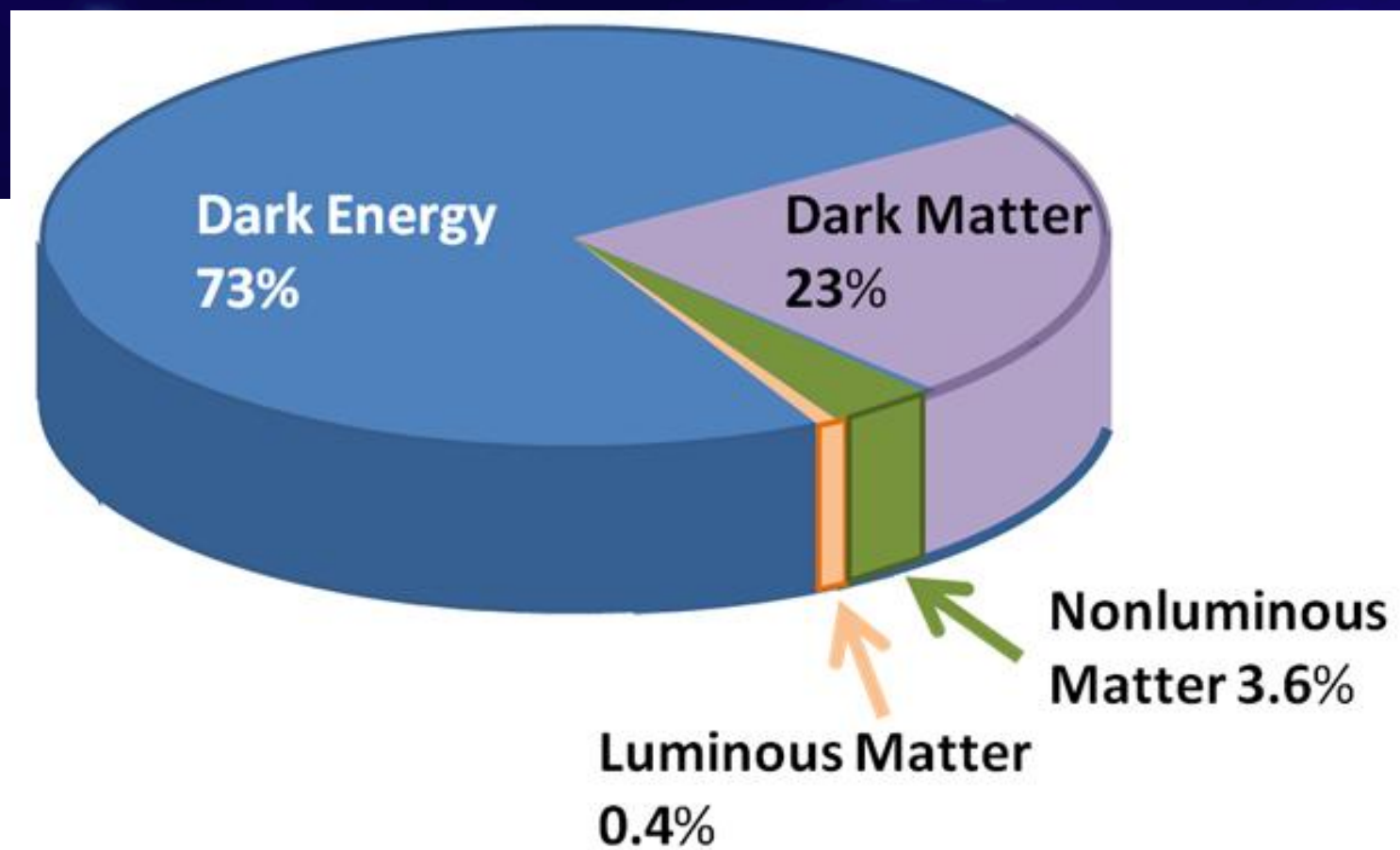
One accurate
measurement is worth
a thousand
expert opinions
Grace Hopper



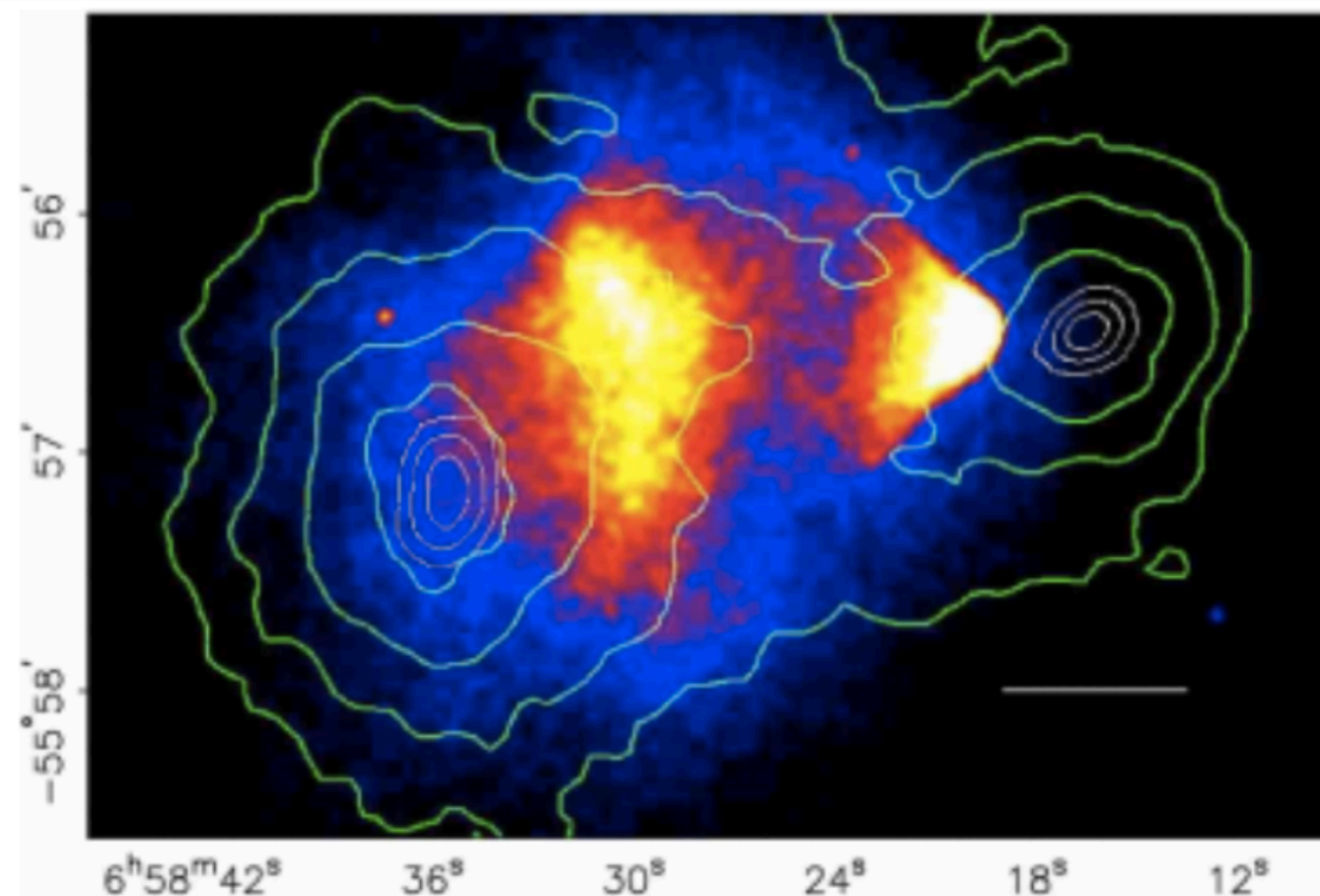
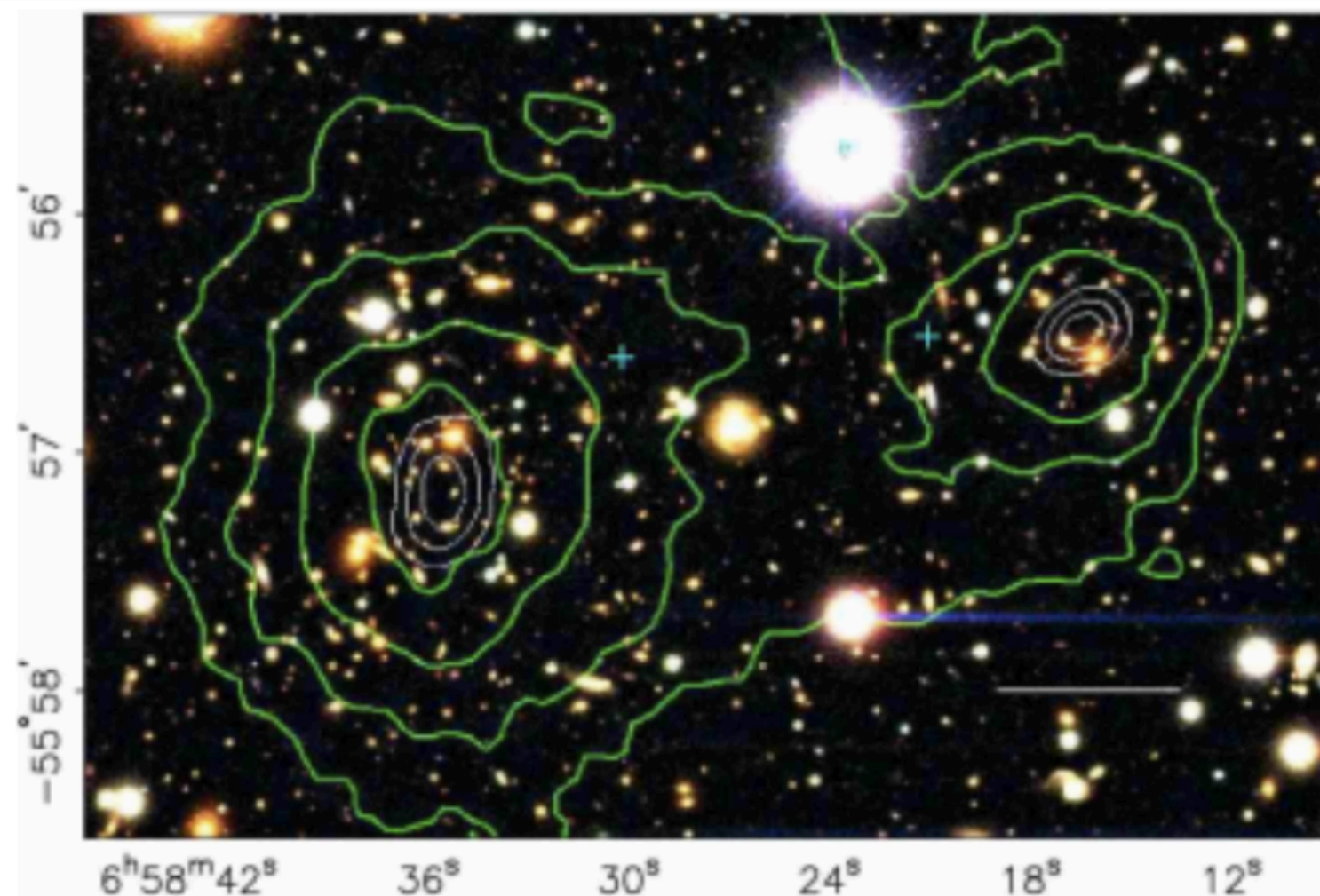
Note: This is not the actual book cover

ARC CENTRE OF EXCELLENCE FOR DARK MATTER PARTICLE PHYSICS

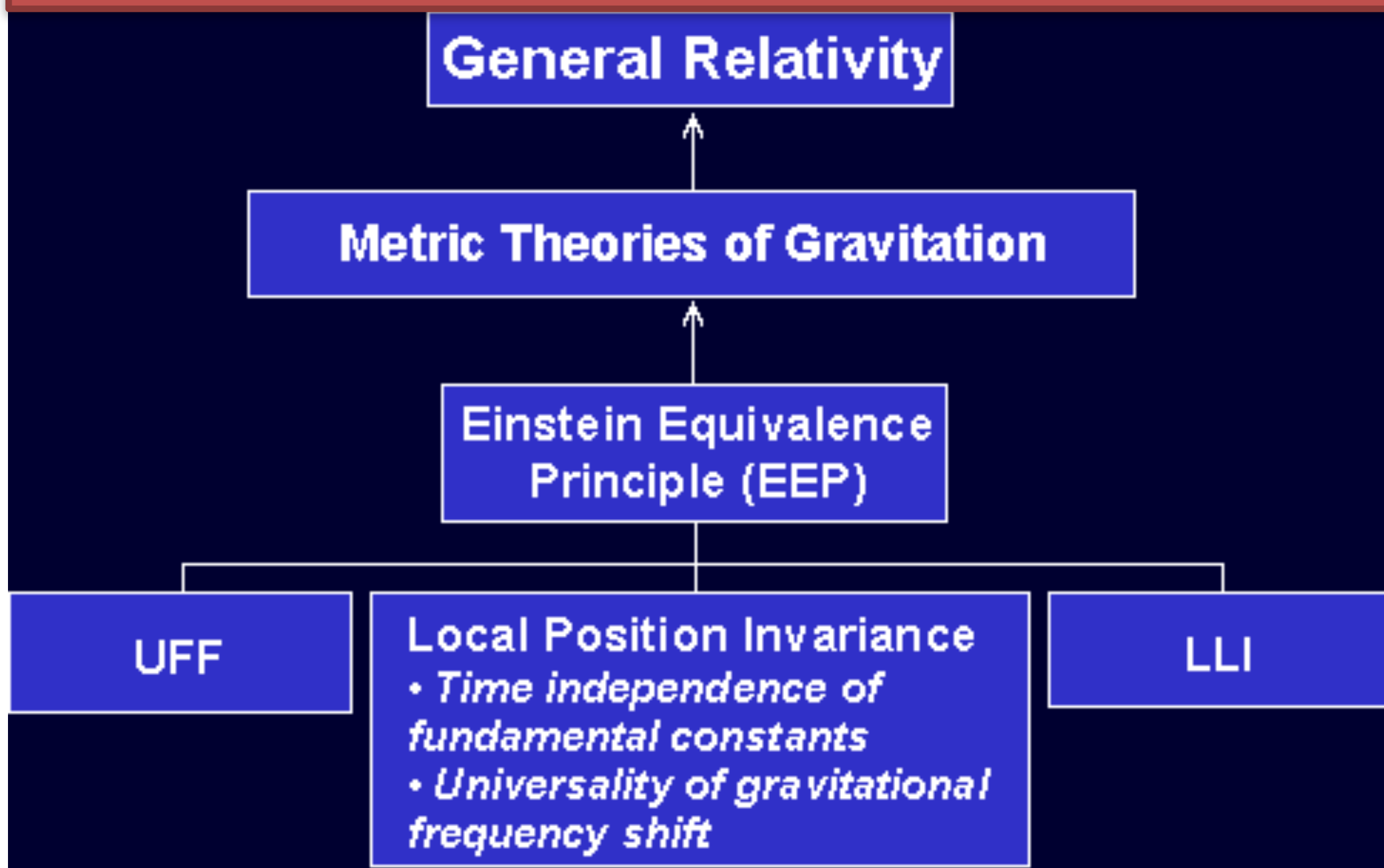
Dark Matter ??



- *Models of dark matter interacting with Standard Model particles.
- *Convert dark matter to standard model particles
- *Use precision measurement to search for dark matter?

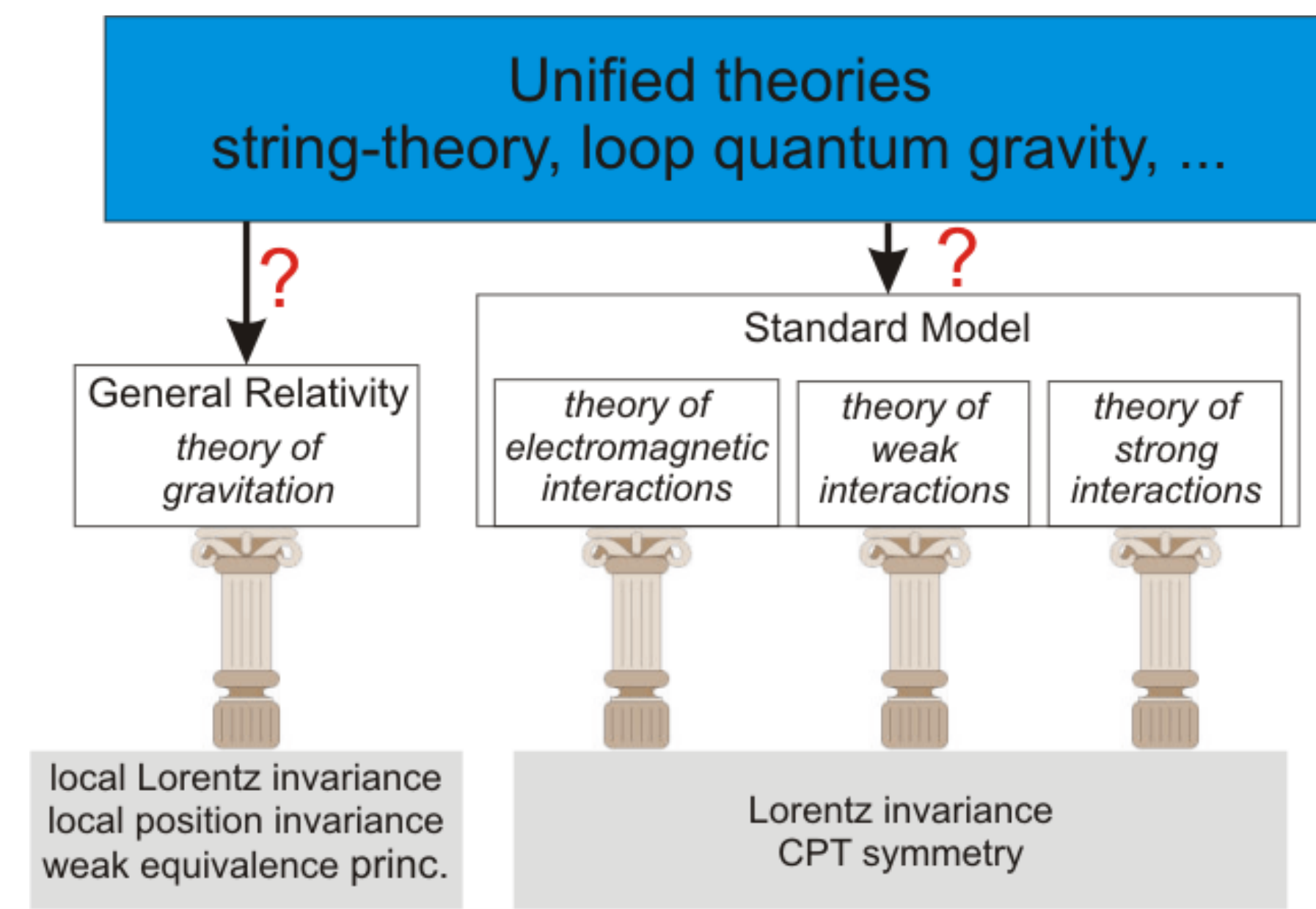


Does Reality Break Down? Is there a Theory of Quantum Gravity?



- *General Relativity rests on the Einstein Equivalence Principle (EEP) which consist of three parts
- *Universality of Free Fall (UFF), also called the Weak Equivalence Principle
- *Local Lorentz Invariance (LLI) which implies the local validity of Special Relativity
- *Local Position Invariance (LPI), which implies the universality of the gravitational red shift

**Approach:
Test Relativity with Precision
Measurement**



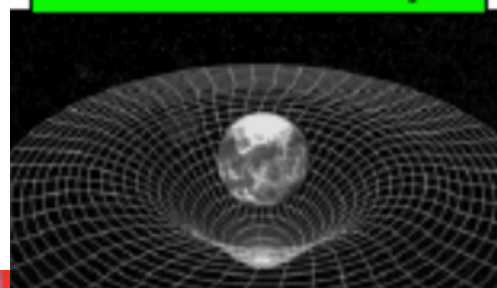


STANDARD MODEL EXTENSION: General Framework for Studying Local Lorentz Violations

SME - effective field theory with lagrangian:

$$\mathcal{L}_{SME} = \mathcal{L}_{GR} + \mathcal{L}_{SM} + \mathcal{L}_{LV}$$

General Relativity



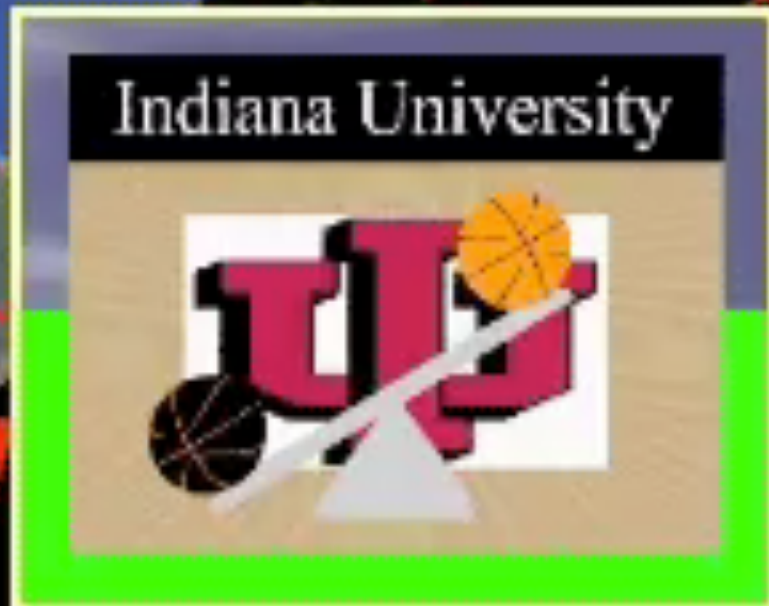
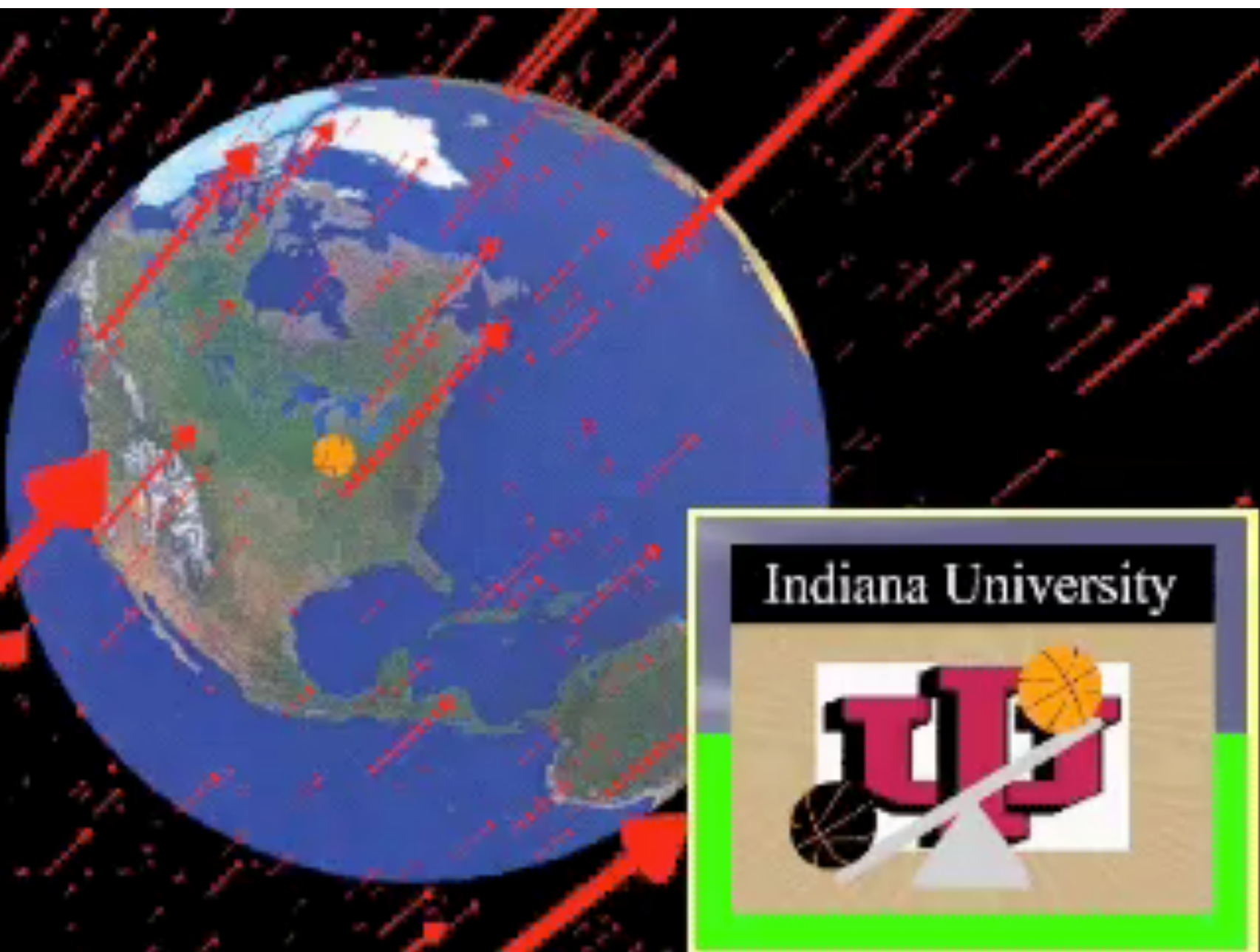
+

Standard Model



+

All possible forms of Lorentz violation
Background fields interacting with known matter



- Photon Sector
- Matter Sector
- Gravity Sector

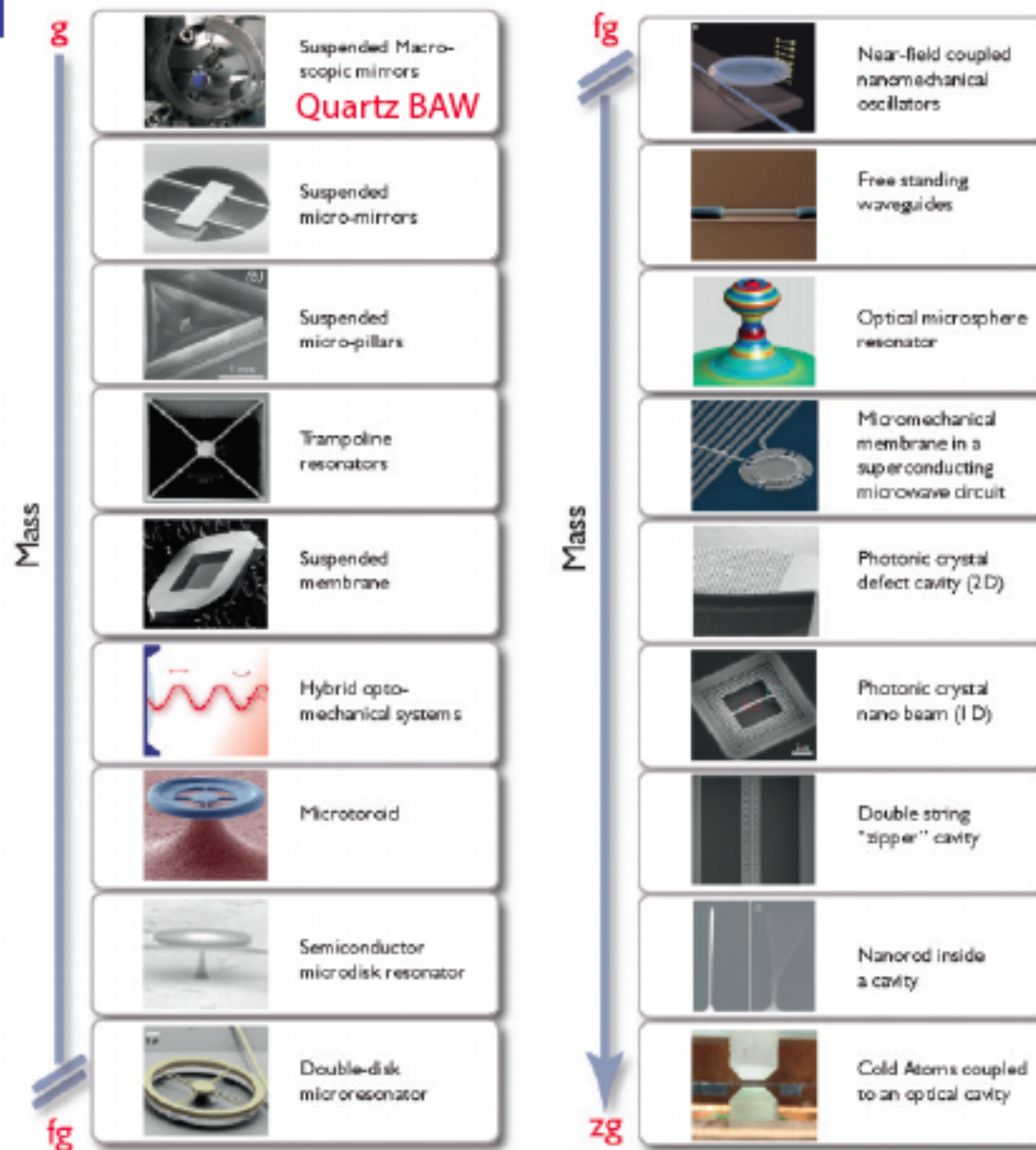
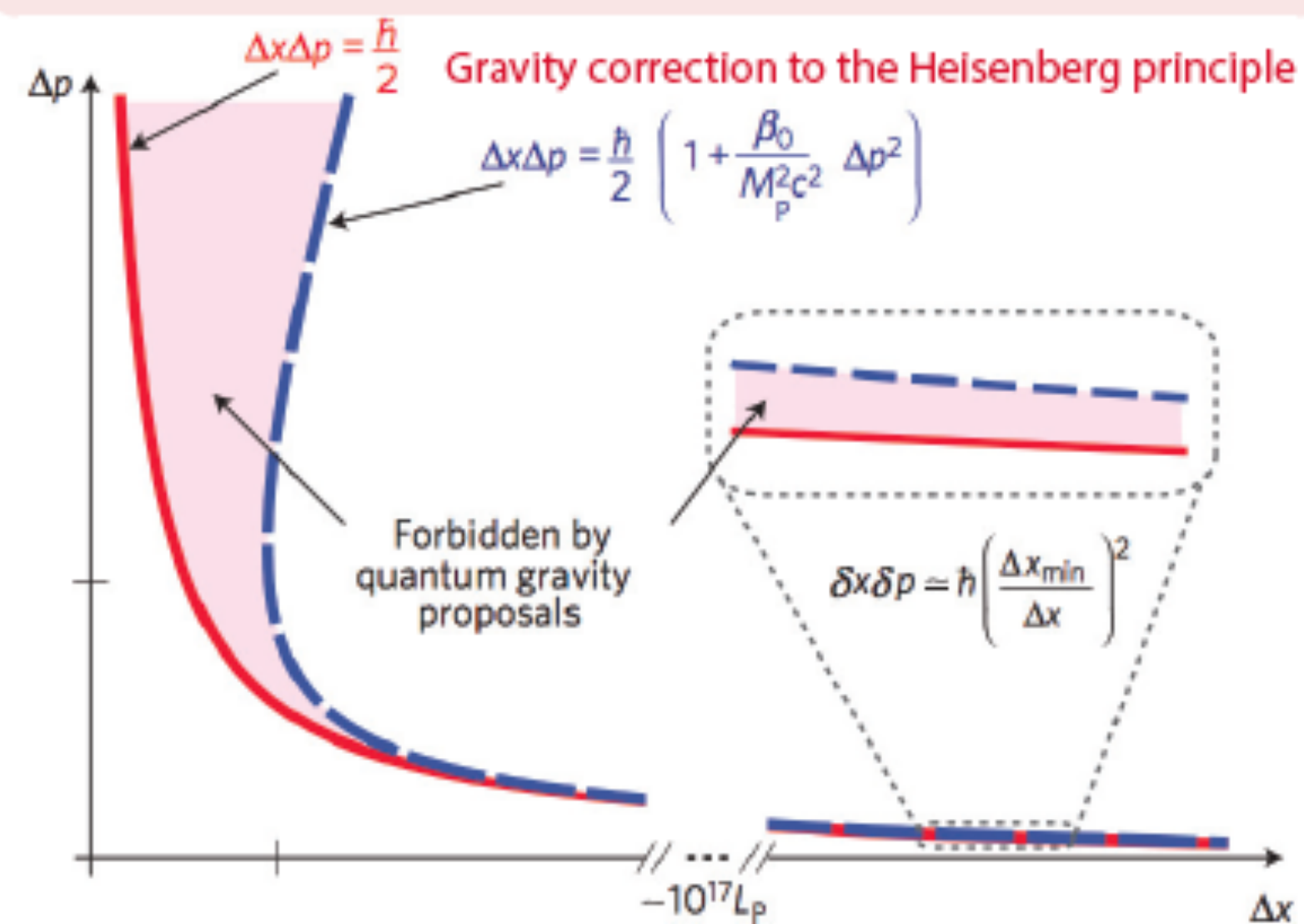
Does Gravity Modify Quantum Mechanics? Is there a Theory of Quantum Gravity?

Quantum Gravity

nature physics **ARTICLES**
 PUBLISHED ONLINE 18 MARCH 2012 | DOI: 10.1038/NPHYS2262
Probing Planck-scale physics with quantum optics
 Igor Pikovski^{1,2*}, Michael R. Vanner^{1,2}, Markus Aspelmeyer^{1,2}, M. S. Kim^{3*} and Časlav Brukner^{2,4}

Theoretically proposed modification:

$$[x, p] = i\hbar \sqrt{1 + 2\eta_0 \frac{(p/c)^2 + m^2}{M_p^2}}$$



Quantum Mechanics Possible with Macroscopic Masses

- *Does mass modify the position and momentum commutator relations?
- *QG theories predict a nonlinear correction to the canonical commutation relation between x , and p .
- *Can test with precision opto-mechanics

Outline

Experimental Techniques

(Clock Zoo)

Dark Matter Experiments

(Wave Like Dark Matter; Axions; Acoustic Oscillators)

(Fermi Lab)

Lorentz Invariance with Phonons

(Search for violations of Lorentz Invariance)

(UC Berkeley)

Quantum Gravity

(Tests of Generalised Uncertainty principle)

(FZ Jülich, U Maastricht)

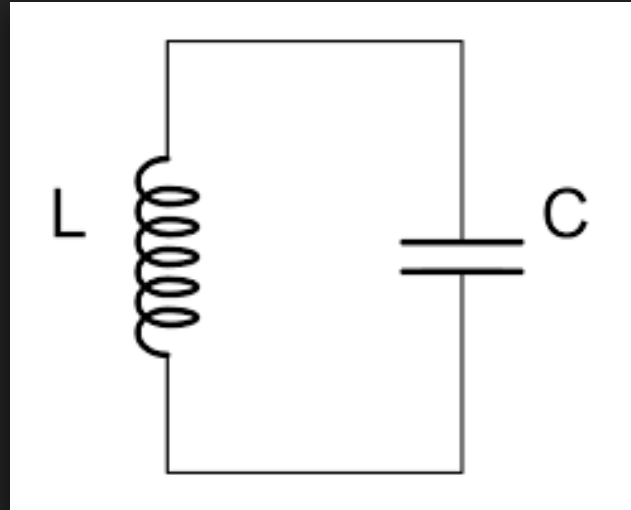
HFGW

(Search for high frequency GW)

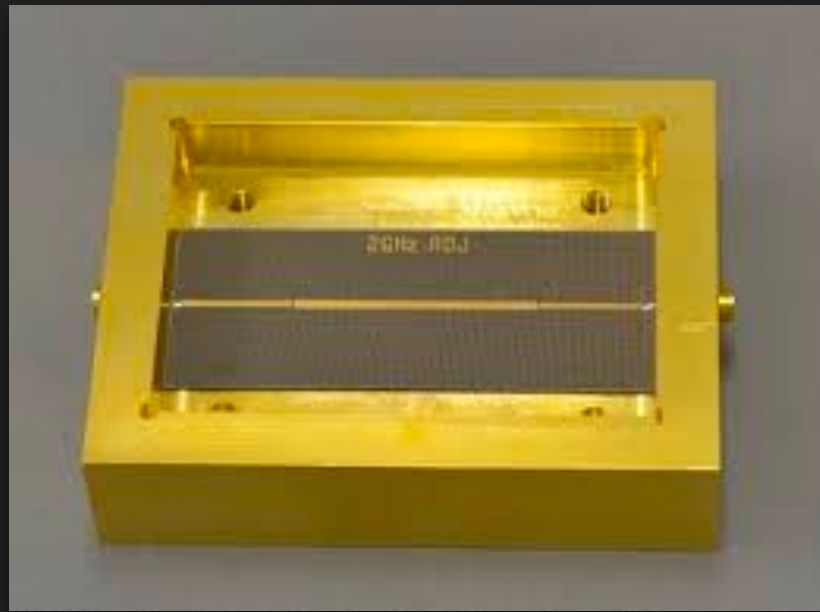
(U Glasgow, FEMTO-ST)

Frequency Metrology

PHOTONS



LC-circuits

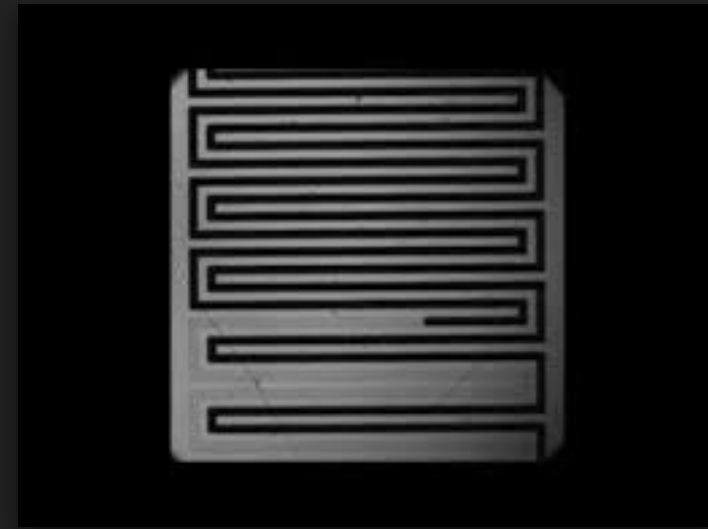


Metallic Cavities



Dielectric Cavities

PHONONS



SAW

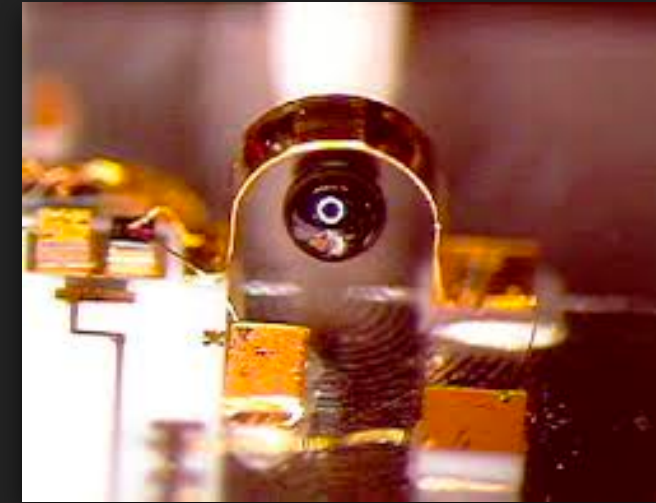


BAW

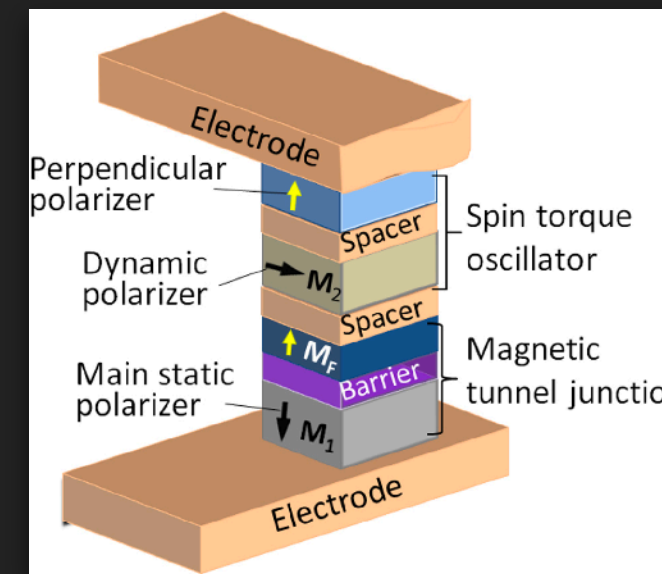


Structures

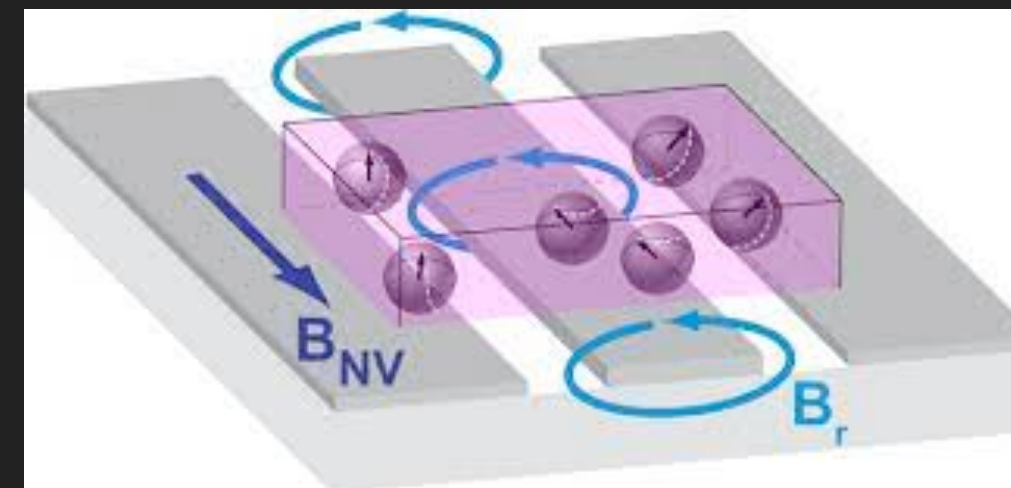
MAGNONS/SPINS



Bulk

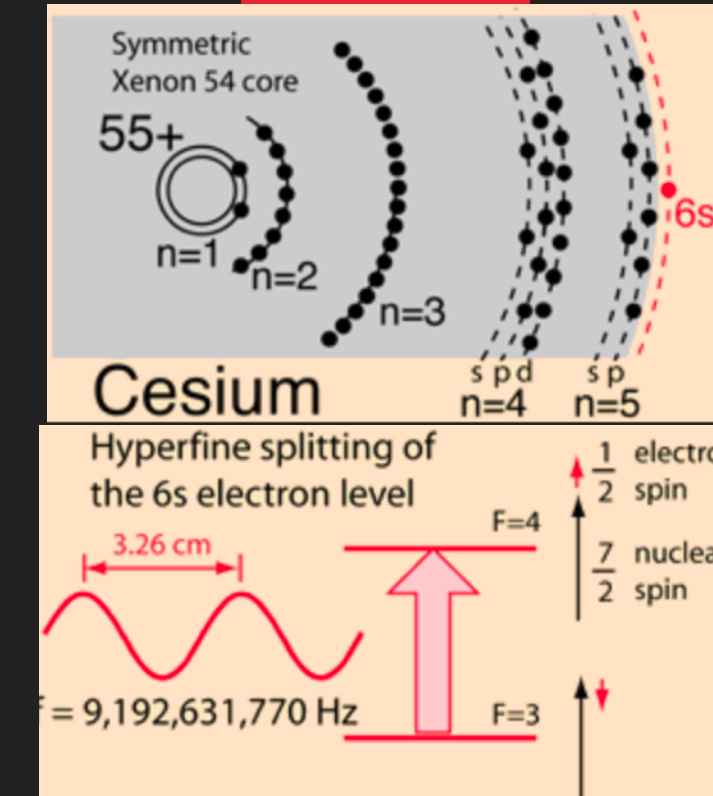


Spin-Torque

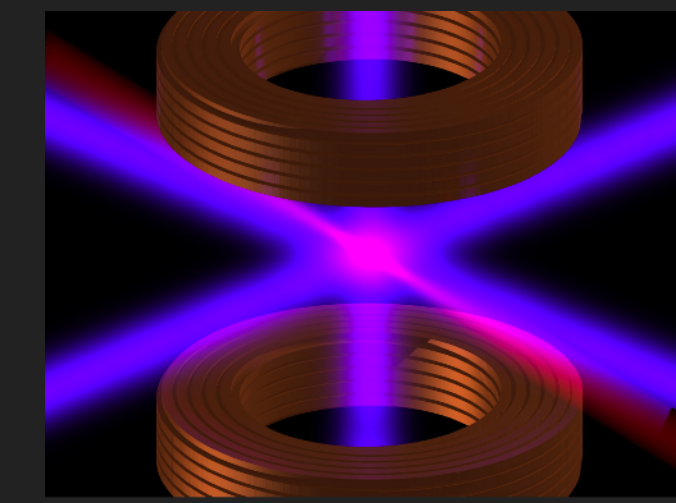


Spin Ensembles

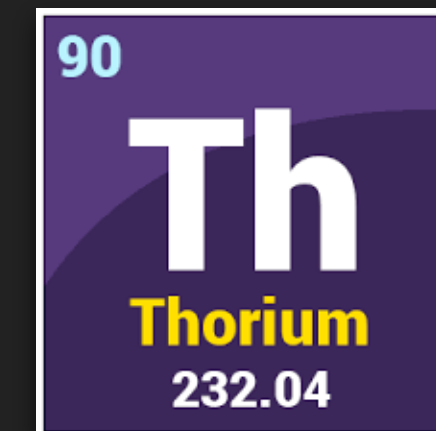
ATOMS



Hyperfine transitions



Electron transitions



Nuclear transitions

WIDE RANGE OF PHYSICAL PHENOMENA

Searching for Putative Wave-Like Dark Matter

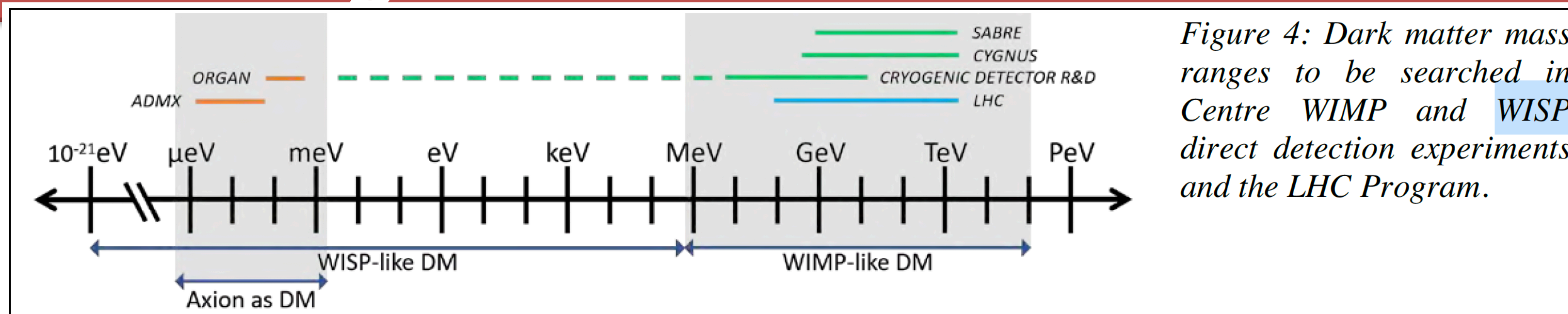


Figure 4: Dark matter mass ranges to be searched in Centre WIMP and WISP direct detection experiments and the LHC Program.

Generic Experiment

Design Physics Package:

Wave like Dark Matter

-> Sensitive to the type of Dark Matter of Interest

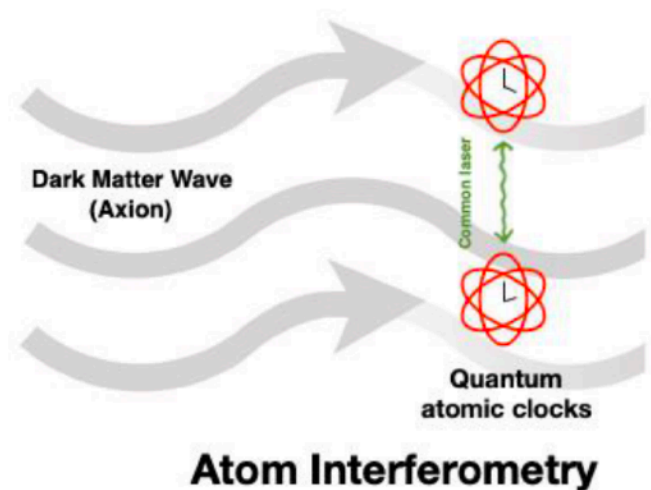
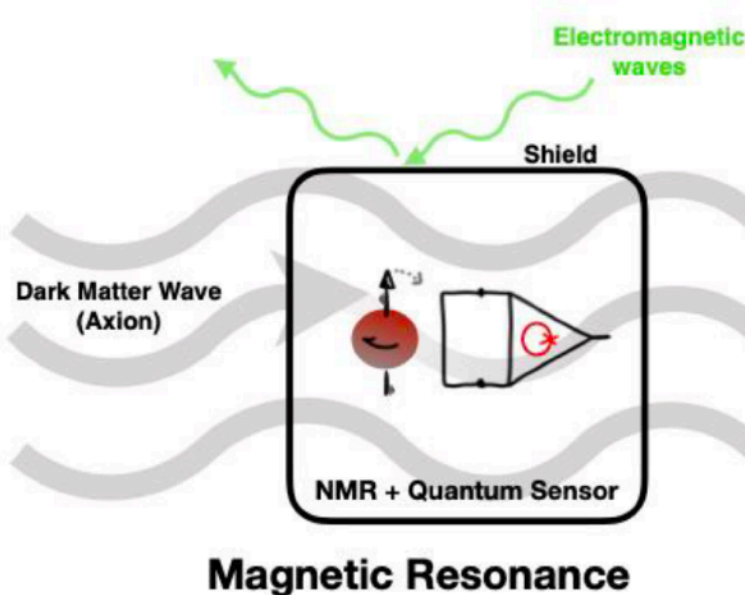
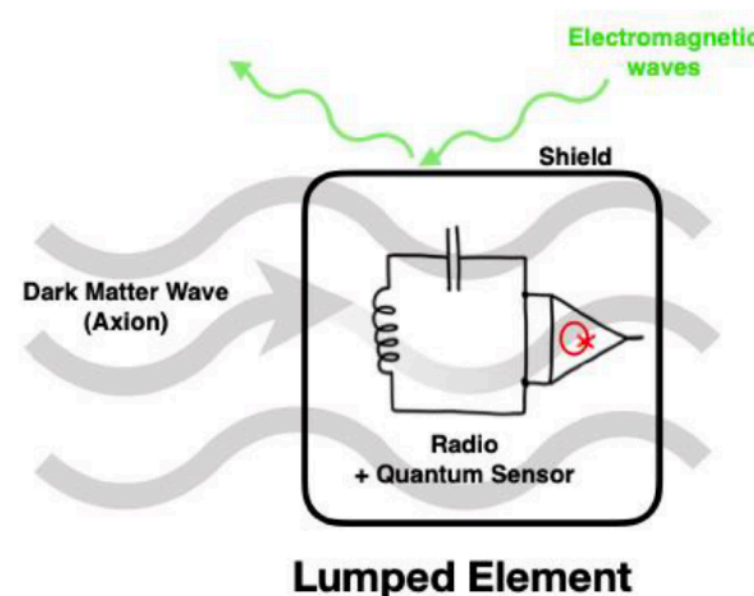
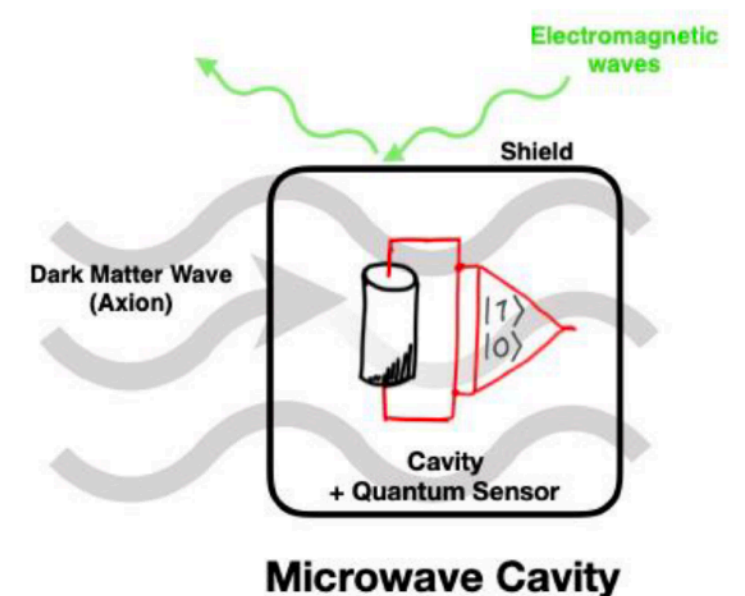
-> Axion, Dilaton etc. (bosonic)

-> Theory interacts with Experiment: How Dark Matter interacts with Standard Model Particles, Optimise Signal

-> Reduce Noise, Fundamental Limit is Quantum Noise

-> Surpass Quantum Limit: Quantum Metrology

-> Broadband and resonant searches via cavities, circuit oscillators, NMR etc



WAVE LIKE DARK MATTER PROGRAM @ UWA

- (1) Axion Dark Matter eXperiment (**ADMX**) Project run by **Fermilab**, run out of **Seattle** at Washington University. **UWA** Officially a group member since 2019. **PI Gray Rybka**
- (2) **Oscillating Resonant Group AxioN experiment (ORGAN)**. The first Axion experiment **at UWA**, will test Axion Cogenesis in 2021, arXiv:2006.04809 [hep-ph]
- (3) Searches for axions through coupling with electron spins; **Magnon-Cavity UWA**
- (4) Low mass detectors for axions with LCR Circuits, **ADMX-SLIC** (Superconducting LC-circuit Investigating Cold axions) **UF** (Sikivie and Tanner) and Broadband Electrical Action Sensing Technique (**BEAST**) **UWA**
- (5) AC Haloscope with Low Noise Oscillators (UPconversion Loop Oscillator Axion Detector (**UPLOAD**) **UWA**
- (6) **Light Scalar Dark Matter** (Dilaton) Clock Comparisons, Acoustic Detectors **UWA**

(2) ORGAN: Cavity Haloscope in DC B-field Dielectric-Boosted Sensitivity to Cylindrical Azimuthally Varying Transverse-Magnetic Resonant Modes in an Axion Haloscope

Aaron P. Quiskamp^{1,*}, Ben T. McAllister¹, Gray Rybka², and Michael E. Tobar^{1,†}

¹ARC Centre of Excellence for Engineered Quantum Systems and ARC Centre of Excellence for Dark Matter Particle Physics, Department of Physics, University of Western Australia, 35 Stirling Highway, Crawley WA 6009, Australia

²Centre for Experimental Nuclear Physics and Astrophysics, University of Washington, 1410 NE Campus Parkway, Seattle, Washington 98195, USA

(Received 15 June 2020; revised 6 August 2020; accepted 28 September 2020; published 27 October 2020)

Axions are a popular dark-matter candidate that are often searched for in experiments known as “haloscopes,” which exploit a putative axion-photon coupling. These experiments typically rely on transverse-magnetic (TM) modes in resonant cavities to capture and detect photons generated via axion conversion. We present a study of a resonant-cavity design for application in haloscope searches, of particular use in the push to higher-mass axion searches (above approximately $60 \mu\text{eV}$). In particular, we take advantage of azimuthally varying TM_{m10} modes that, while typically insensitive to axions due to field nonuniformity, can be made axion sensitive (and frequency tunable) through the strategic placement of dielectric wedges, becoming a type of resonator known as a dielectric-boosted axion-sensitivity (DBAS) resonator. Results from finite-element modeling are presented and compared with a simple proof-of-concept experiment. The results show a significant increase in axion sensitivity for these DBAS resonators over their empty-cavity counterparts and high potential for application in high-mass axion searches when benchmarked against simpler more traditional designs that rely on fundamental TM modes.

DOI: 10.1103/PhysRevApplied.14.044051

PHYSICAL REVIEW LETTERS 126, 081803 (2021)

(5) UPLOAD: Two Modes in one Cavity

Upconversion Loop Oscillator Axion Detection Experiment: A Precision Frequency Interferometric Axion Dark Matter Search with a Cylindrical Microwave Cavity

Catriona A. Thomson^{1,*}, Ben T. McAllister, Maxim Goryachev, Eugene N. Ivanov, and Michael E. Tobar^{1,†}
ARC Centre of Excellence for Engineered Quantum Systems and ARC Centre of Excellence for Dark Matter Particle Physics,
Department of Physics, University of Western Australia, 35 Stirling Highway, Crawley, Western Australia 6009, Australia

(Received 5 January 2020; revised 11 November 2020; accepted 15 January 2021; published 23 February 2021)

First experimental results from a room-temperature tabletop phase-sensitive axion haloscope experiment are presented. The technique exploits the axion-photon coupling between two photonic resonator oscillators excited in a single cavity, allowing low-mass axions to be upconverted to microwave frequencies, acting as a source of frequency modulation on the microwave carriers. This new pathway to axion detection has certain advantages over the traditional haloscope method, particularly in targeting axions below $1 \mu\text{eV}$ (240 MHz) in energy. At the heart of the dual-mode oscillator, a tunable cylindrical microwave cavity supports a pair of orthogonally polarized modes ($\text{TM}_{0,2,0}$ and $\text{TE}_{0,1,1}$), which, in general, enables simultaneous sensitivity to axions with masses corresponding to the sum and difference of the microwave frequencies. However, in the reported experiment, the configuration was such that the sum frequency sensitivity was suppressed, while the difference frequency sensitivity was enhanced. The results place axion exclusion limits between 7.44–19.38 neV, excluding a minimal coupling strength above $5 \times 10^{-7} \text{ 1/GeV}$, after a measurement period of two and a half hours. We show that a state-of-the-art frequency-stabilized cryogenic implementation of this technique, ambitious but realizable, may achieve the best limits in a vast range of axion space.

DOI: 10.1103/PhysRevLett.126.081803



(4) LCR Circuits in DC Magnetic Field

Physics of the Dark Universe 30 (2020) 100624



ELSEVIER

Contents lists available at ScienceDirect

Physics of the Dark Universe

journal homepage: www.elsevier.com/locate/dark



Broadband electrical action sensing techniques with conducting wires for low-mass dark matter axion detection

Michael E. Tobar^{*}, Ben T. McAllister, Maxim Goryachev

ARC Centre of Excellence For Engineered Quantum Systems, Department of Physics, University of Western Australia, 35 Stirling Highway, Crawley WA 6009, Australia



(6) SCALAR DARK MATTER: Includes Phonons

PHYSICAL REVIEW LETTERS 126, 071301 (2021)

Searching for Scalar Dark Matter via Coupling to Fundamental Constants with Photonic, Atomic, and Mechanical Oscillators

William M. Campbell¹, Ben T. McAllister, Maxim Goryachev, Eugene N. Ivanov², and Michael E. Tobar^{1,*}
ARC Centre of Excellence for Engineered Quantum Systems and ARC Centre of Excellence for Dark Matter Particle Physics,
Department of Physics, University of Western Australia, 35 Stirling Highway, Crawley, Western Australia 6009, Australia

(Received 16 October 2020; revised 25 November 2020; accepted 15 January 2021; published 18 February 2021)

We present a way to search for light scalar dark matter (DM), seeking to exploit putative coupling between dark matter scalar fields and fundamental constants, by searching for frequency modulations in direct comparisons between frequency stable oscillators. Specifically we compare a cryogenic sapphire oscillator (CSO), hydrogen maser (HM) atomic oscillator, and a bulk acoustic wave quartz oscillator (OCXO). This work includes the first calculation of the dependence of acoustic oscillators on variations of the fundamental constants, and demonstration that they can be a sensitive tool for scalar DM experiments. Results are presented based on 16 days of data in comparisons between the HM and OCXO, and 2 days of comparison between the OCXO and CSO. No evidence of oscillating fundamental constants consistent with a coupling to scalar dark matter is found, and instead limits on the strength of these couplings as a function of the dark matter mass are determined. We constrain the dimensionless coupling constant d_e and combination $|d_{m_e} - d_g|$ across the mass band $4.4 \times 10^{-19} \lesssim m_\phi \lesssim 6.8 \times 10^{-14} \text{ eV } c^{-2}$, with most sensitive limits $d_e \gtrsim 1.59 \times 10^{-1}$, $|d_{m_e} - d_g| \gtrsim 6.97 \times 10^{-1}$. Notably, these limits do not rely on maximum reach analysis (MRA), instead employing the more general coefficient separation technique. This experiment paves the way for future, highly sensitive experiments based on state-of-the-art acoustic oscillators, and we show that these limits can be competitive with the best current MRA-based exclusion limits.

DOI: 10.1103/PhysRevLett.126.071301



Searching for Scalar Dark matter from Oscillating Fundamental Constants

Atomic ionization by scalar dark matter and solar scalars

H. B. Tran Tan, A. Derevianko, V. Dzuba, V.V. Flambaum

hhu.



Limits on oscillating fundamental constants from laser spectroscopy of molecular ensembles

R. Oswald, A. Nevsky, V. Vogt, S. Schiller
Heinrich-Heine-Universität Düsseldorf (Germany)

N. L. Figueroa, Ke Zhang, O. Tretiak, D. Antypas, D. Budker*
Johannes Gutenberg-Universität Mainz (Germany)

Helmholtz-Institut, GSI Helmholtzzentrum für Schwerionenforschung (Germany)

A. Banerjee, G. Perez

Department of Particle Physics and Astrophysics, Weizmann Institute of Science (Israel)

* also: Department of Physics, University of California, Berkeley (USA)

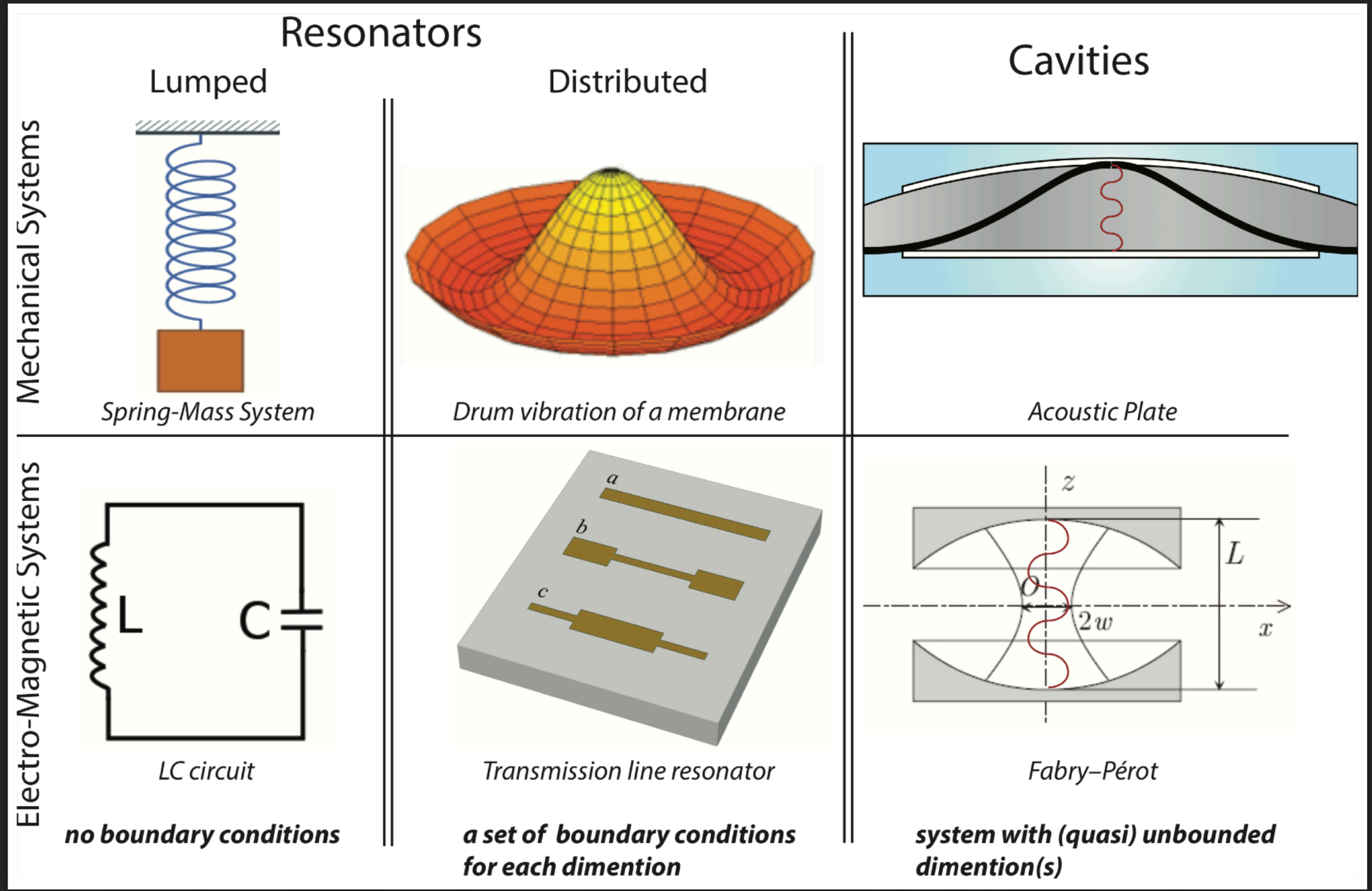
Some dependencies

- Optical transition frequency $f \propto m_e c^2 \alpha^2 H(\alpha)$
- Hyperfine transition frequency $f \propto m_e c^2 \alpha^4 F(\alpha) \left(\frac{m_e}{m_p}\right) \mu_{nuc}$
- Molecular vibrational transition frequency $f \propto m_e c^2 \alpha^2 \left(\frac{m_e}{M_{nuc}}\right)^{\frac{1}{2}} G\left(\frac{m_e}{M_{nuc}}\right)$
- Electromagnetic cavity mode frequency (empty cavity) $f \propto m_e c^2 \alpha$
- Mechanical mode frequency $f \propto m_e c^2 \alpha^2 \left(\frac{m_e}{M_{nuc}}\right)^{\frac{1}{2}}$

$$\frac{\delta M_{nuc}}{M_{nuc}} = \frac{\delta \Lambda_{QCD}}{\Lambda_{QCD}} + 0.093 \frac{\delta \hat{m}}{\hat{m}} + 0.043 \frac{\delta m_s}{m_s}, \quad \hat{m} = (m_u + m_d)/2$$

μ_{nuc} has a small/modest dependence on m_s (~ 0.01) and on \hat{m} (~ 0.1)

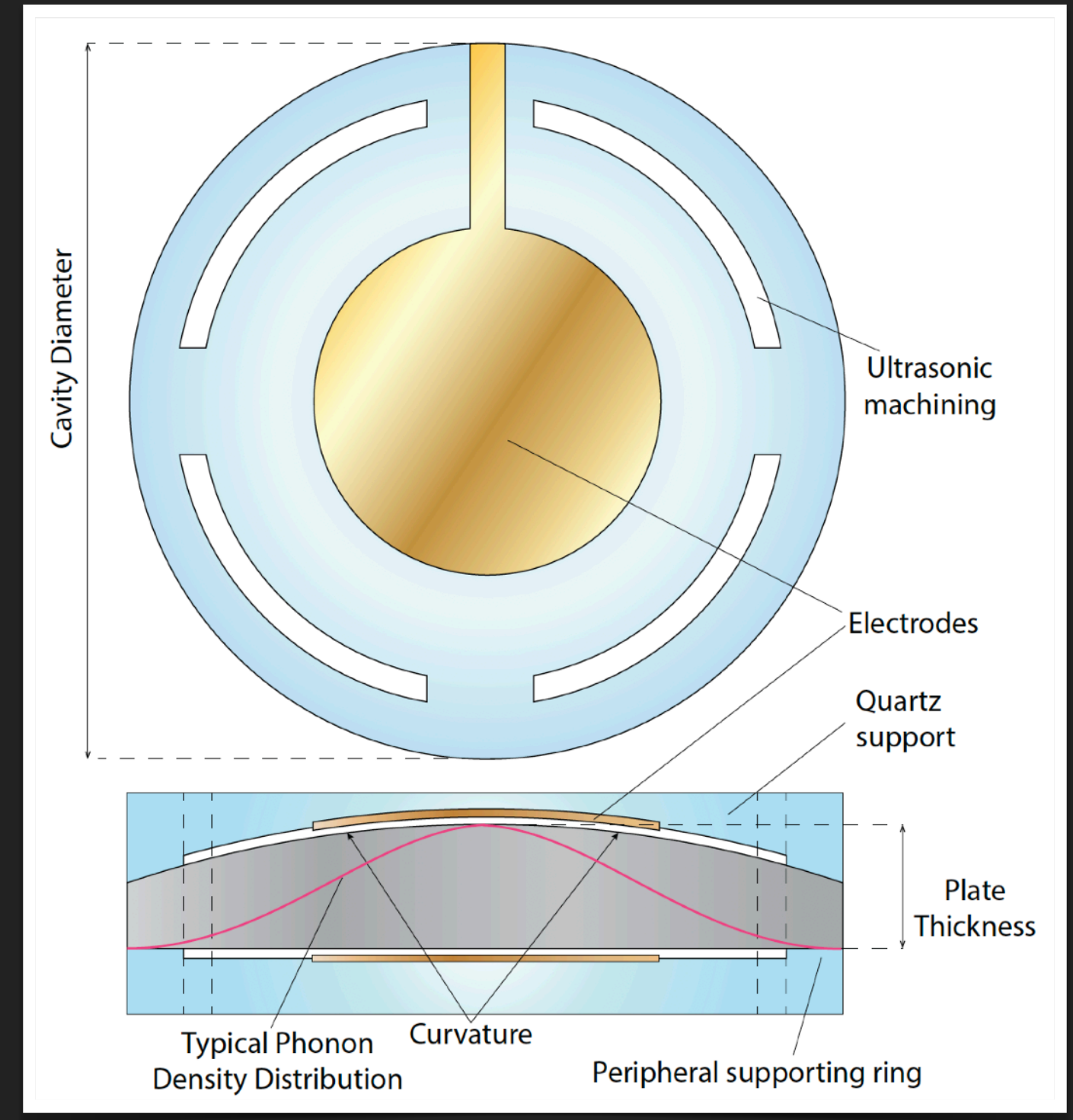
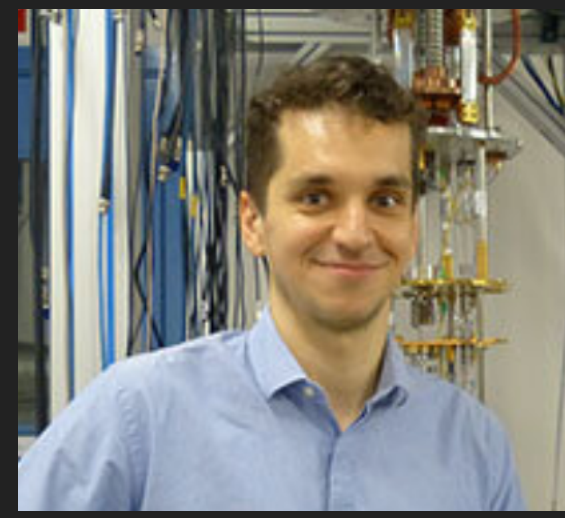
Photons (Electromagnetic) vs Phonons (Acoustic) BAW Cavities



APPLIED PHYSICS LETTERS 100, 243504 (2012)

PRL 111, 085502 (2013)

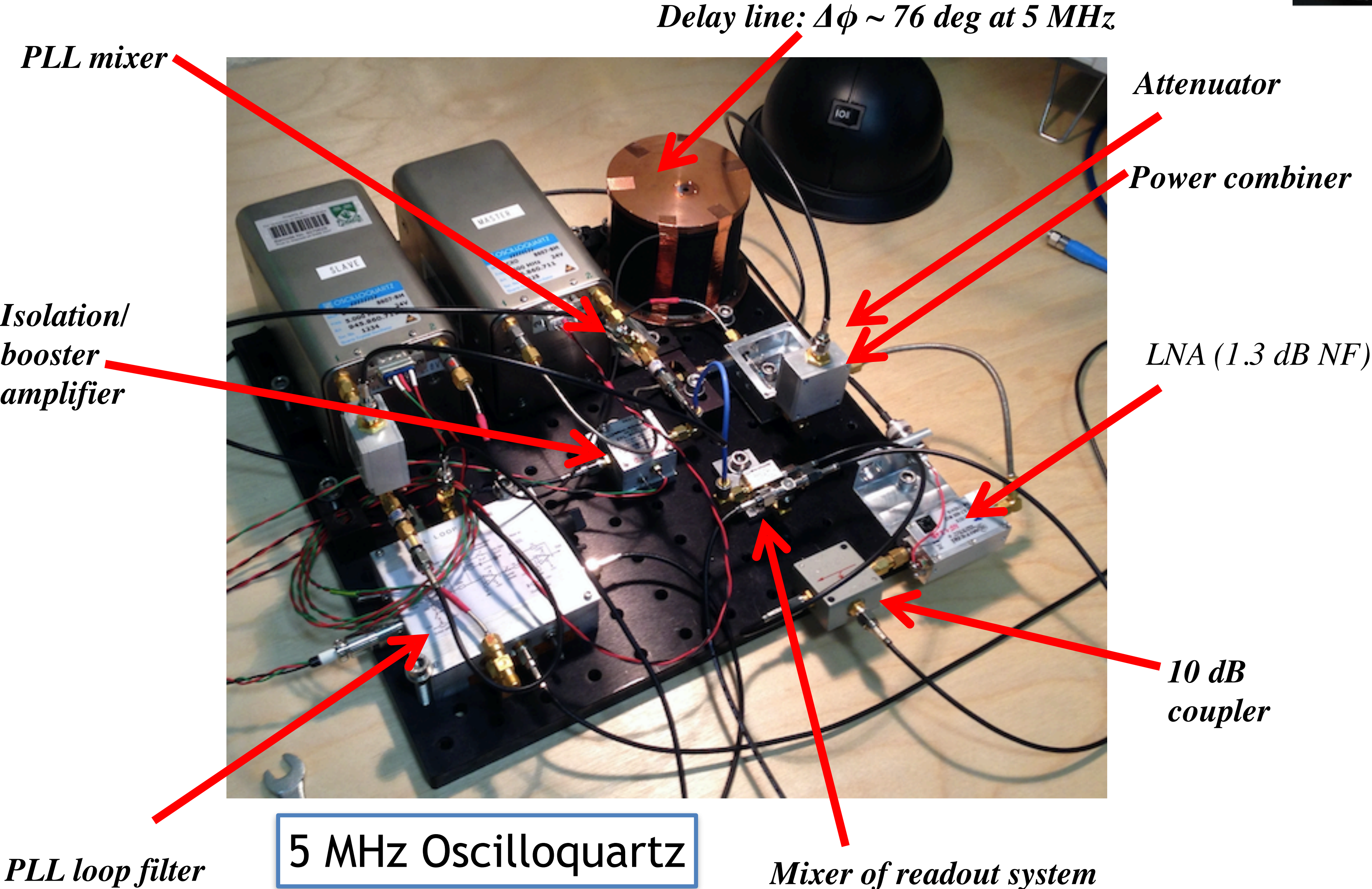
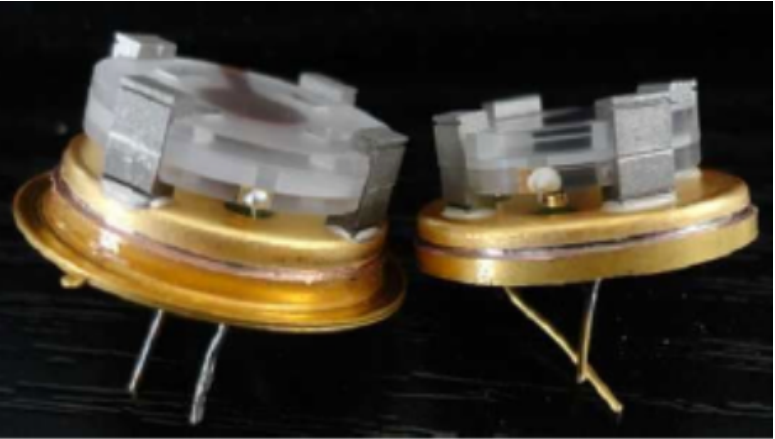
Scientific Reports Vol. 3, 2132 (2013)



- * Frequency range: 1-1000 MHz
- * Three mode family types: 2 transverse and 1 longitudinal
- * Piezoelectric Coupling
- * Established technology (>70 years for time keeping applications)
- * Record high Quality factors $\sim 10^{10}$

Room Temperature Resonator-Oscillators using Quartz BAWs

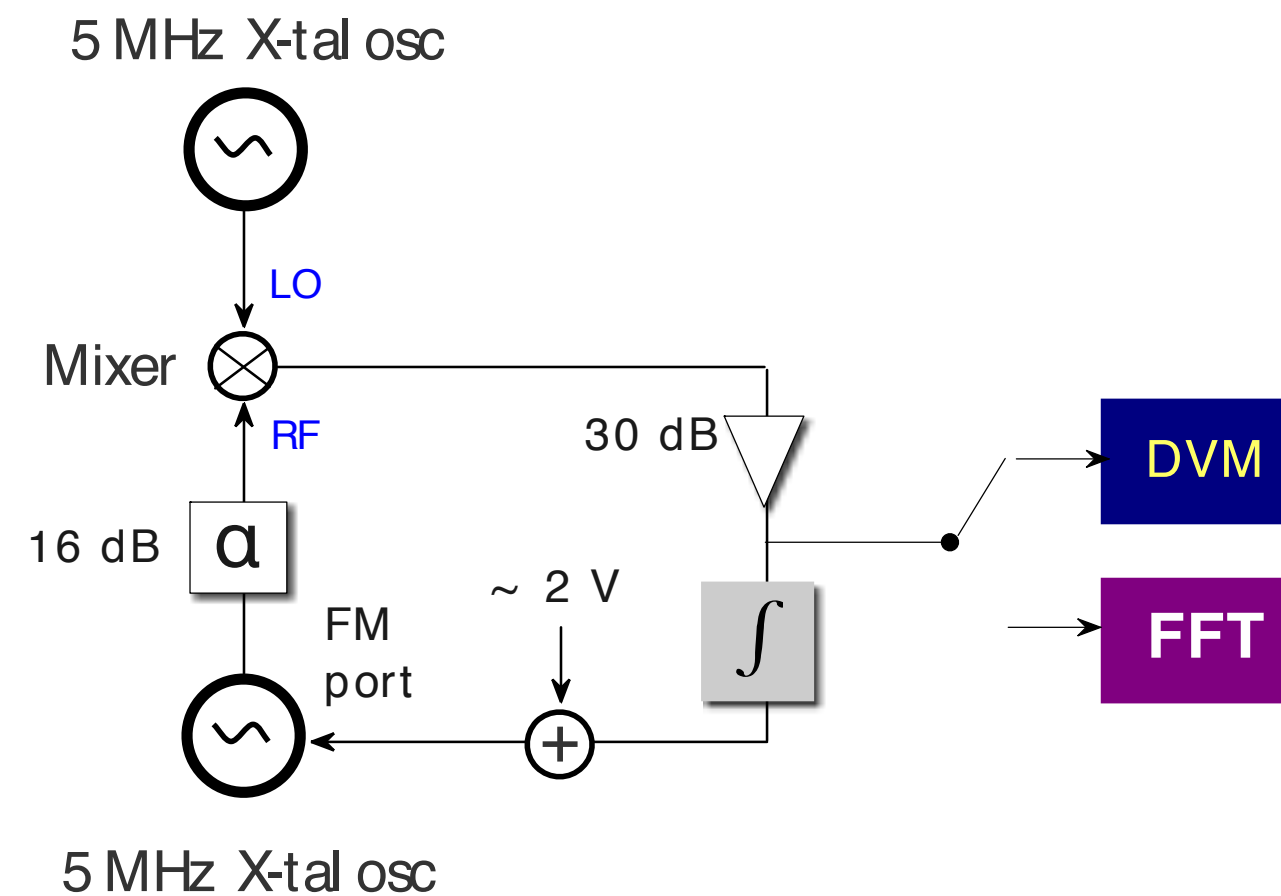
$Q \sim 10^6$



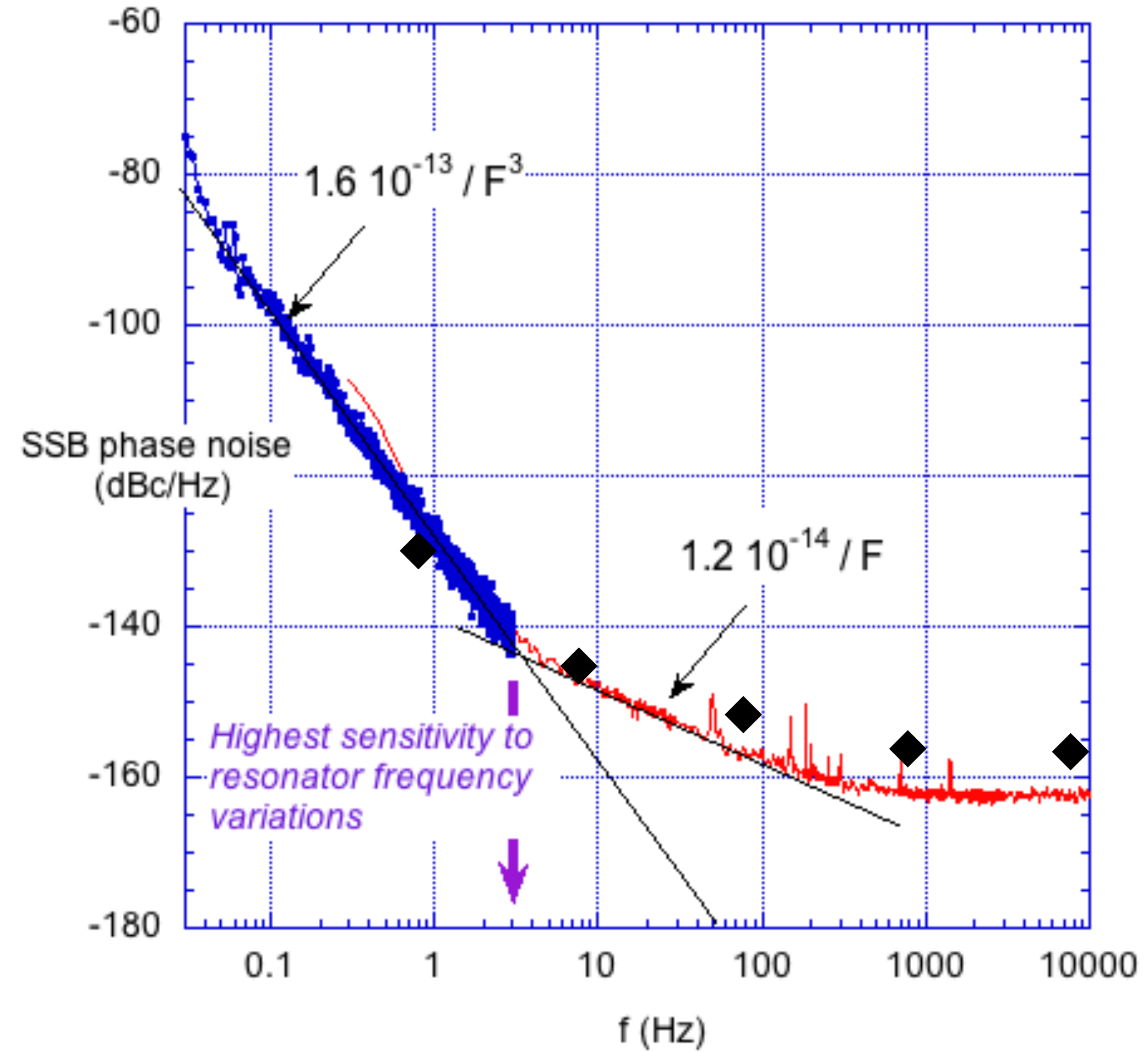
Phase Noise Spectrum of 5 MHz Oscilloquartz

Phase noise (BW = 1 Hz) Options

Frequencies	5 MHz		10 MHz	
	Standard	Option L	Standard	Option L
Phase noise 1 Hz	-125 dBc	-130 dBc	-118 dBc	-122 dBc
10 Hz	-145 dBc	-145 dBc	-137 dBc	-137 dBc
100 Hz	-153 dBc	-153 dBc	-143 dBc	-143 dBc
1'000 Hz	-156 dBc	-156 dBc	-145 dBc	-145 dBc
10'000 Hz	-156 dBc	-156 dBc	-145 dBc	-145 dBc



SSB phase noise of individual 5 MHz Oscilloquartz



Used to test fundamental Physics - LLI, Dark Matter
Can we Improve Using a Cryogenic Oscillator???

Sapphire Low Noise Oscillators under Development at UWA

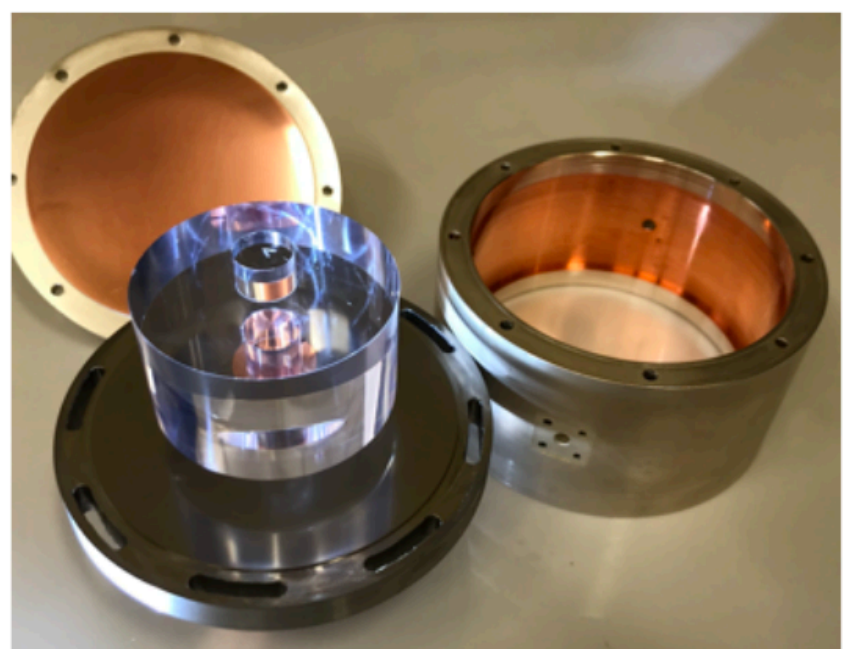


Cryogenic Version Under development < -180 dBc/Hz.

IEEE MICROWAVE AND WIRELESS COMPONENTS LETTERS, VOL. 31, NO. 4, APRIL 2021

Noise Suppression With Cryogenic Resonators

Eugene N. Ivanov^{1b} and Michael E. Tobar^{1b}, *Fellow, IEEE*



IEEE TRANSACTIONS ON ULTRASONICS, FERROELECTRICS, AND FREQUENCY CONTROL, VOL. 56, NO. 2, FEBRUARY 2009

263

Low Phase-Noise Sapphire Crystal Microwave Oscillators: Current Status

Eugene N. Ivanov and Michael E. Tobar, *Senior Member, IEEE*

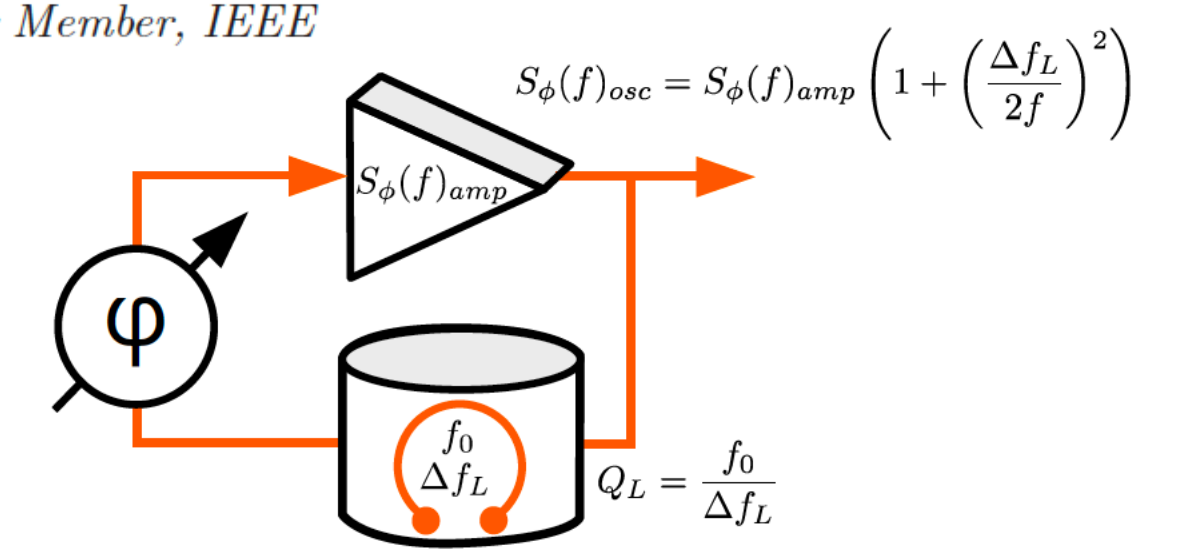


FIG. 4: Schematic of a simple feedback oscillator, with resonator loaded Q-factor Q_L , and amplifier phase noise of $S_\phi(f)_{amp}$. Shown is the simple relation to the oscillator phase noise, $S_\phi(f)_{osc}$.

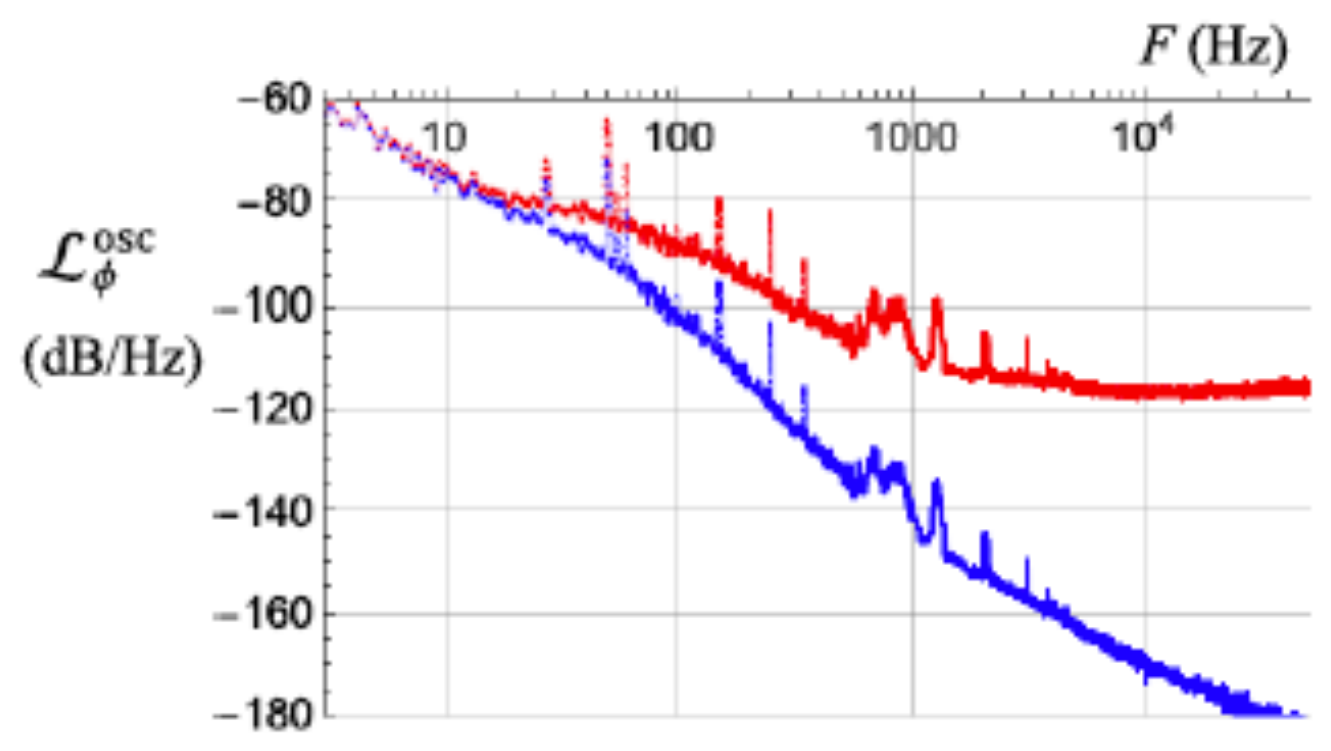


Fig. 7. SSB phase noise spectra of the E8257D at 11.2 GHz: top trace is the measured phase noise of the incident signal; bottom trace is the inferred phase noise of the transmitted signal.

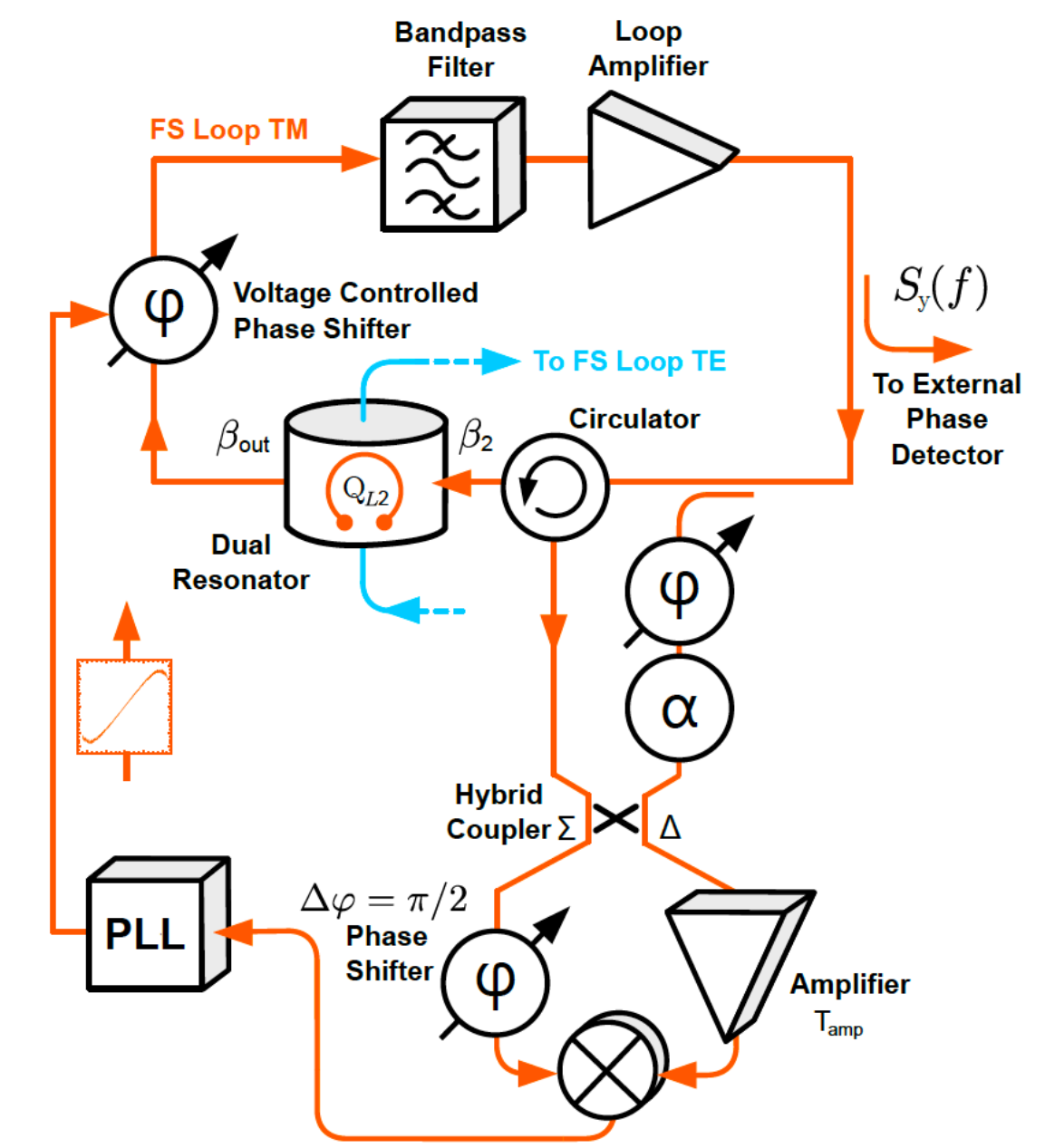
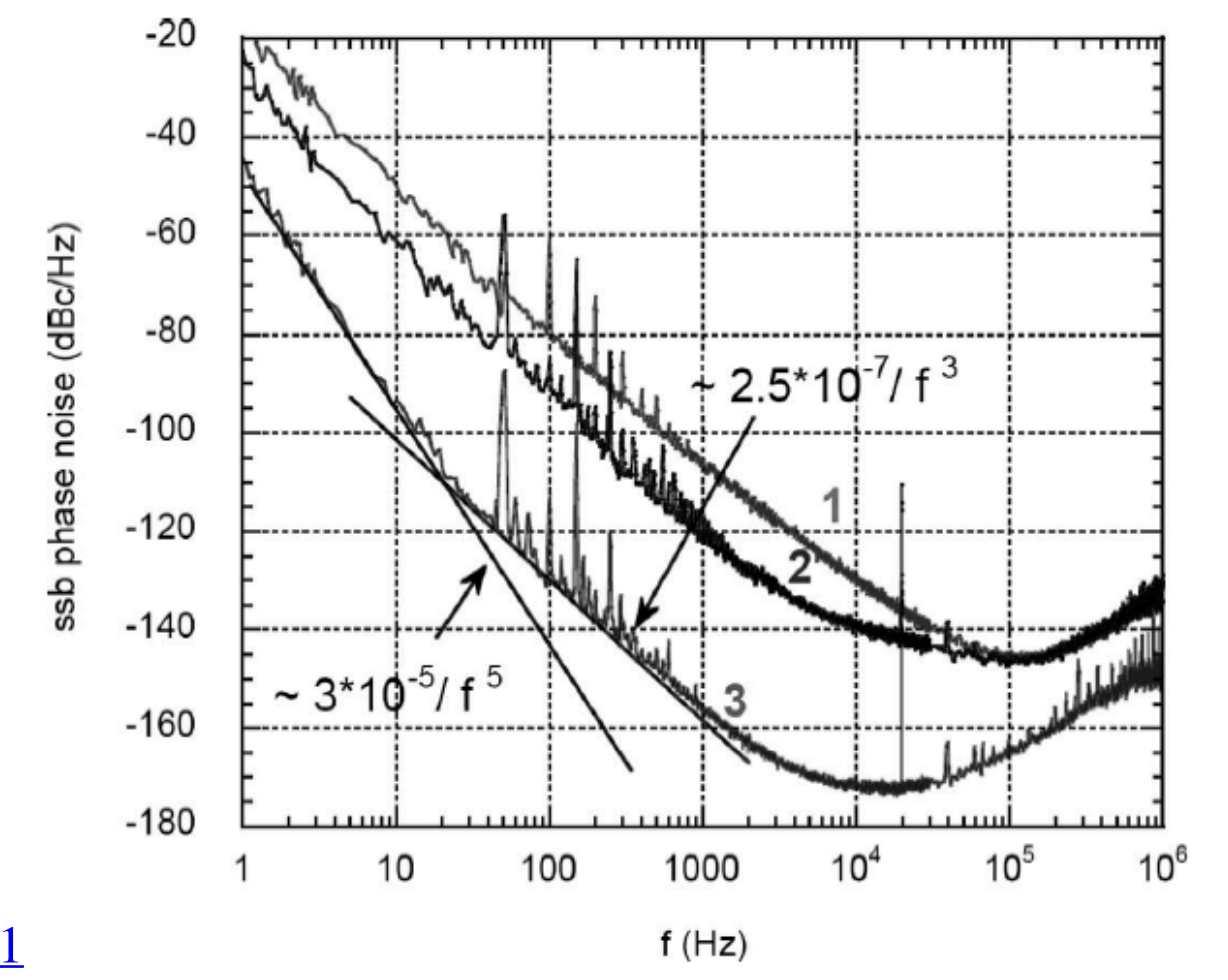


FIG. 5: Schematic of a frequency stabilized (FS) feedback oscillator with interferometric signal processing.



E. N. Ivanov and M. E. Tobar, "Noise Suppression with Cryogenic Resonators," in [IEEE Microwave and Wireless Components Letters](https://doi.org/10.1109/LMWC.2021.3059291), doi: 10.1109/LMWC.2021.3059291, 2021

CA Thomson, BT McAllister, M Goryachev, EN Ivanov, ME Tobar, "Upconversion Loop Oscillator Axion Detection Experiment: A Precision Frequency Interferometric Axion Dark Matter Search with a Cylindrical Microwave Cavity," [Phys. Rev. Lett.](https://arxiv.org/abs/2108.08180), vol. 126, 081803, 2021.

Searching Scalar Dark Matter with Oscillators

Searching for Scalar Dark Matter via Coupling to Fundamental Constants with Photonic, Atomic, and Mechanical Oscillators

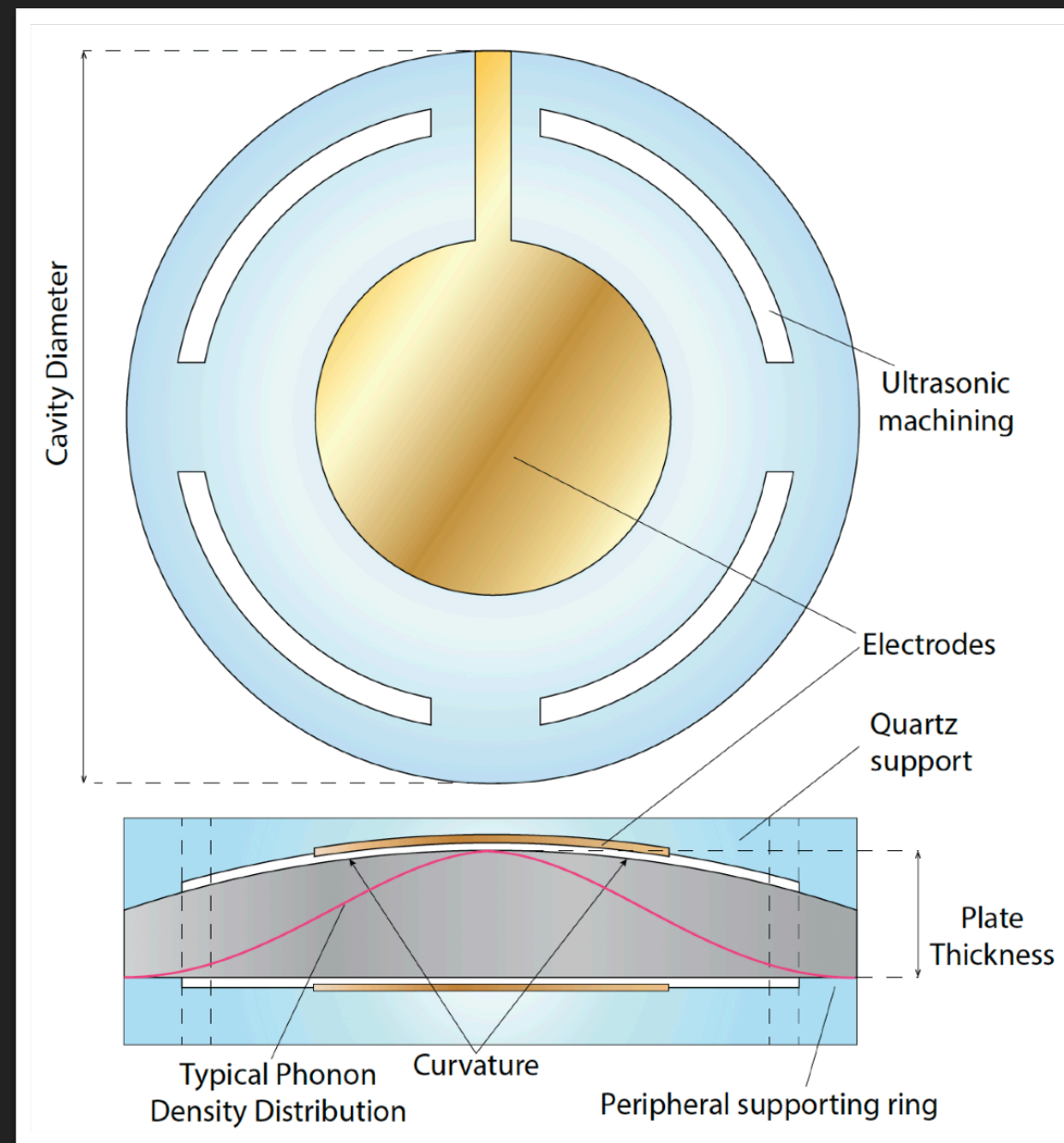
William M. Campbell, Ben T. McAllister, Maxim Goryachev, Eugene N. Ivanov, and Michael E. Tobar
Phys. Rev. Lett. **126**, 071301 – Published 18 February 2021

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We present a way to search for light scalar dark matter (DM), seeking to exploit putative coupling between dark matter scalar fields and fundamental constants, by searching for frequency modulations in direct comparisons between frequency stable oscillators. Specifically we compare a cryogenic sapphire oscillator (CSO), hydrogen maser (HM) atomic oscillator, and a bulk acoustic wave quartz oscillator (OCXO). This work includes the first calculation of the dependence of acoustic oscillators on variations of the fundamental constants, and demonstration that they can be a sensitive tool for scalar DM experiments. Results are presented based on 16 days of data in comparisons between the HM and OCXO, and 2 days of comparison between the OCXO and CSO. No evidence of oscillating fundamental constants consistent with a coupling to scalar dark matter is found, and instead limits on the strength of these couplings as a function of the dark matter mass are determined. We constrain the dimensionless coupling constant d_e and combination $|d_{m_e} - d_g|$ across the mass band $4.4 \times 10^{-19} \lesssim m_\phi \lesssim 6.8 \times 10^{-14} \text{ eV } c^{-2}$, with most sensitive limits $d_e \gtrsim 1.59 \times 10^{-1}$, $|d_{m_e} - d_g| \gtrsim 6.97 \times 10^{-1}$. Notably, these limits do not rely on maximum reach analysis (MRA), instead employing the more general coefficient separation technique. This experiment paves the way for future, highly sensitive experiments based on state-of-the-art acoustic oscillators, and we show that these limits can be competitive with the best current MRA-based exclusion limits.

Bulk Acoustic Wave (BAW) Oscillator Fundamental Constant Dependence

Quartz Oscillator Limits the Experiment Stability: 10^{-13} to 10^{-16} possible



$$f = \frac{nv}{L}$$

vis the speed of sound in the material
 L resonator length parameter
 n is a constant.

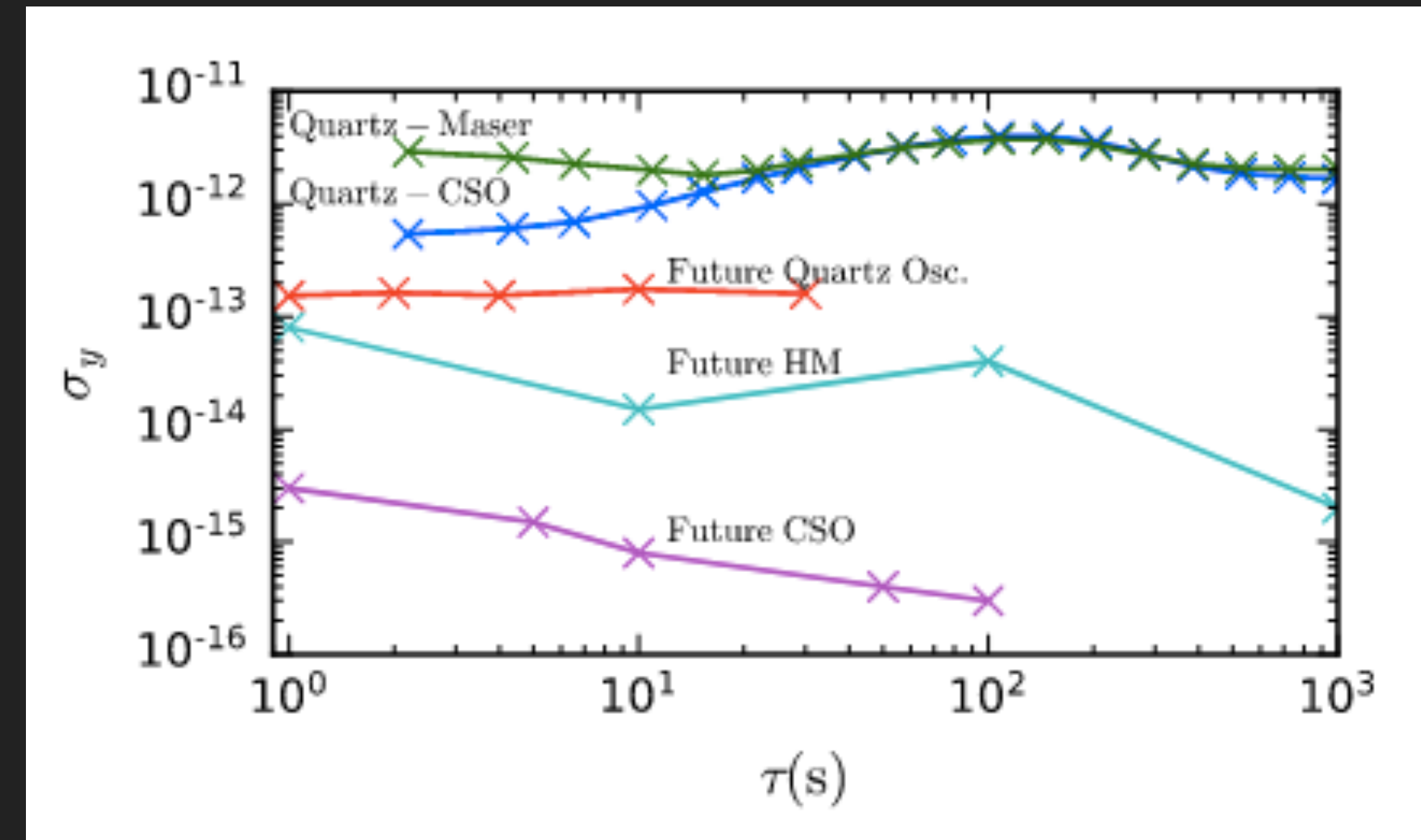
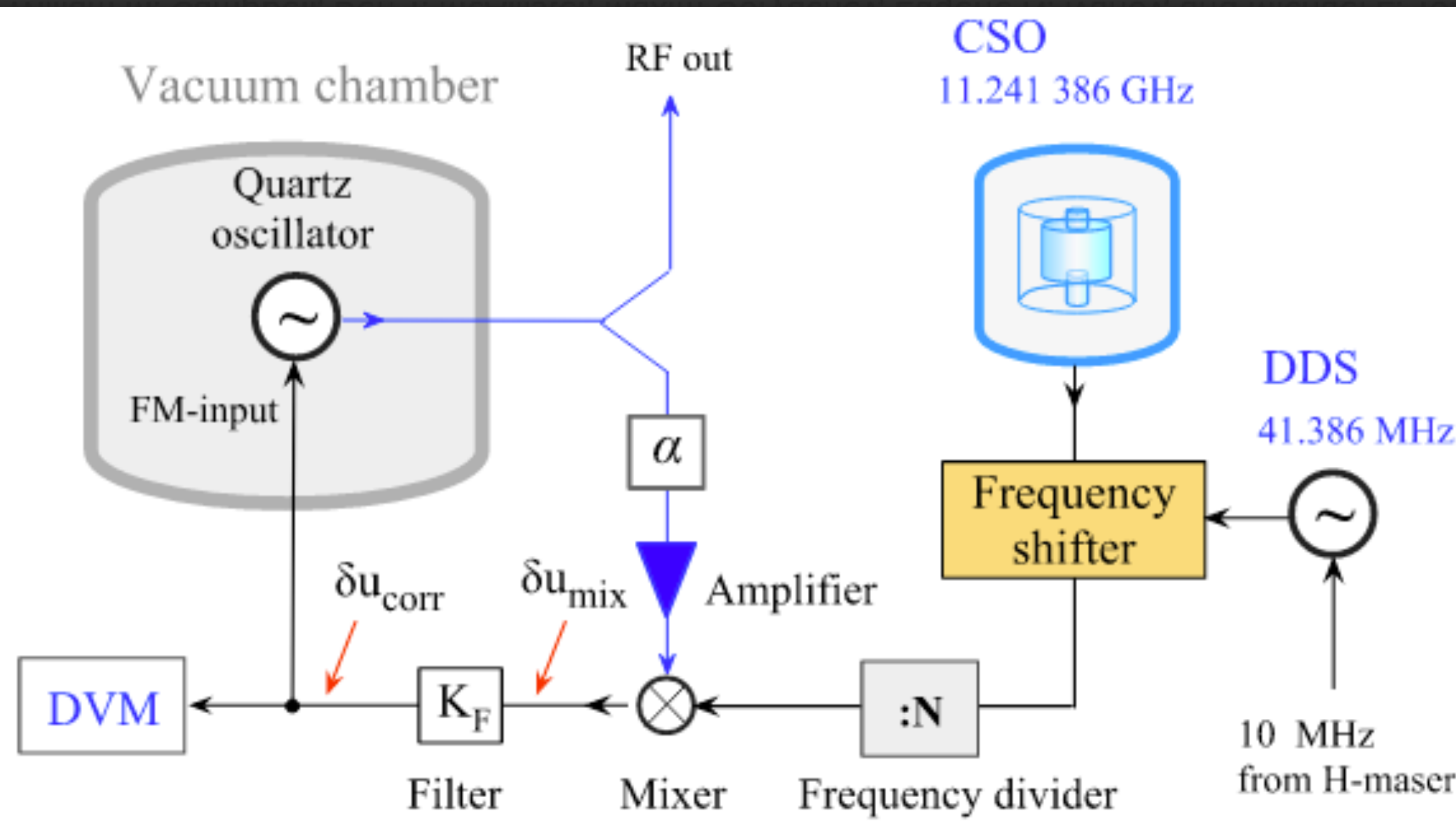
$$f \propto \sqrt{\frac{KL}{m}} \propto \sqrt{\frac{\alpha^4 m_e^3}{m_p}} \propto m_e \alpha^2 \sqrt{\frac{m_e}{m_p}}$$

Scalar Dark Matter

Searching for Scalar Dark Matter via Coupling to Fundamental Constants with Photonic, Atomic, and Mechanical Oscillators

William M. Campbell, Ben T. McAllister, Maxim Goryachev, Eugene N. Ivanov, and Michael E. Tobar
 Phys. Rev. Lett. **126**, 071301 – Published 18 February 2021

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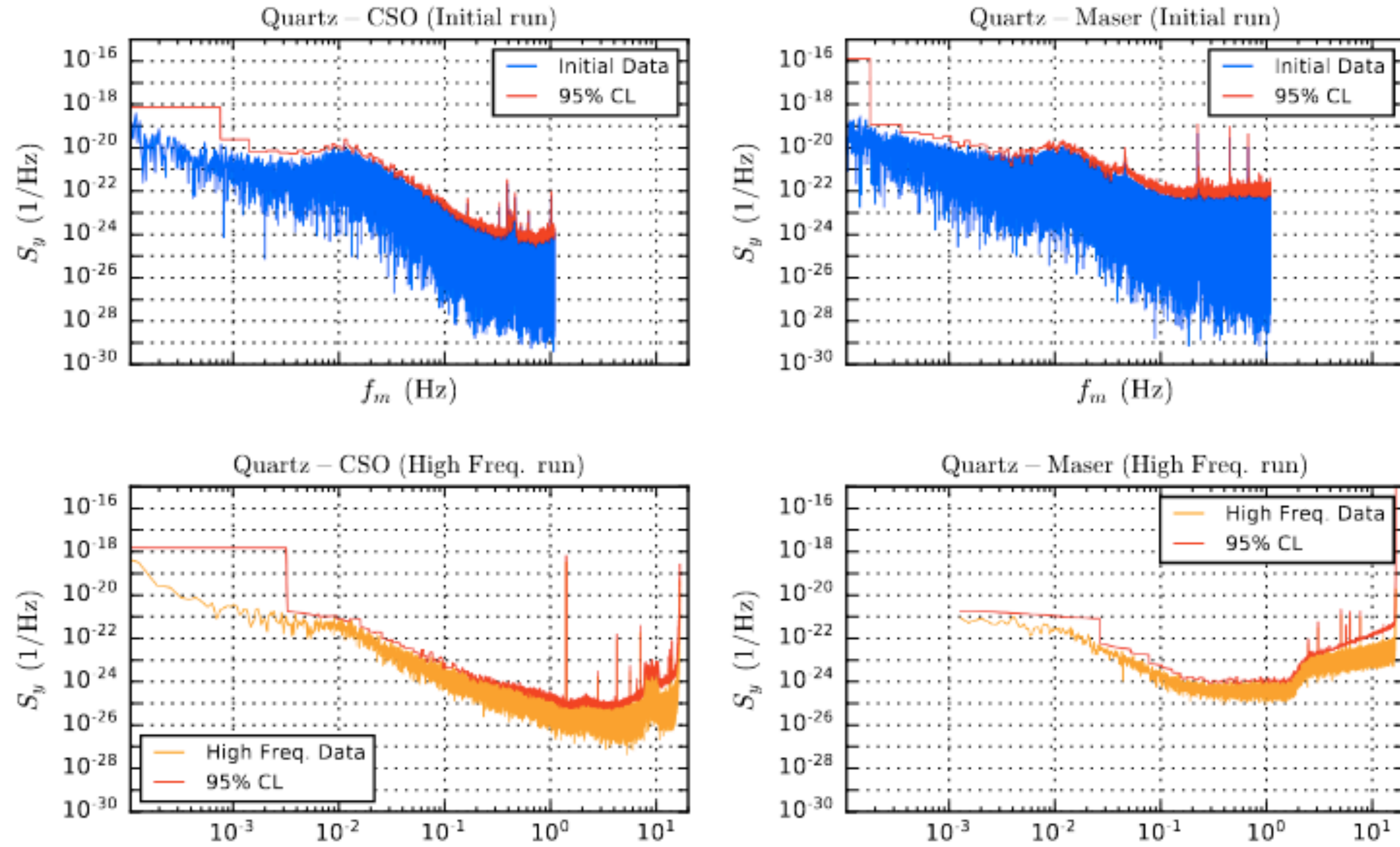
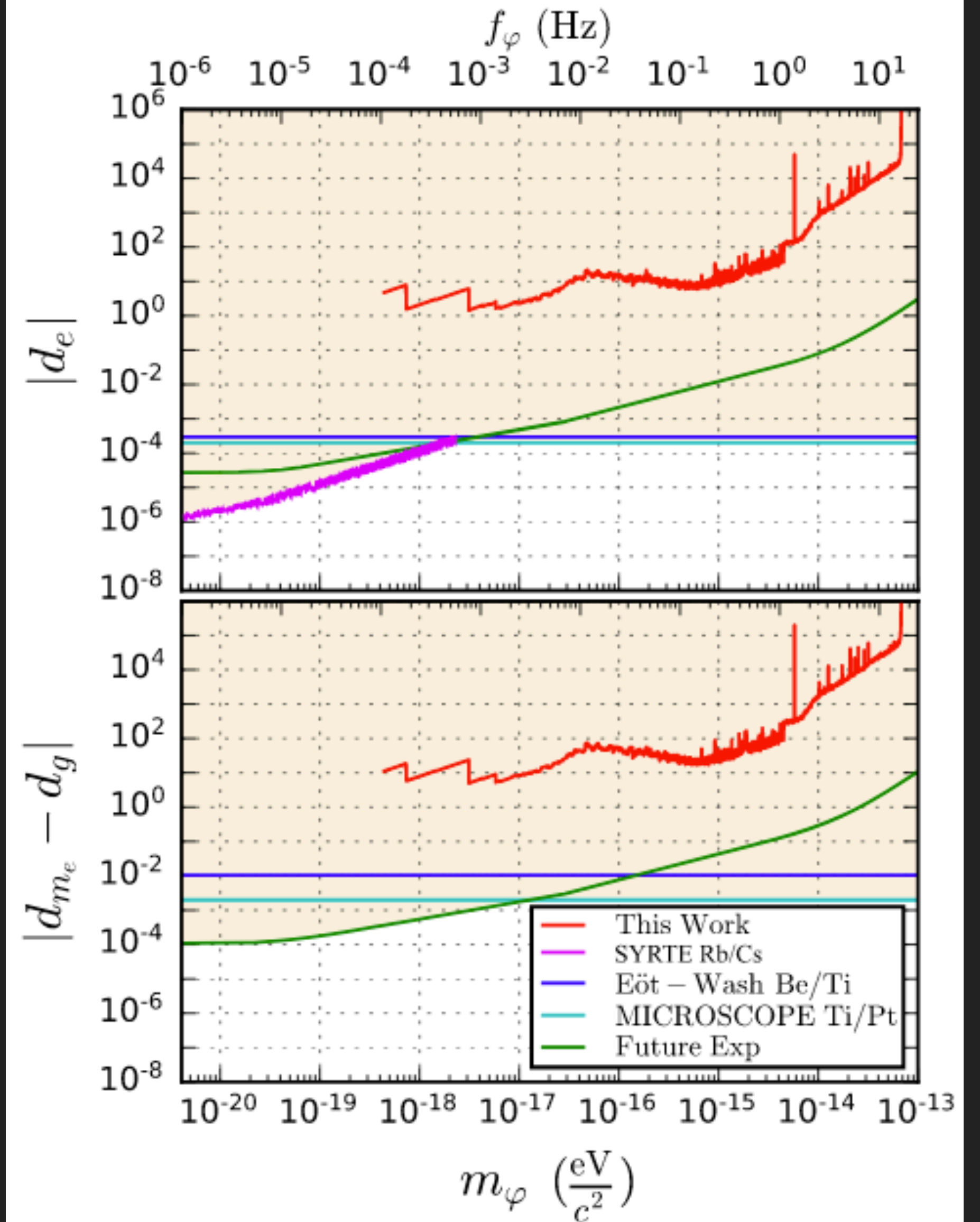
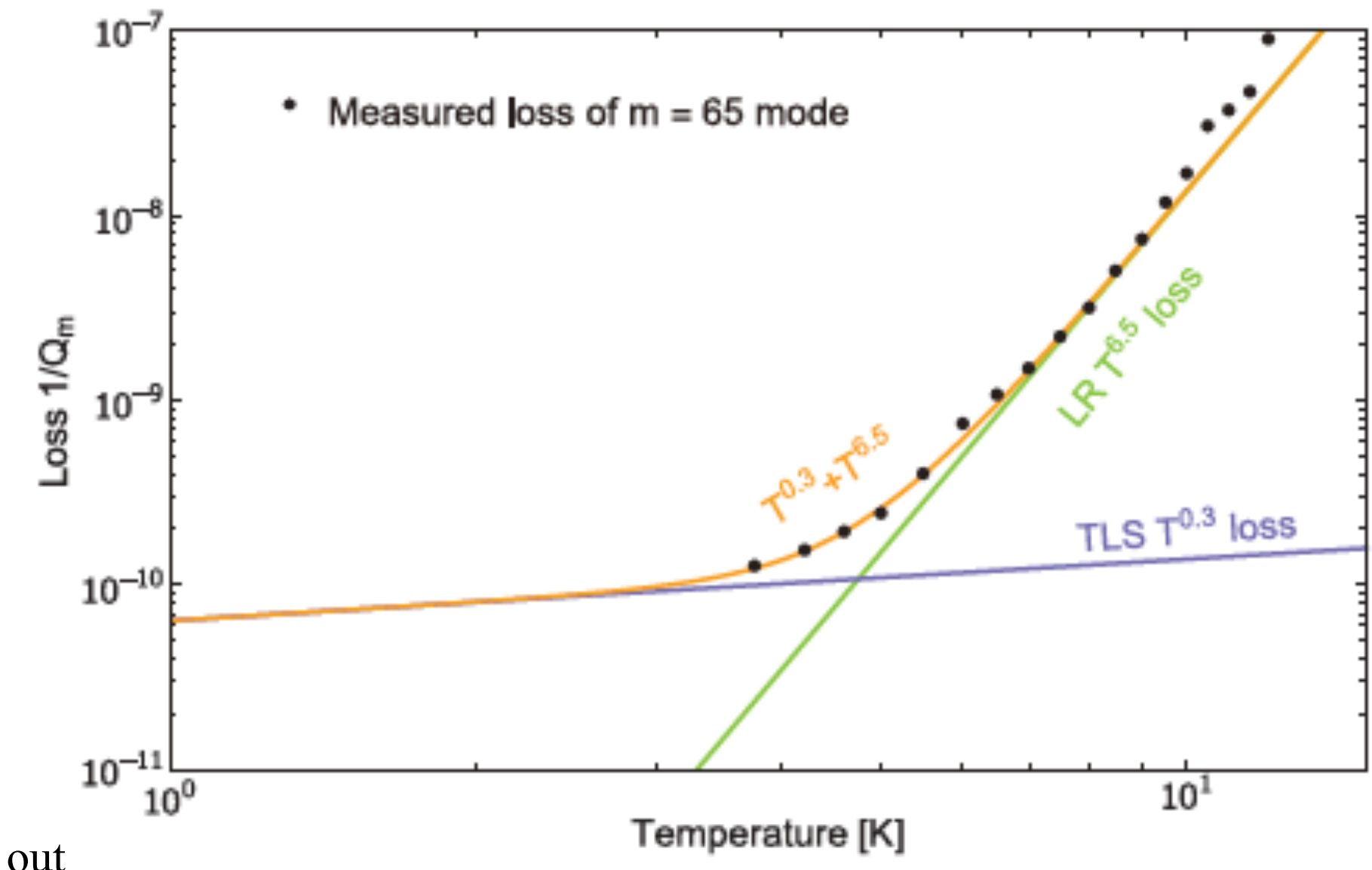
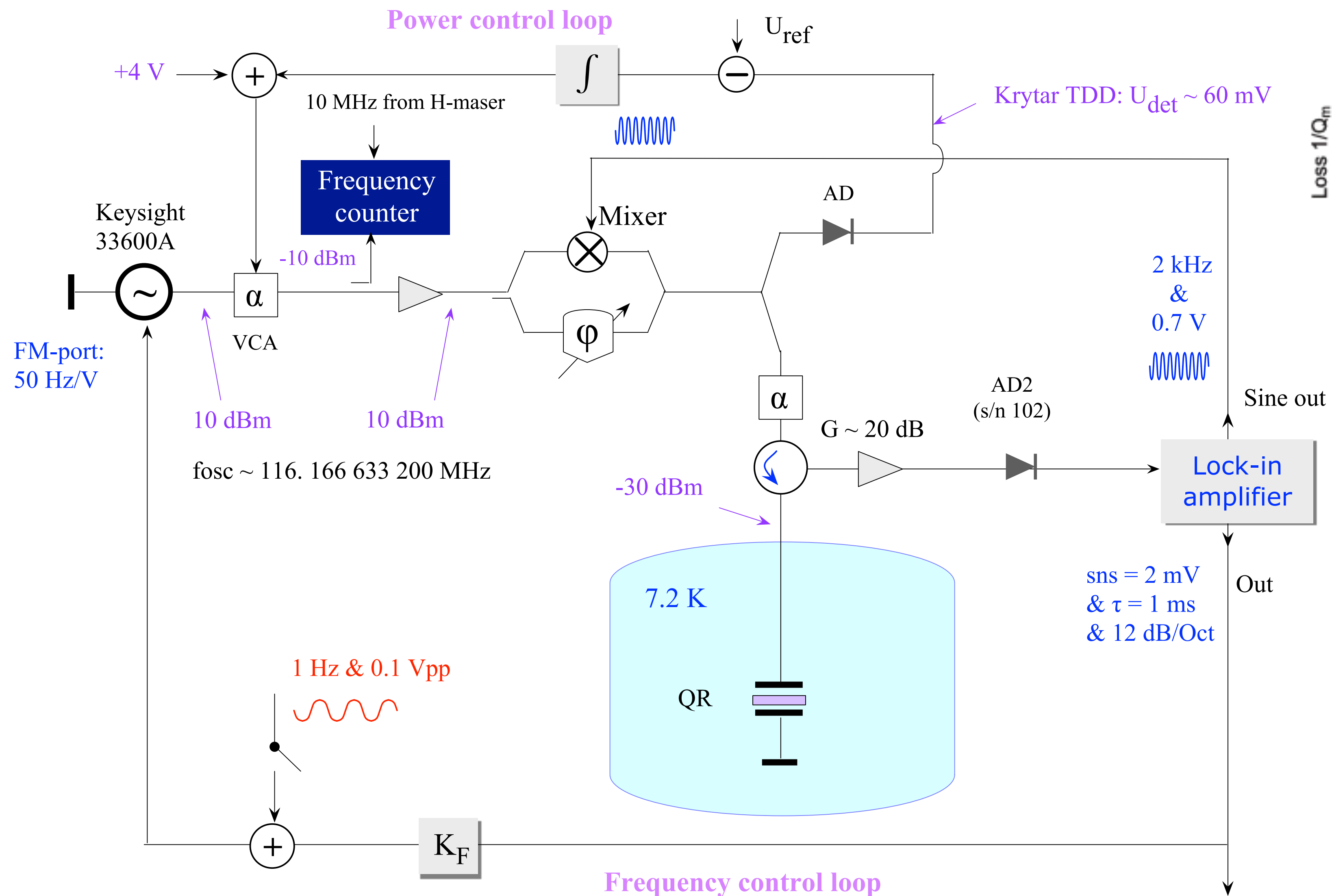


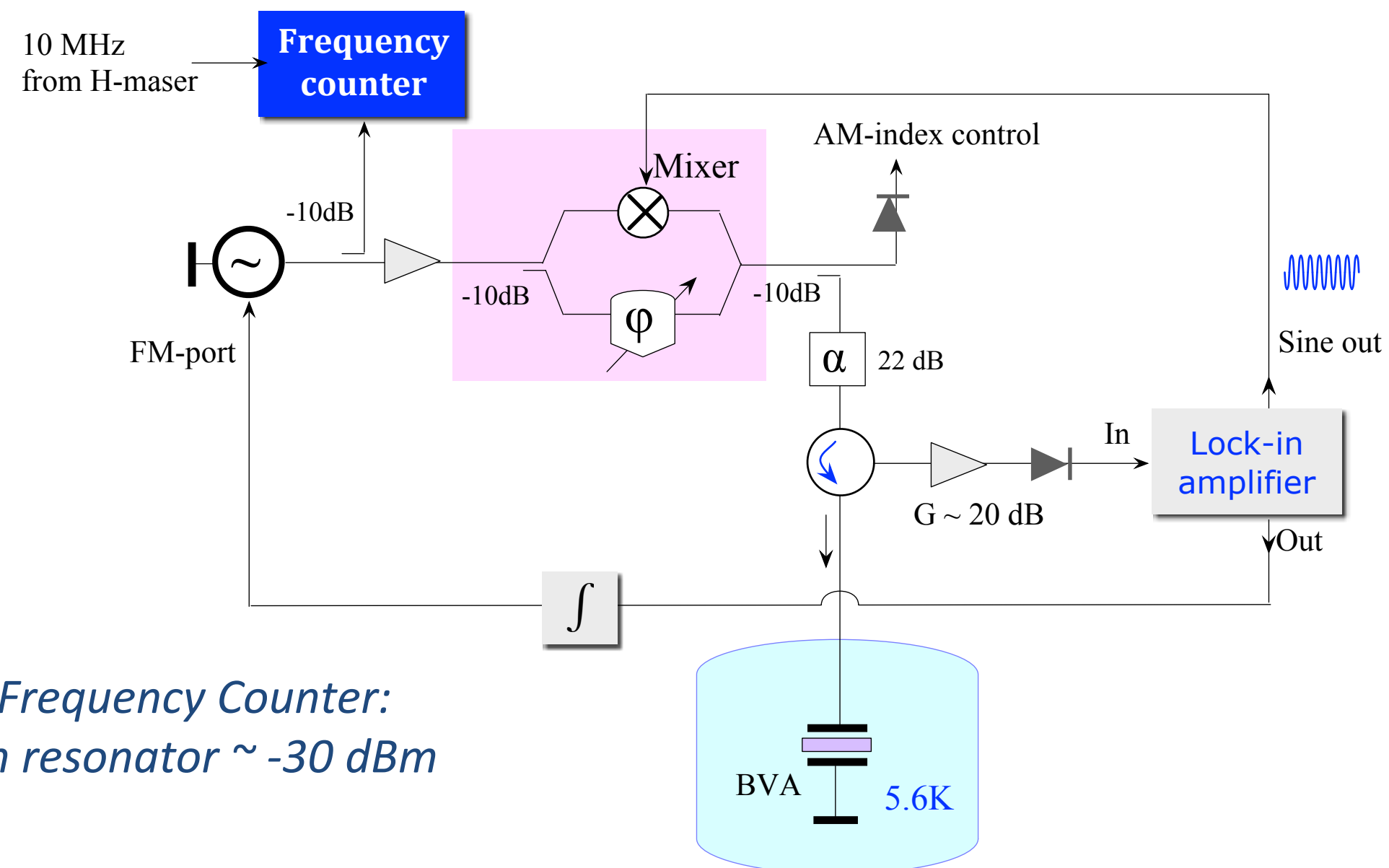
FIG. 6. PSD of frequency noise for all initial and later runs are shown by the blue and orange traces respectively. Also shown in Red is the excluded power to 95% confidence.



Next Generation Experiment? Cryogenic BAW (4K) -> $Q \sim 10^{10}$ 4 orders better than room temp



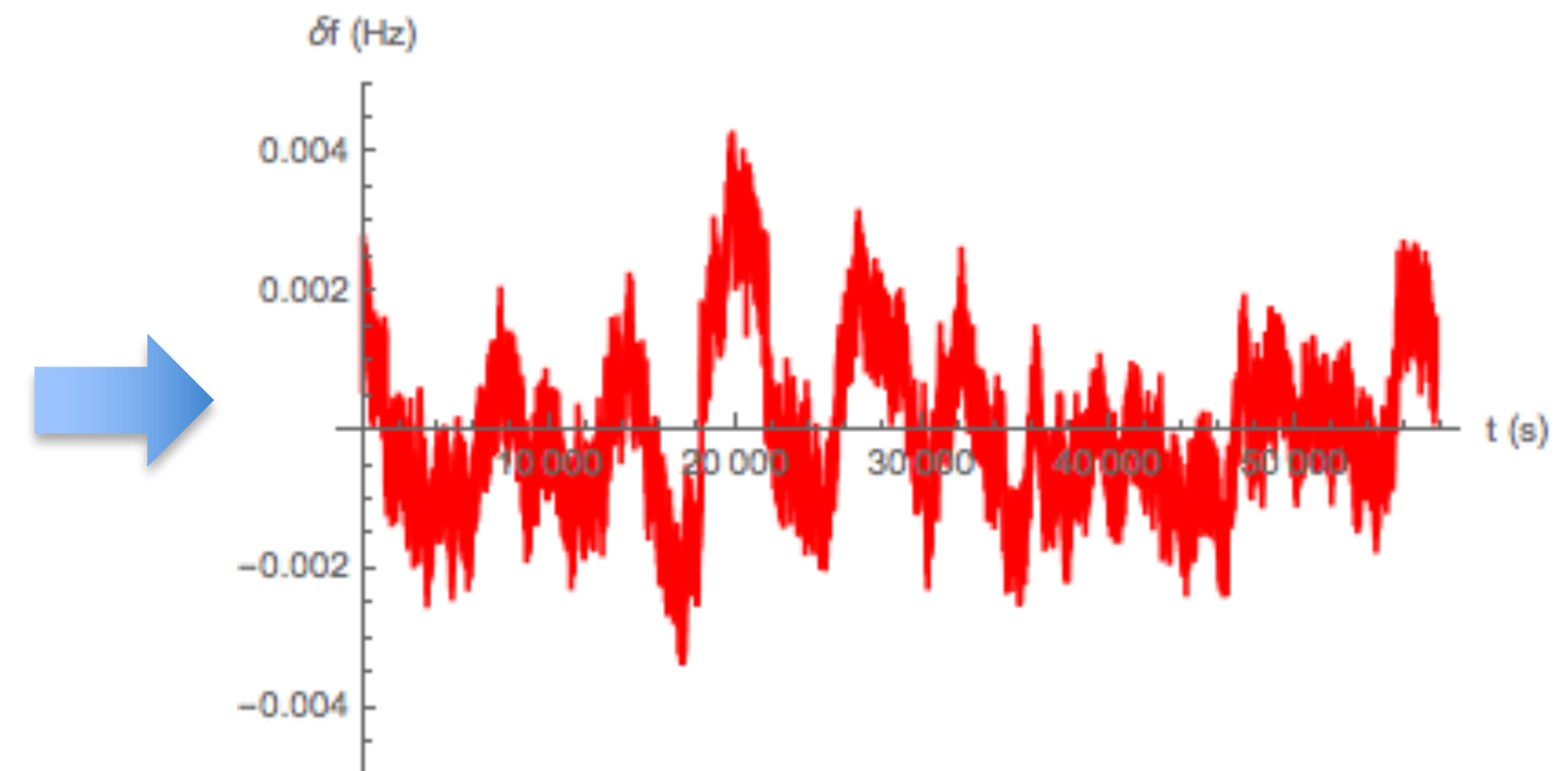
Frequency Stability Measurements



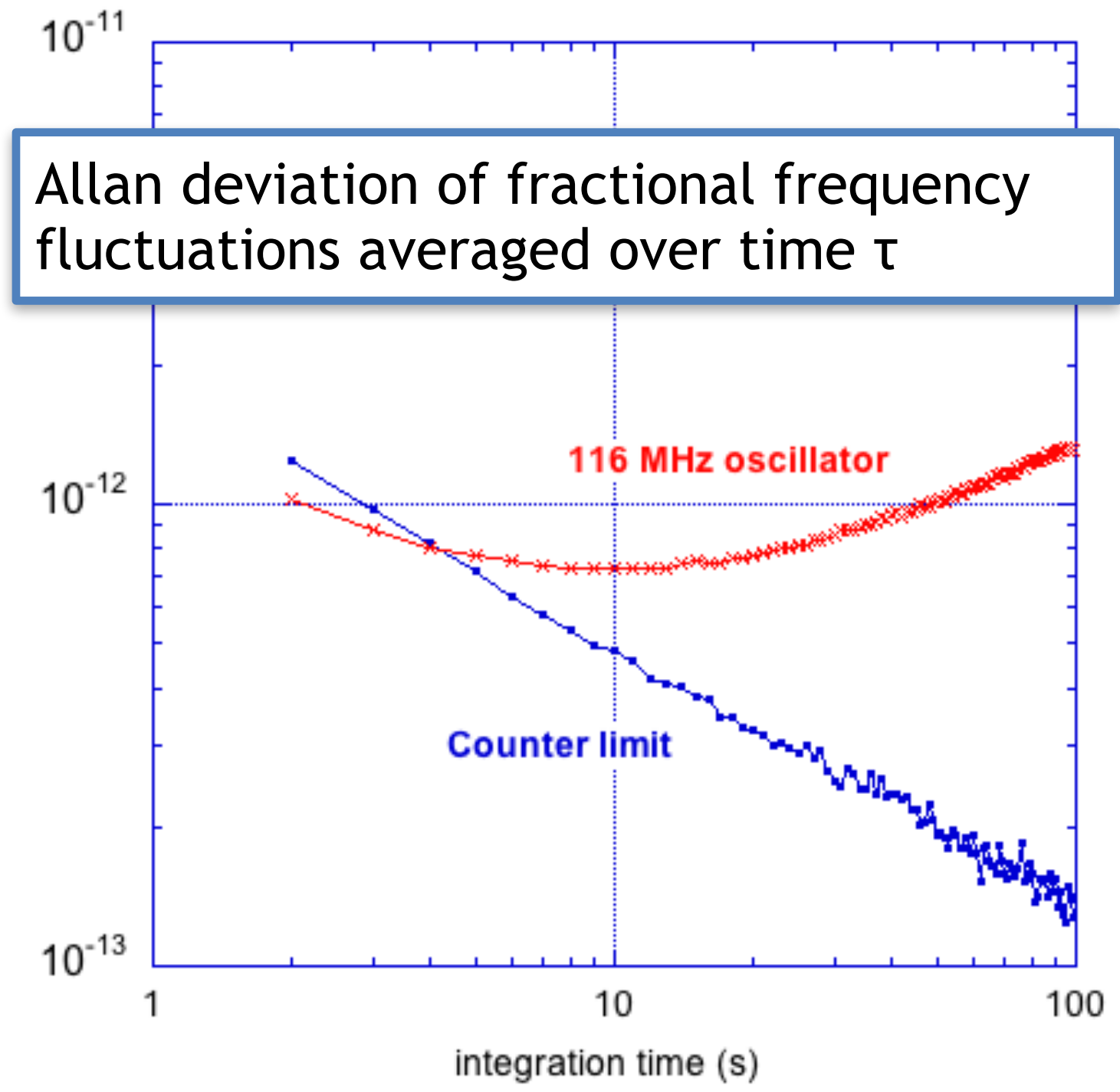
RF oscillator & Frequency Counter:
Power incident on resonator ~ -30 dBm

NB

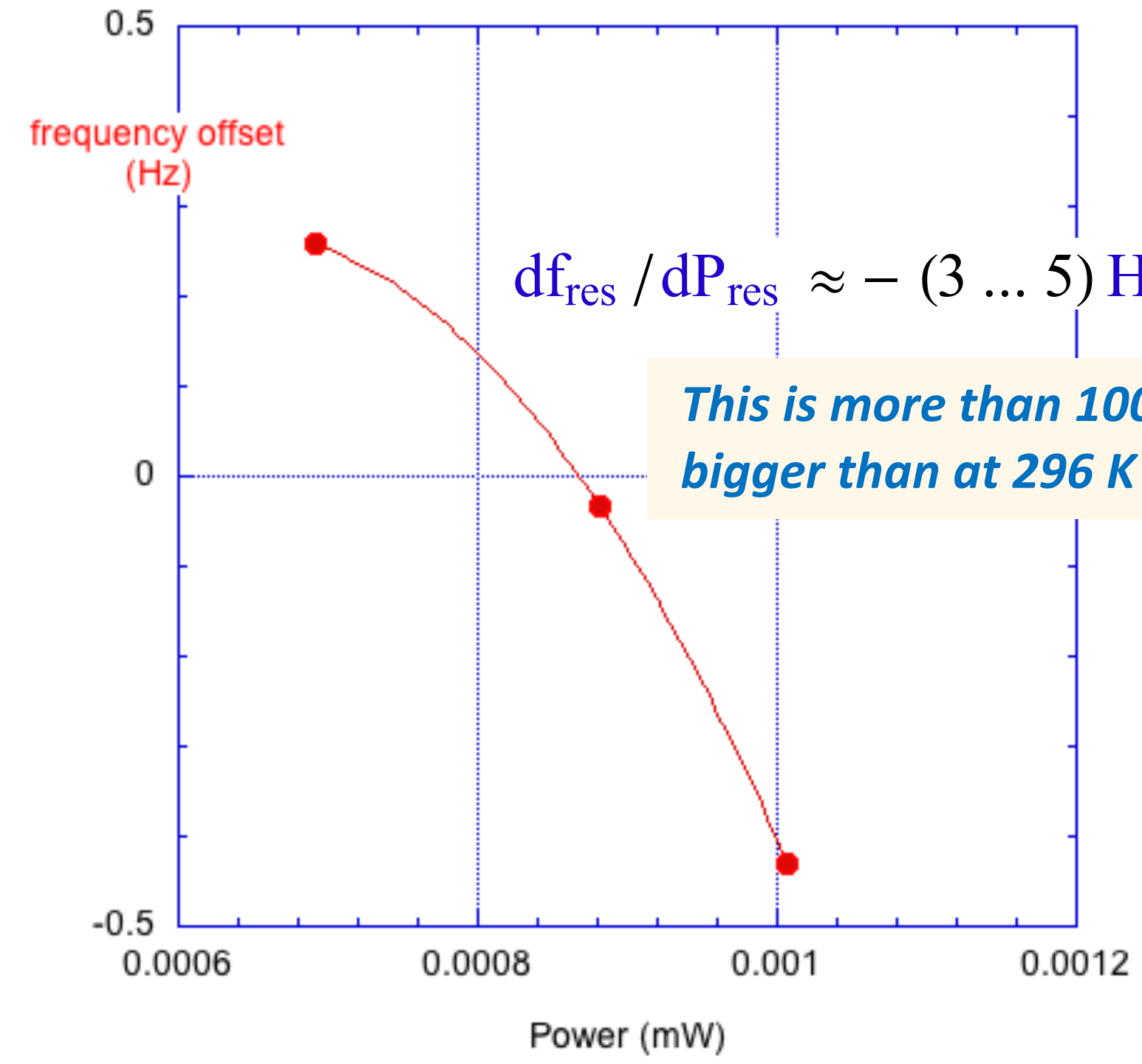
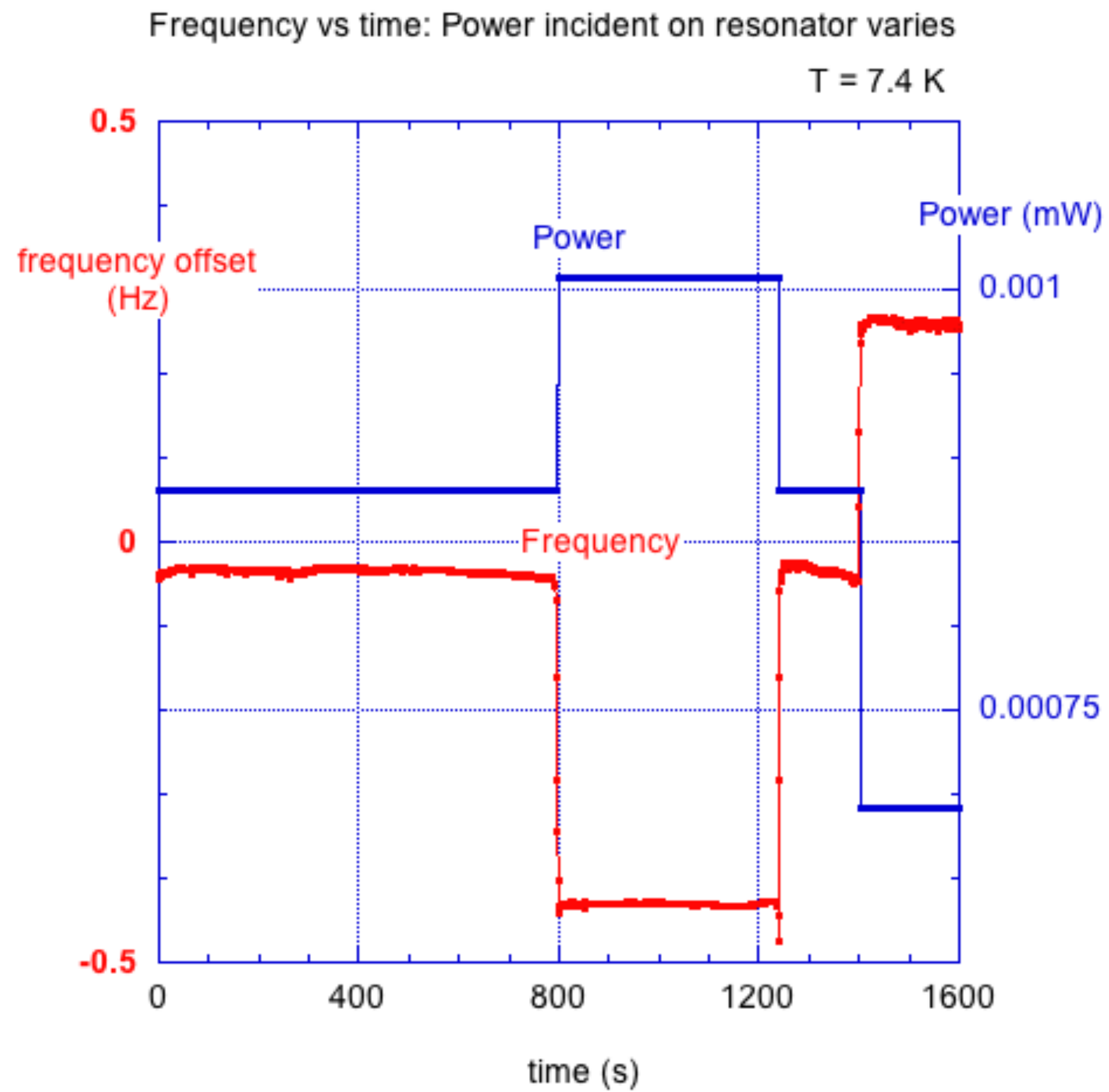
1. There is no noticeable frequency drift despite the use of room temperature detection system.
2. Frequency stability improves by $\sim 10\%$ when LA sensitivity increases from 20 to 5 mV



Allan deviation of fractional frequency fluctuations averaged over time τ



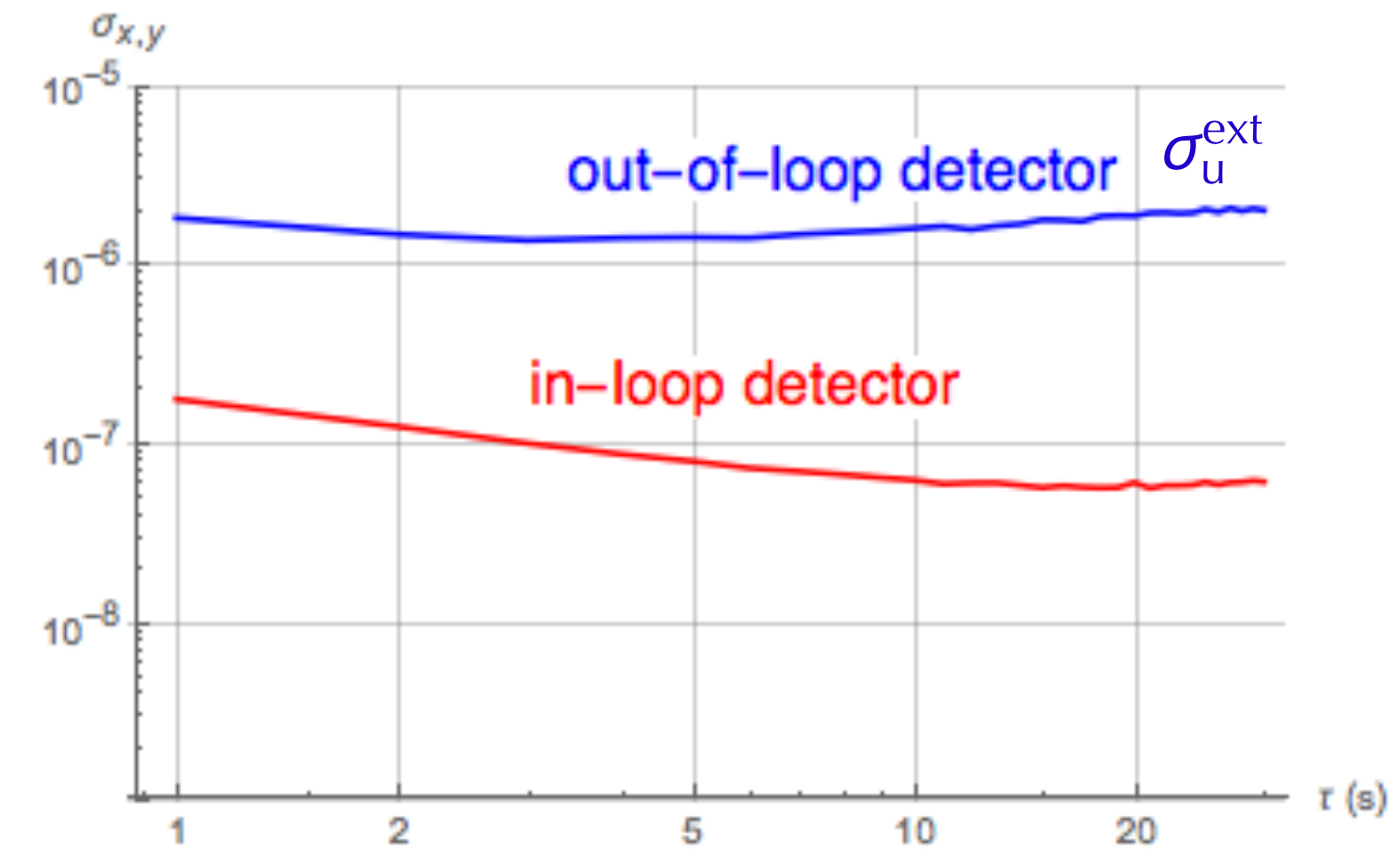
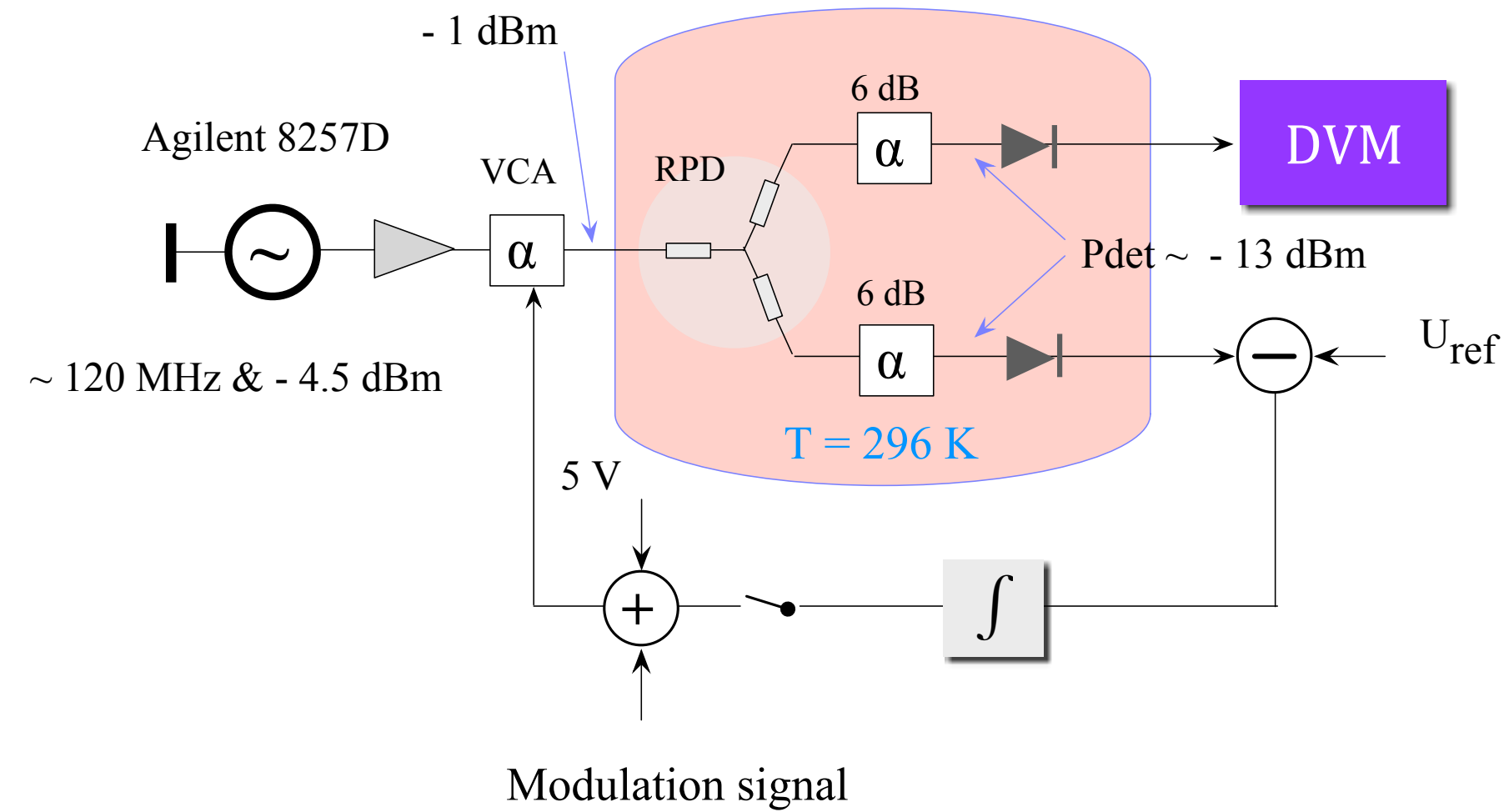
Cryogenic Quartz Oscillator: Power-to-Frequency Conversion: Duffing Nonlinearity



Oscillator frequency vs power

Power incident on cryogenic BVA resonator (blue) and oscillator frequency (red) vs time

Oscillator Frequency Stability due to Power Fluctuations



Oscillator fractional frequency stability due to power fluctuations

$$\sigma_y = \frac{1}{f_{\text{res}}} \frac{df_{\text{res}}}{dP_{\text{res}}} \eta \frac{\sigma_u^{\text{ext}}}{\sqrt{(S_{\text{AD}}^{\text{in-loop}})^2 + (S_{\text{AD}}^{\text{ext}})^2}}$$

$\eta = P_{\text{res}} / P_{\text{det}} \ll 1$

Allan deviation of fractional voltage fluctuations at the output of out-of-loop detector

Power-to-Voltage conversions of amplitude detectors (see next slides)

Oscillator parameters:

$$P_{\text{det}} = -10 \text{ dBm}$$

$$P_{\text{res}} = -30 \text{ dBm}$$

$$df_{\text{res}} / dP_{\text{res}} \approx - (3 \dots 5) \text{ Hz}/\mu\text{W}$$




$$du / dP_{\text{det}} \approx 1000 \text{ mV}/\text{mW}$$

$$\sigma_u^{\text{ext}} (1 \dots 30 \text{ s}) \approx 2 \times 10^{-6}$$



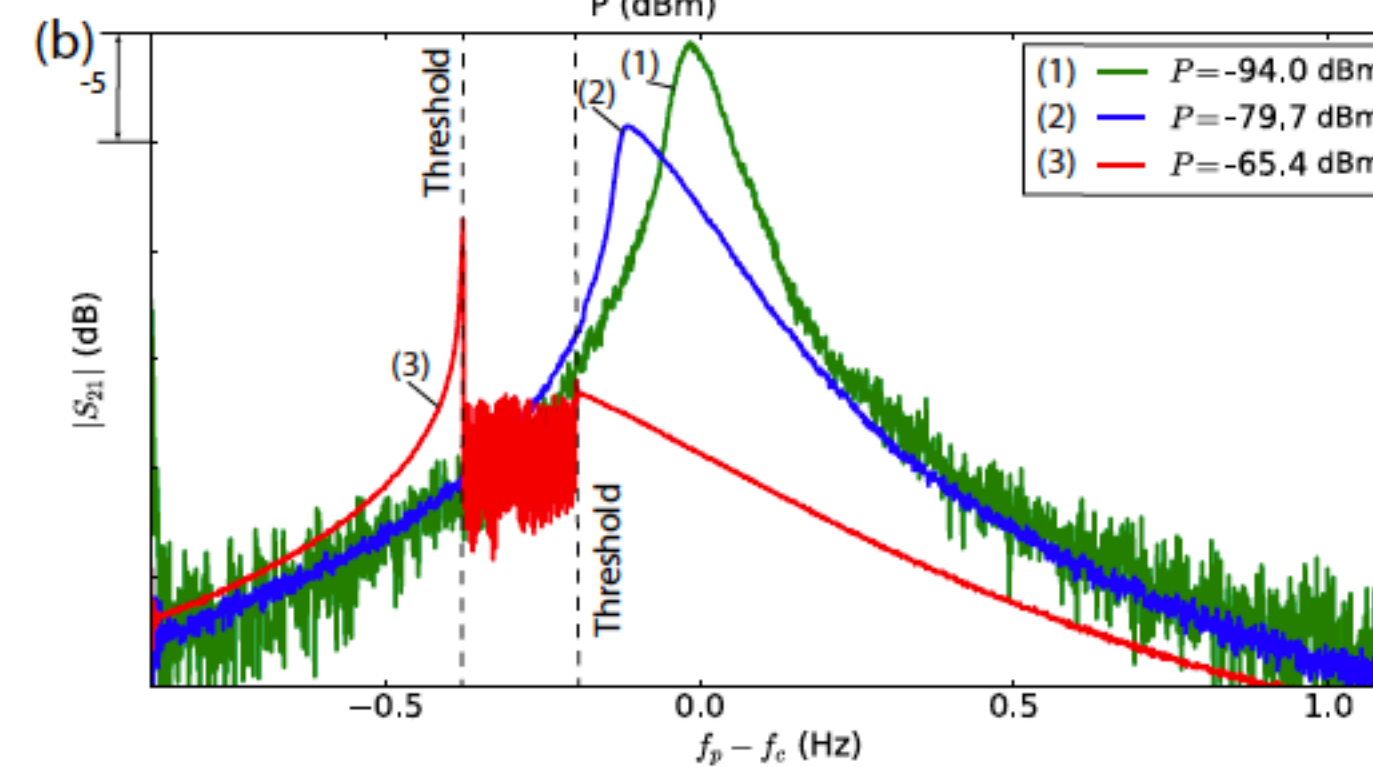
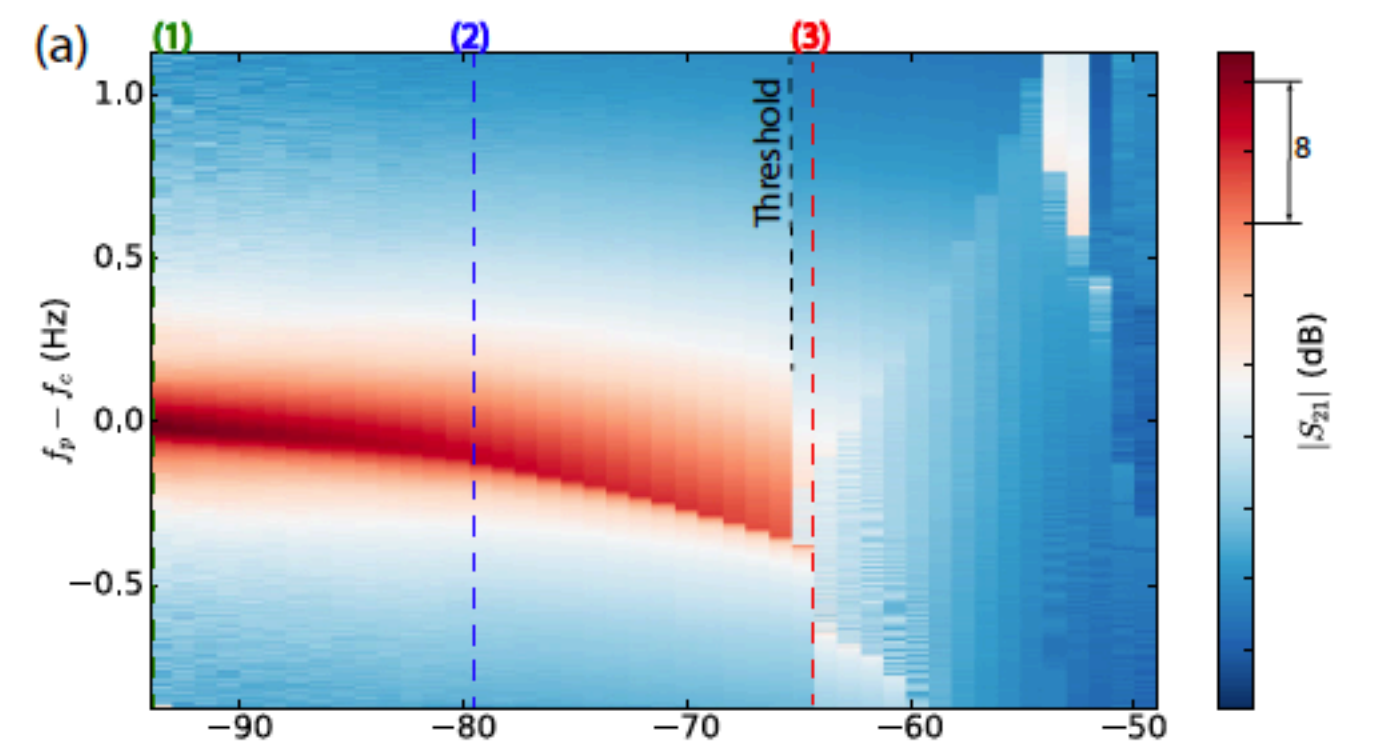
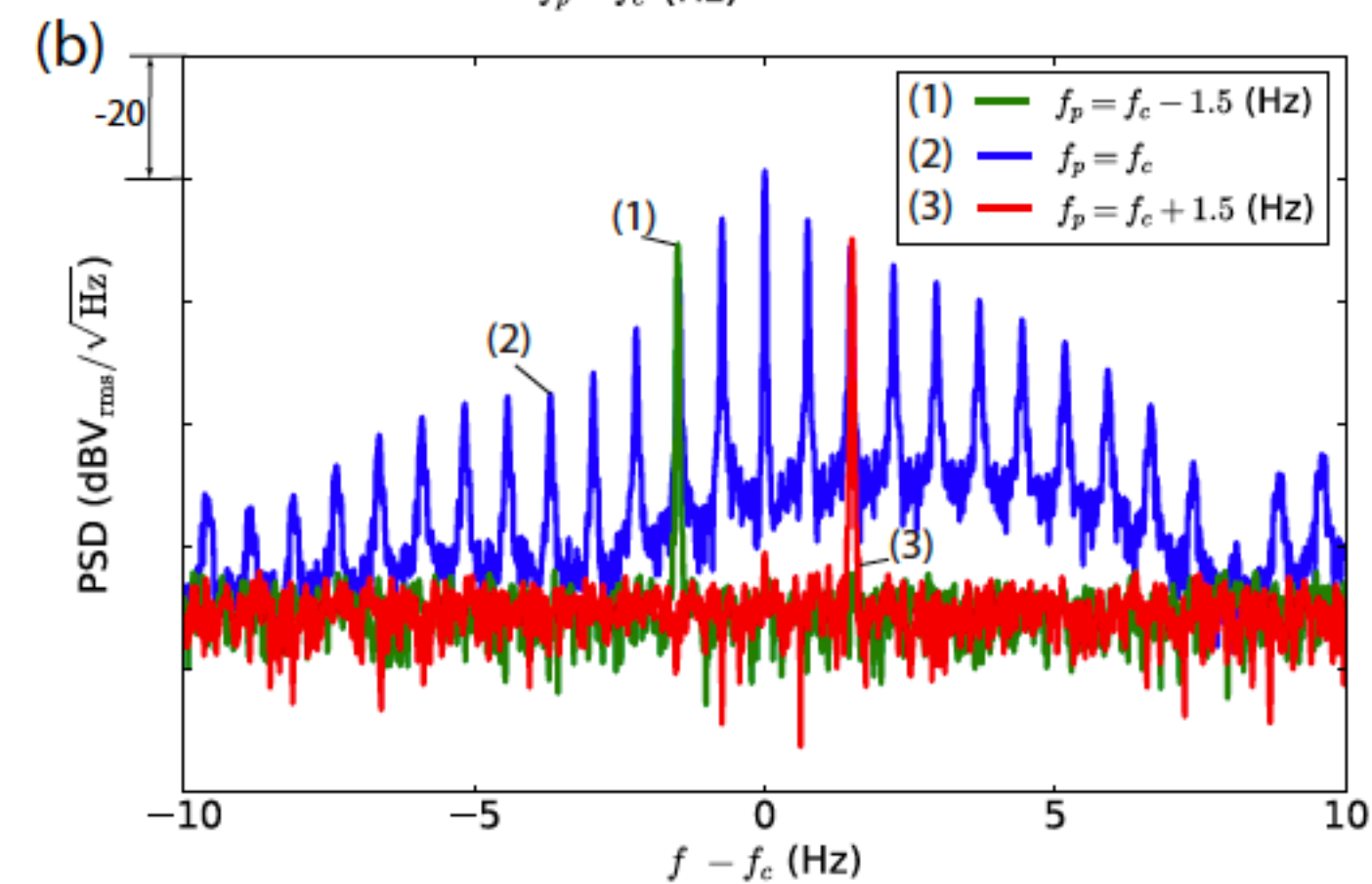
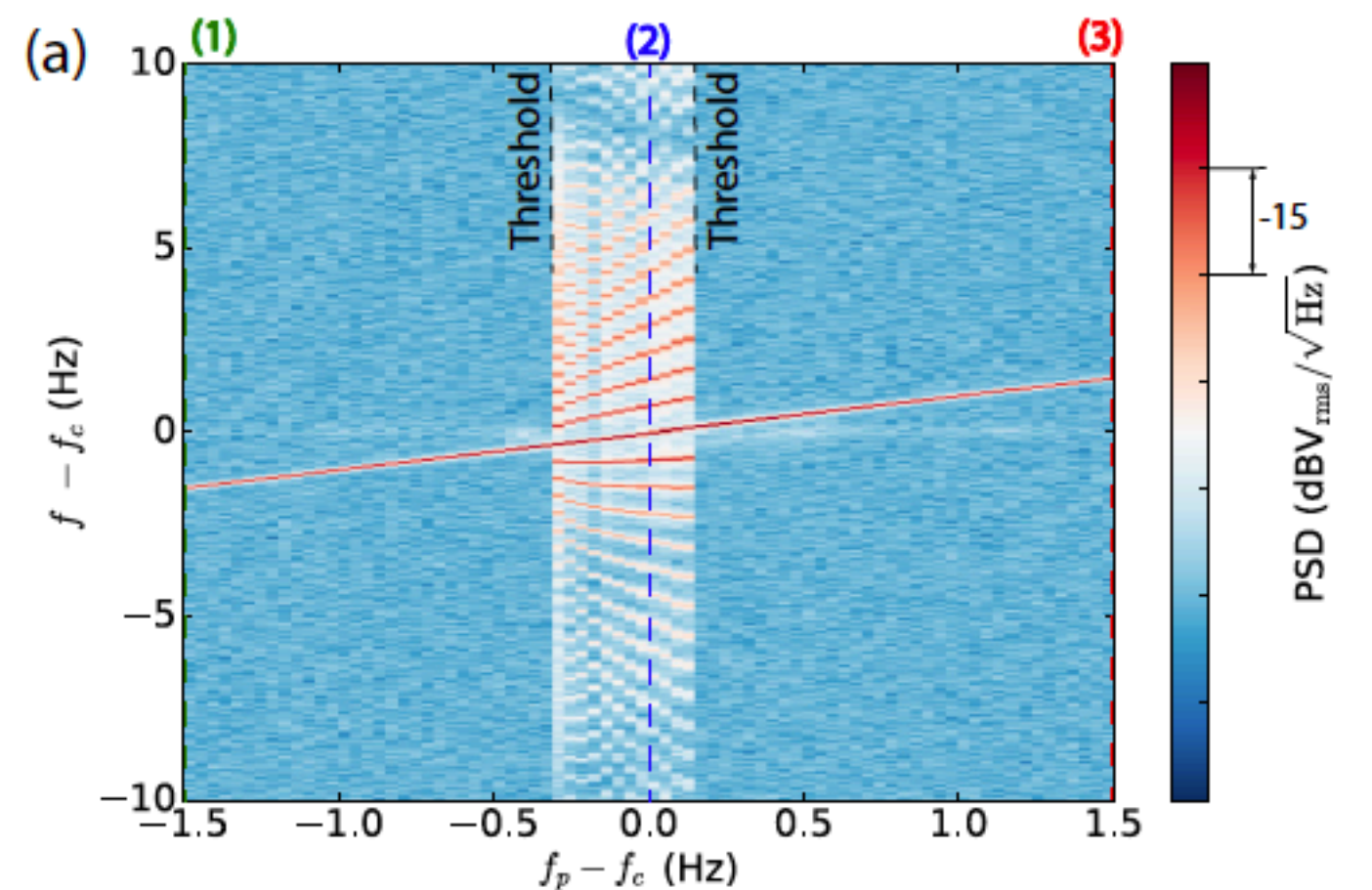
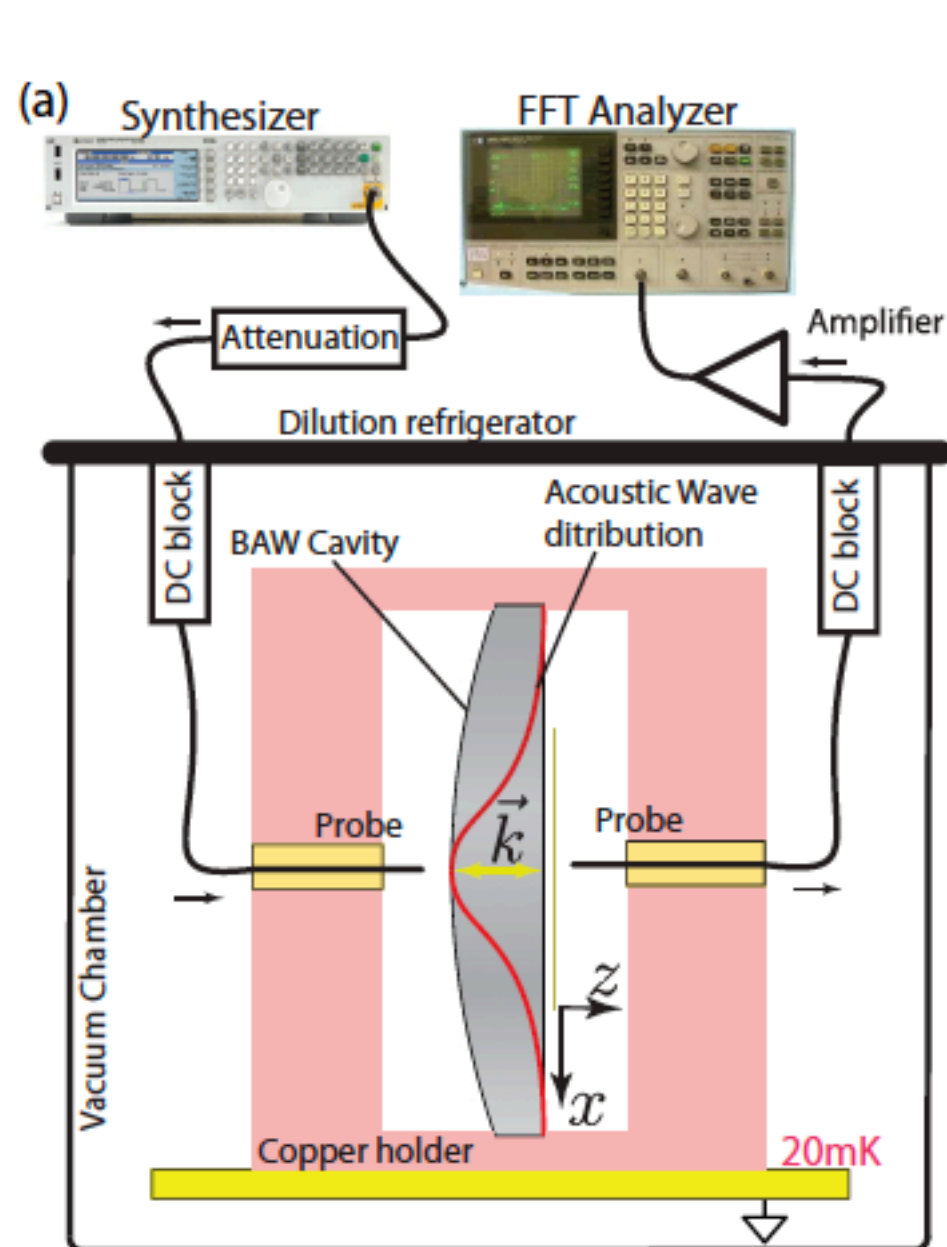
$$\sigma_y (1 \dots 30 \text{ s}) \approx 5 \dots 7 \times 10^{-13}$$

Generation of ultralow power phononic combs

Maxim Goryachev ^{1,*}, Serge Galliou ², and Michael E. Tobar ¹




¹ARC Centre of Excellence for Engineered Quantum Systems, University of Western Australia, 35 Stirling Highway, Crawley, WA 6009, Australia

²FEMTO-ST Institute, Université Bourgogne Franche-Comté, Centre National de la Recherche Scientifique, ENSMM, 26 Rue de l'Épitaphe, 25000 Besançon, France



Article

Inducing Strong Non-Linearities in a Phonon Trapping Quartz Bulk Acoustic Wave Resonator Coupled to a Superconducting Quantum Interference Device

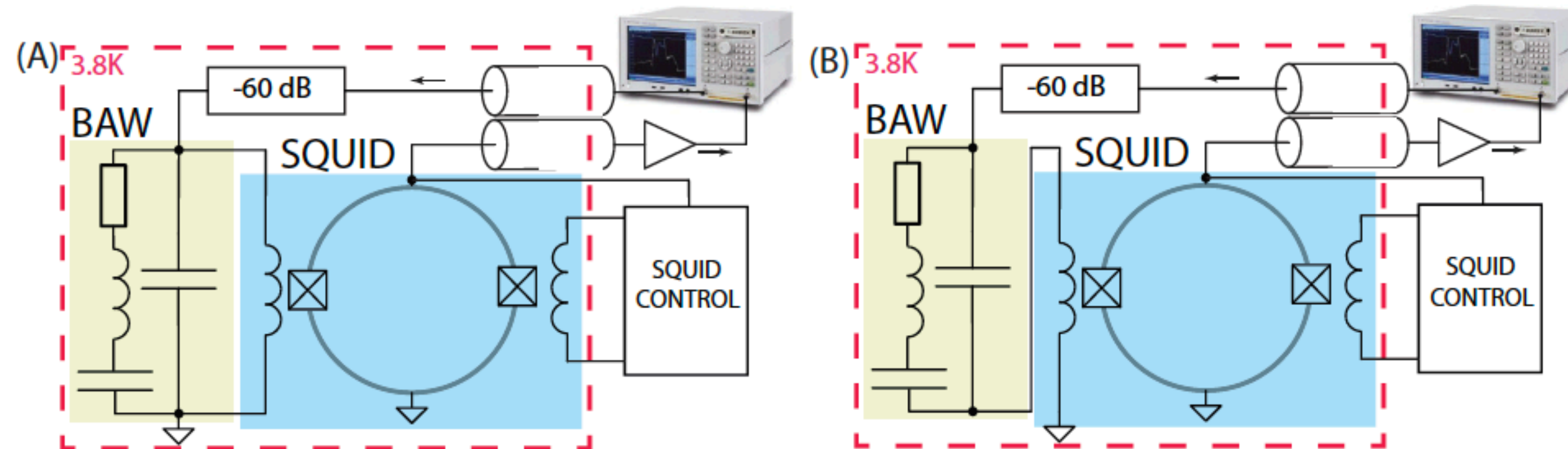
Maxim Goryachev ¹, Eugene N. Ivanov ¹ , Serge Galliou ²  and Michael E. Tobar ^{1,*} 

¹ ARC Centre of Excellence for Engineered Quantum Systems, University of Western Australia, 35 Stirling Highway, Crawley WA 6009, Australia; maxim.goryachev@uwa.edu.au (M.G.); eugene.ivanov@uwa.edu.au (E.N.I.)

² FEMTO-ST Institute, CNRS, Univ. Bourgogne Franche Comte, ENSMM, 26 Chemin de l'Épitaphe, 25000 Besançon, France; serge.galliou@ens2m.fr

* Correspondence: michael.tobar@uwa.edu.au

Received: 22 March 2018; Accepted: 4 April 2018; Published: 11 April 2018

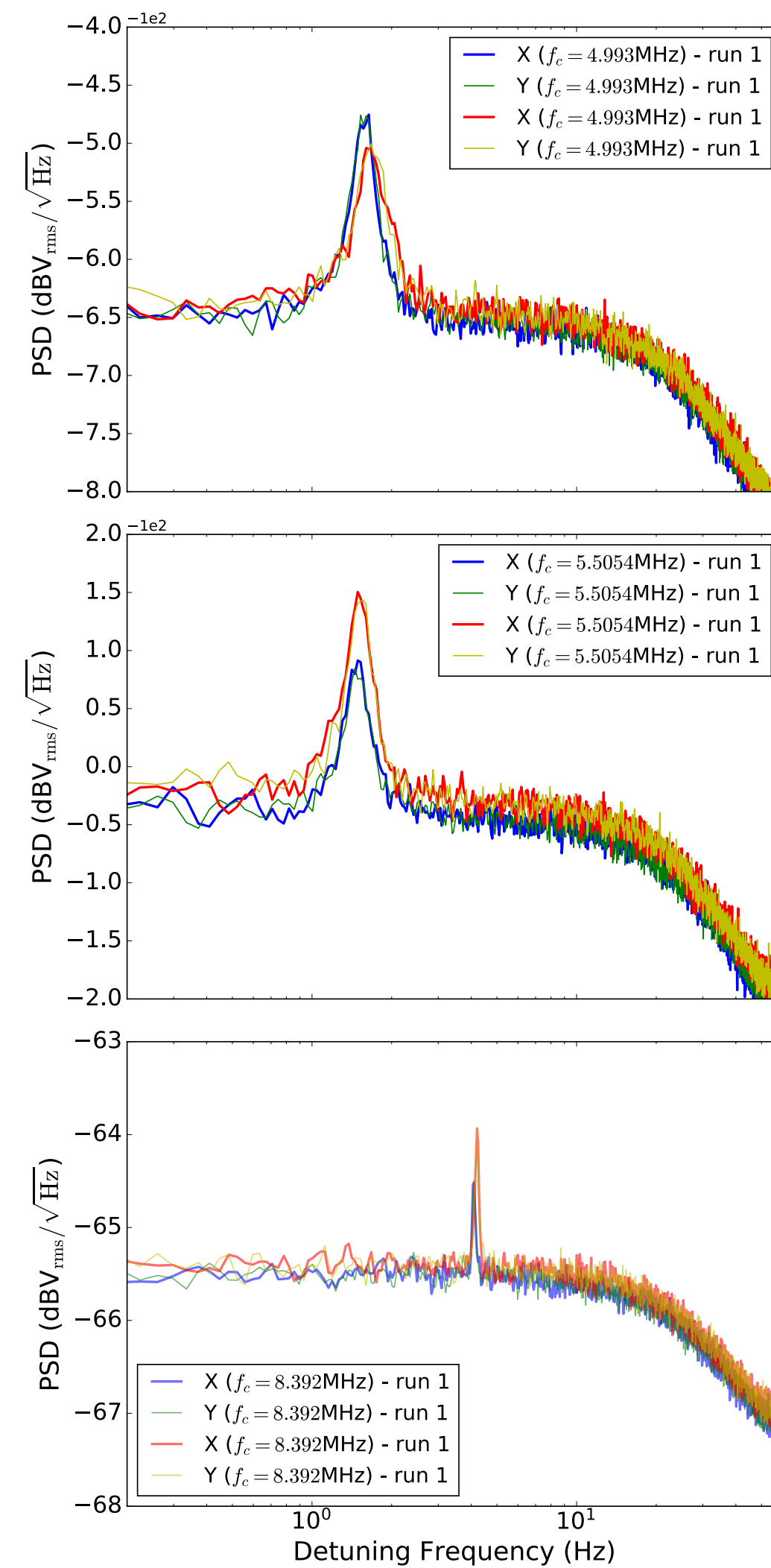


Can use Passive Bulk Acoustic Wave Resonators at Low Temperatures

- * Avoid the non-linear regime
- * Measuring Thermal Noise; Continuous; ~ 1 year
- * Search: GWs; Scalar Dark Matter; Test Quantum Gravity



20 mK



Scalar Dark Matter

PRL 116, 031102 (2016)

PHYSICAL REVIEW LETTERS

week ending
22 JANUARY 2016

Sound of Dark Matter: Searching for Light Scalars with Resonant-Mass Detectors

Asimina Arvanitaki,^{1,*} Savas Dimopoulos,^{2,†} and Ken Van Tilburg^{2,‡}

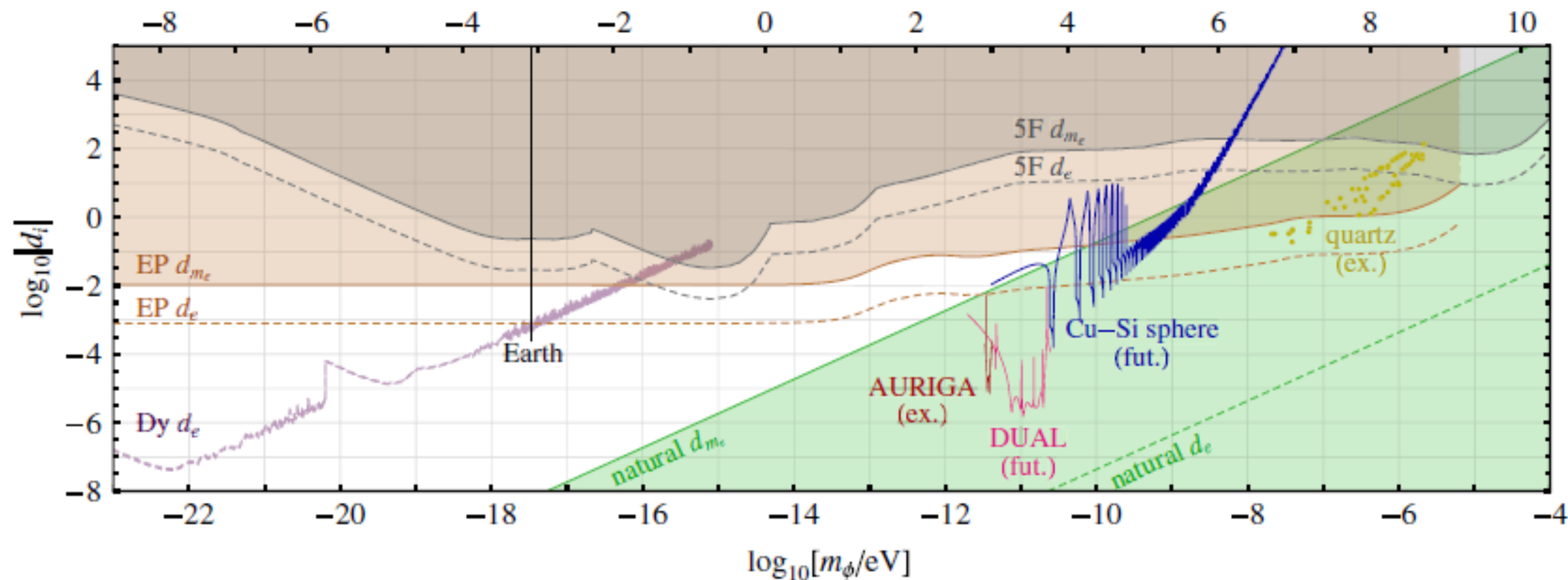
¹Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada

²Stanford Institute for Theoretical Physics, Stanford University, Stanford, California 94305, USA

(Received 23 August 2015; revised manuscript received 17 October 2015; published 22 January 2016)

The fine-structure constant and the electron mass in string theory are determined by the values of scalar fields called moduli. If the dark matter takes on the form of such a light modulus, it oscillates with a frequency equal to its mass and an amplitude determined by the local dark-matter density. This translates into an oscillation of the size of a solid that can be observed by resonant-mass antennas. Existing and planned experiments, combined with a dedicated resonant-mass detector proposed in this Letter, can probe dark-matter moduli with frequencies between 1 kHz and 1 GHz, with much better sensitivity than searches for fifth forces.

DOI: 10.1103/PhysRevLett.116.031102



Scalar Dark Matter

Searching for Scalar Dark Matter with Compact Mechanical Resonators

Jack Manley¹, Dalziel J. Wilson², Russell Stump¹, Daniel Grin³, and Swati Singh^{1,*}

¹Department of Electrical and Computer Engineering, University of Delaware, Newark, Delaware 19716, USA

²College of Optical Sciences, University of Arizona, Tucson, Arizona 85721, USA

³Department of Physics and Astronomy, Haverford College, Haverford, Pennsylvania 19041, USA

(Received 21 November 2019; accepted 18 March 2020; published 16 April 2020)

Ultralight scalars are an interesting dark matter candidate that may produce a mechanical signal by modulating the Bohr radius. Recently it has been proposed to search for this signal using resonant-mass antennas. Here, we extend that approach to a new class of existing and near term compact (gram to kilogram mass) acoustic resonators composed of superfluid helium or single crystal materials, producing displacements that are accessible with opto- or electromechanical readout techniques. We find that a large unprobed parameter space can be accessed using ultrahigh- Q , cryogenically cooled centimeter-scale mechanical resonators operating at 100 Hz–100 MHz frequencies, corresponding to 10^{-12} – 10^{-6} eV scalar mass range.

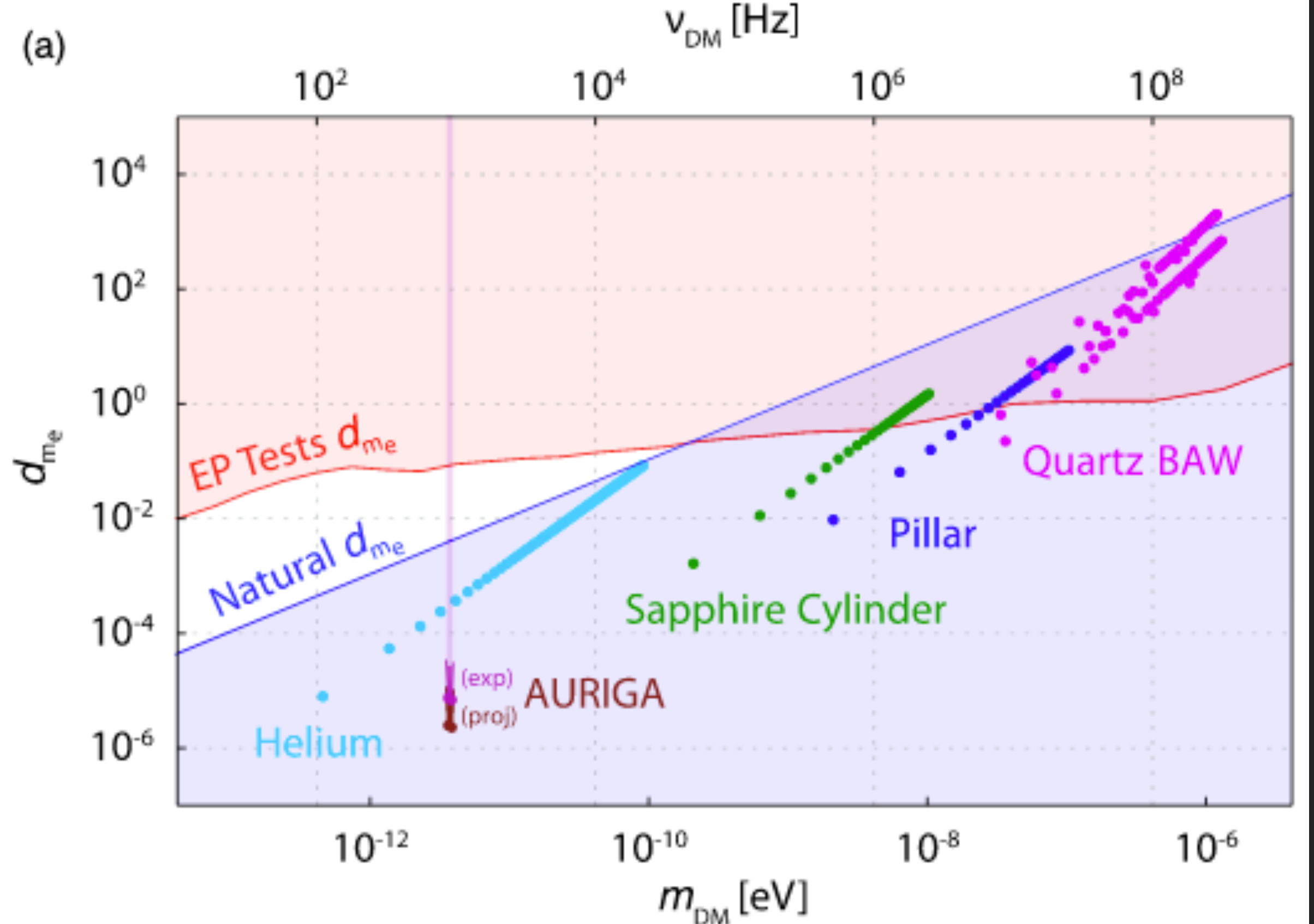
DOI: 10.1103/PhysRevLett.124.151301

DOI: 10.1103/PhysRevLett.124.151301

mass range

mechanical resonators operating at 100 Hz–100 MHz frequencies, corresponding to 10^{-12} – 10^{-6} eV scalar mass range.

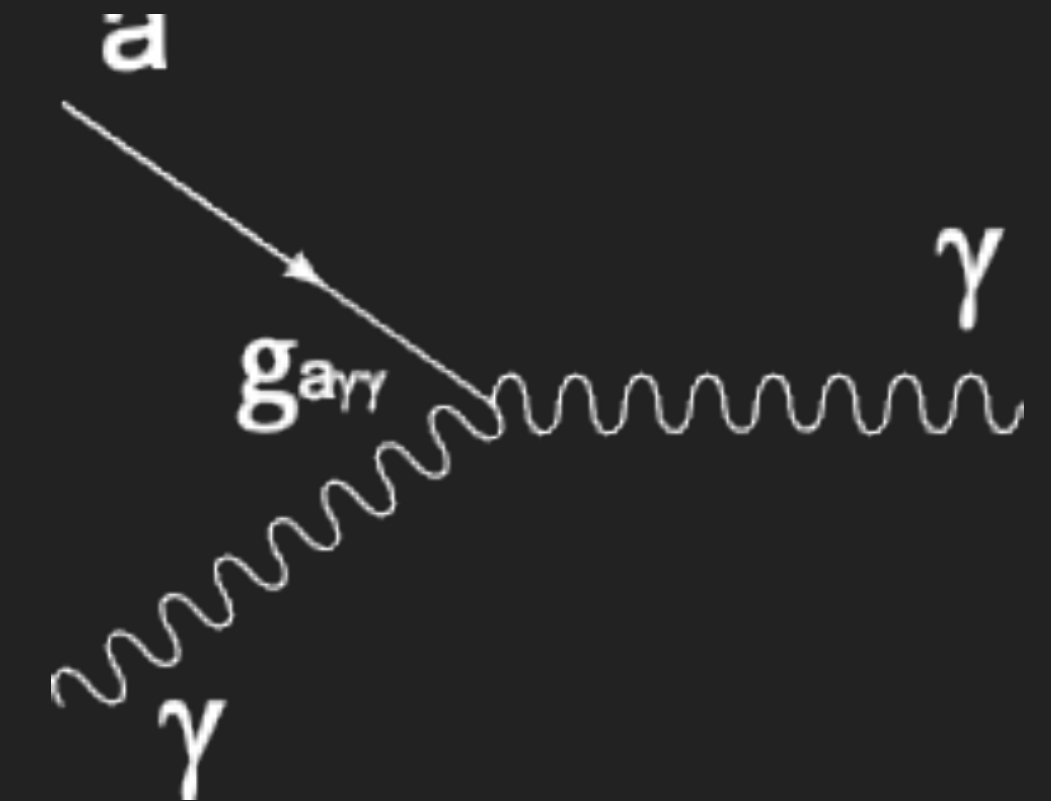
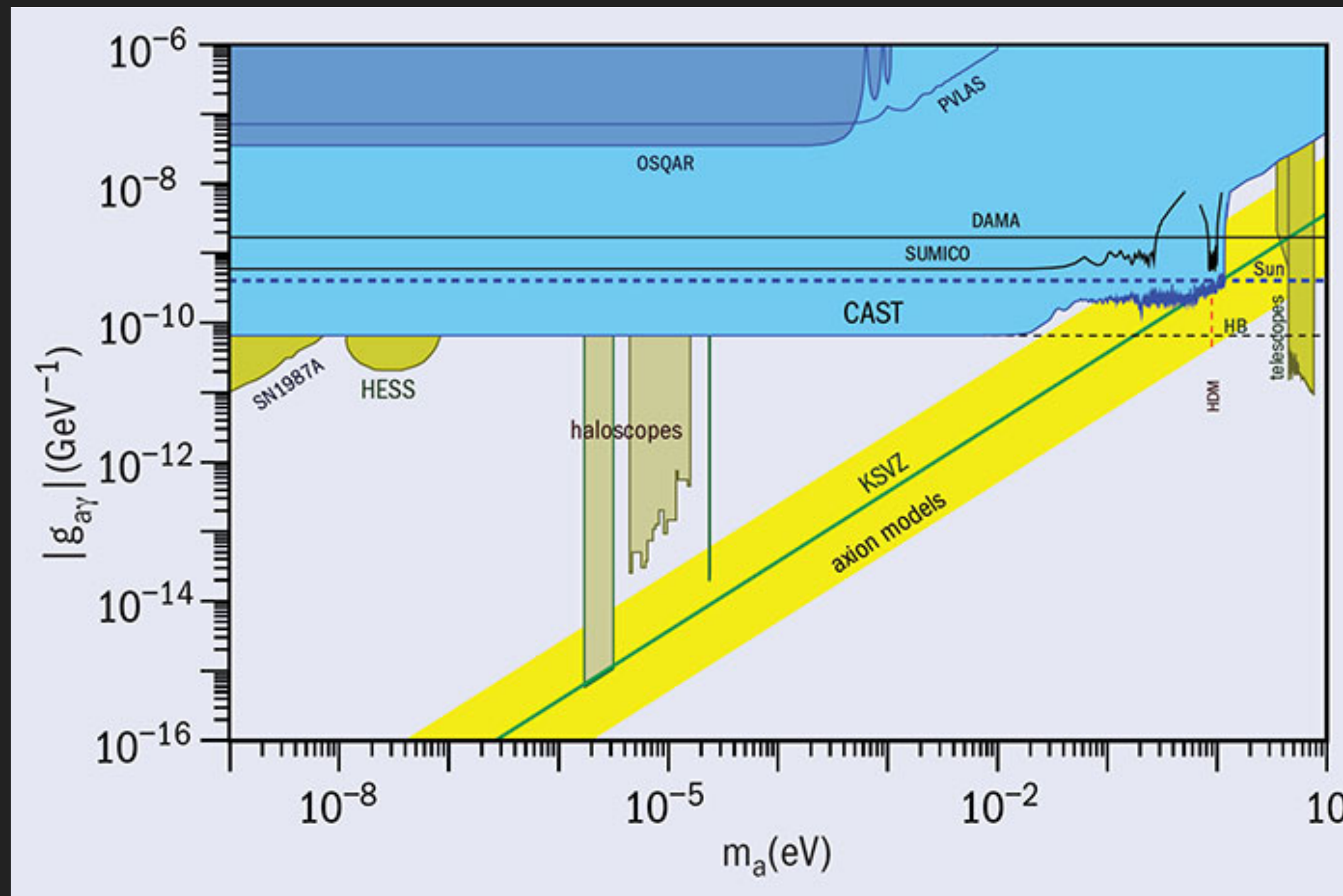
$$(d_{\text{DM}})_{\text{min}} \approx \sqrt{\frac{c^2}{8\pi G \rho_{\text{DM}}}} \omega_n h_{\text{min}}$$



AXION as Dark Matter

Two most popular classes of phenomenologically viable invisible axion models:

- Hadronic or Kim-Shifman-Vainshtein-Zakharov (KSVZ) model
- Dine-Fischler-Srednicki-Zhitnitskii (DFSZ) model



ORGAN: Oscillating Resonant Group Axion Experiment

McGillivray Organ at UWA



Covers 15-50 GHz range

Promising theoretical models point to high mass ranges:

- Lattice QCD Simulations
- Encompassing theories like “SMASH” - SM Axion Seesaw Higgs Portal Inflation
- Strange experimental observations in Josephson Junctions (albeit quite controversial)
- Mature experiments already covering large portions of the lower mass range

with higher frequencies everything gets more complicated!

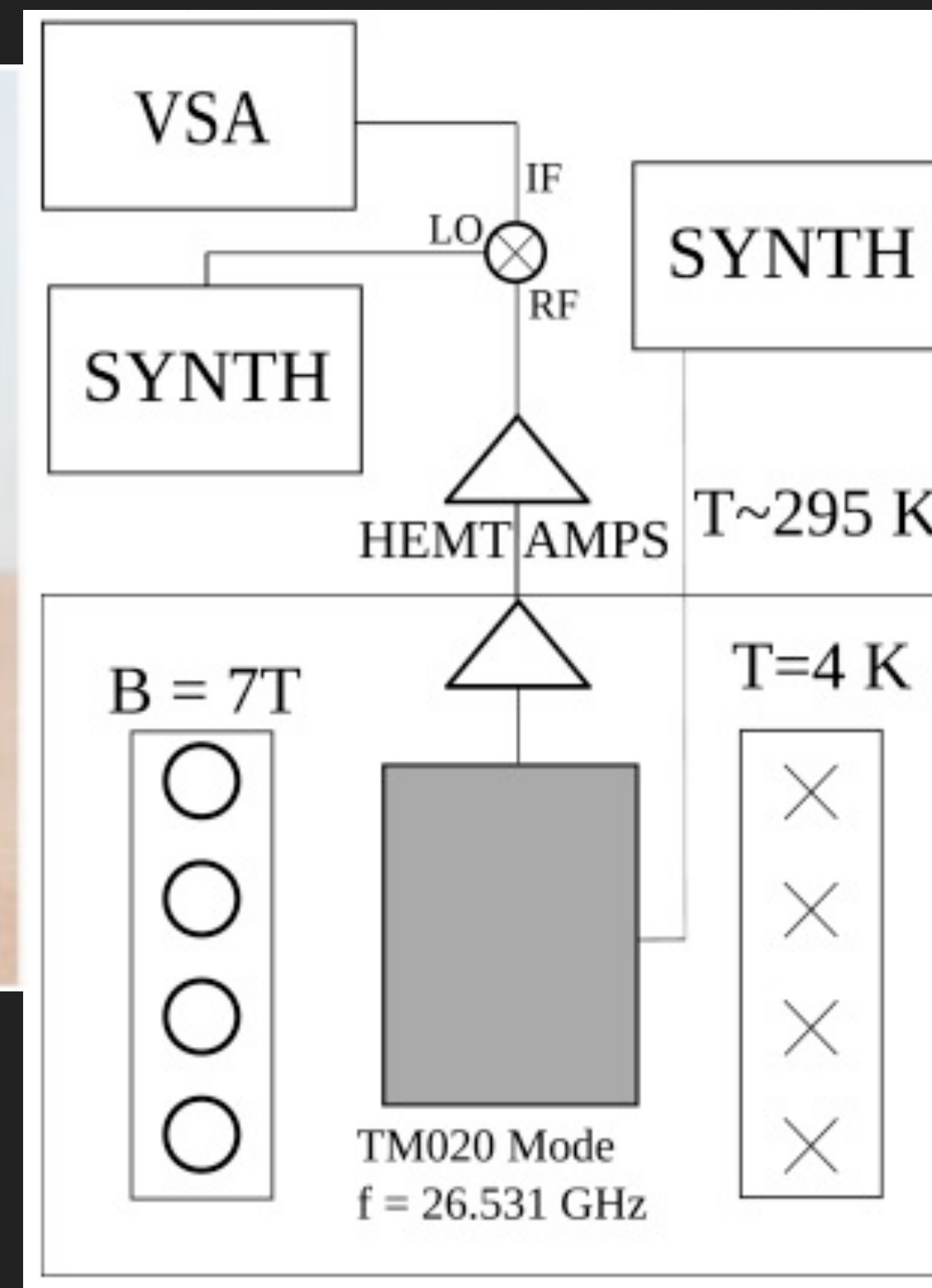
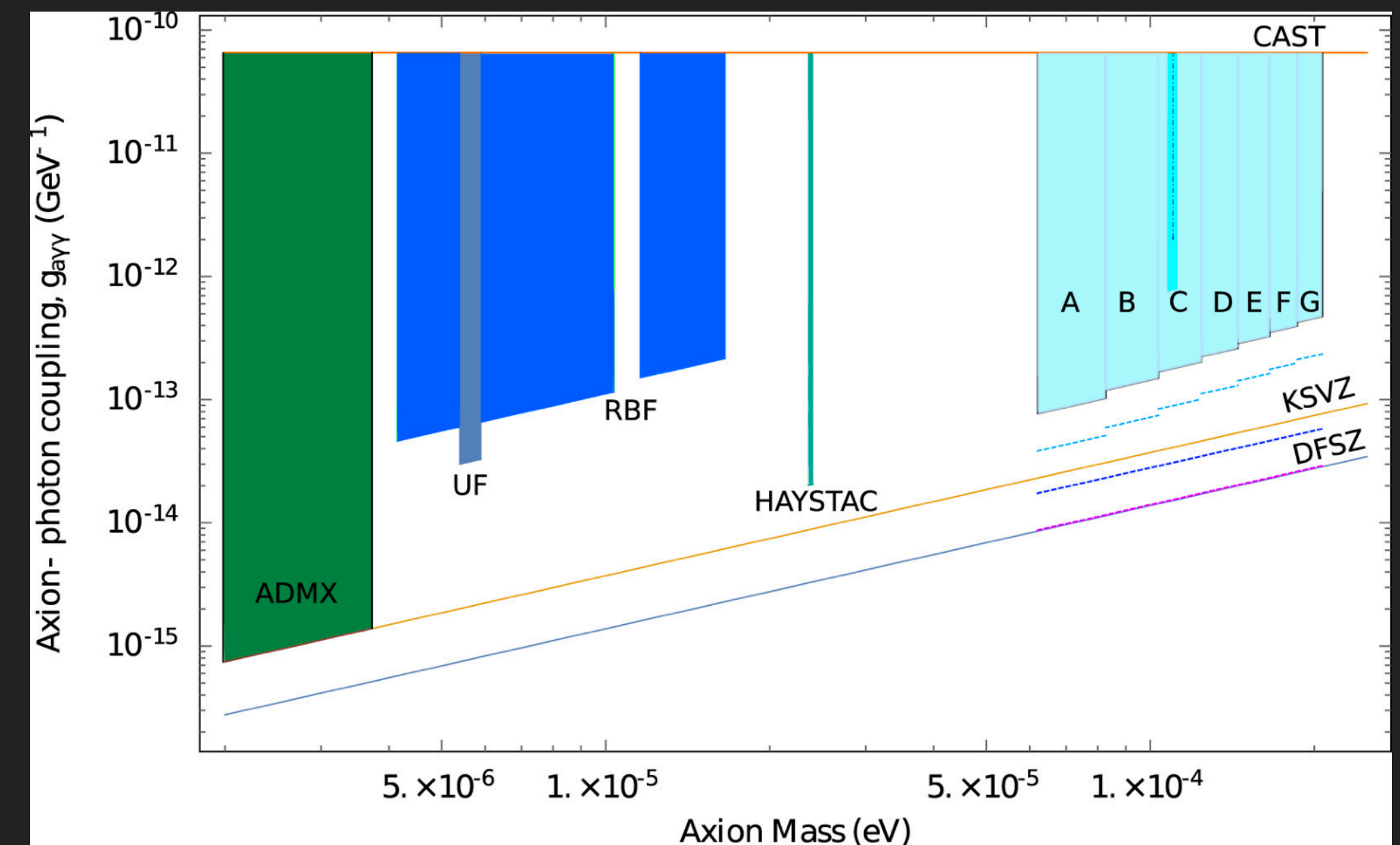
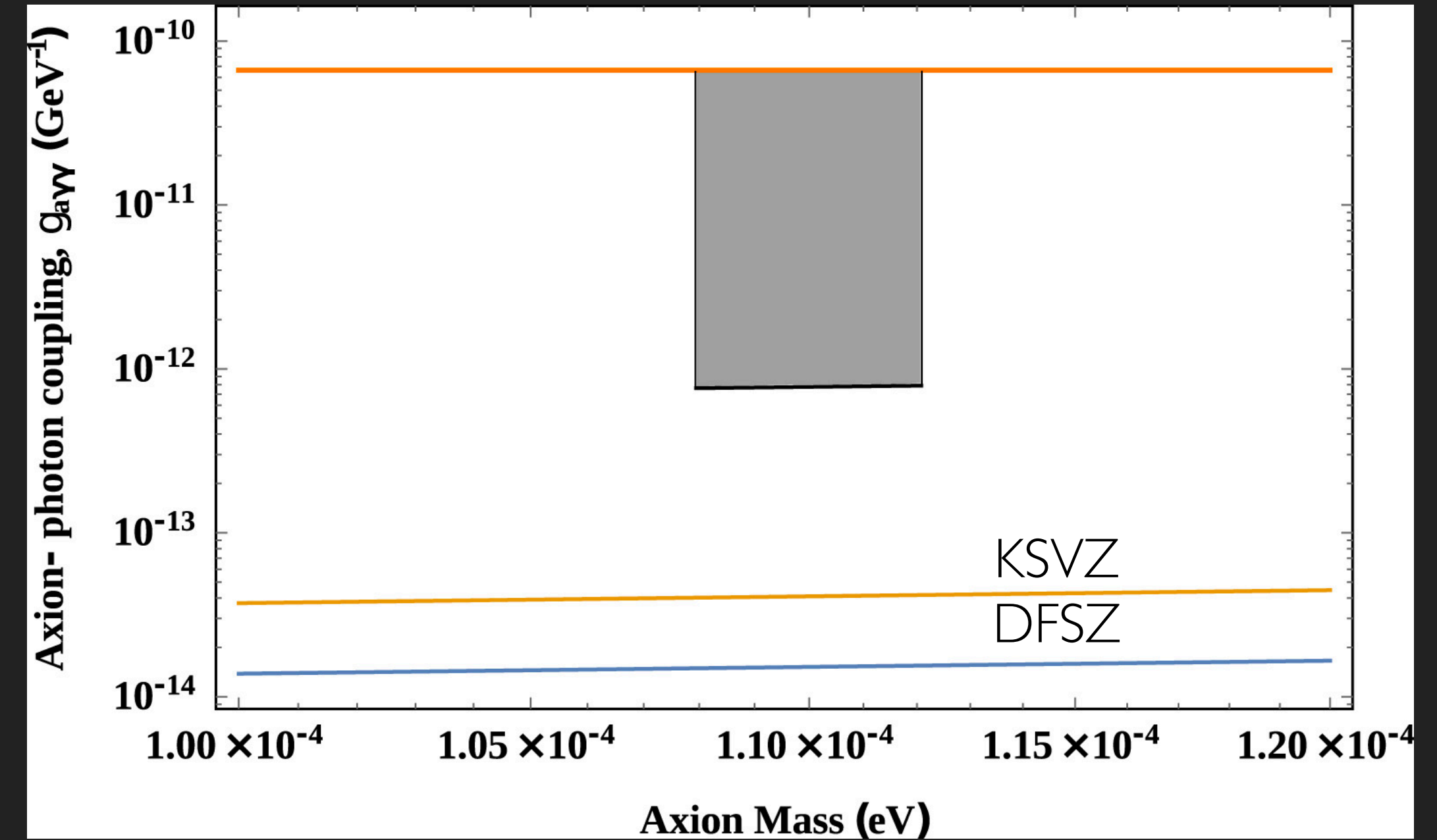
$$\frac{df}{dt} \propto \frac{1}{SNR_{goal}^2} \frac{g_{a\gamma\gamma}^4 B^4 C^2 V^2 \rho_a^2 Q_L Q_a}{m_a^2 (k_B T_n)^2}$$

ORGAN: Oscillating Resonant Group AxioN Experiment


 Physics of the Dark Universe
 Volume 18, December 2017, Pages 67-72

The ORGAN experiment: An axion haloscope above 15 GHz

Ben T. McAllister ^a, Graeme Flower ^a, Eugene N. Ivanov ^b, Maxim Goryachev ^a, Jeremy Bourhill ^a, Michael E. Tobar ^a



ORGAN: Oscillating Resonant Group Axion Experiment

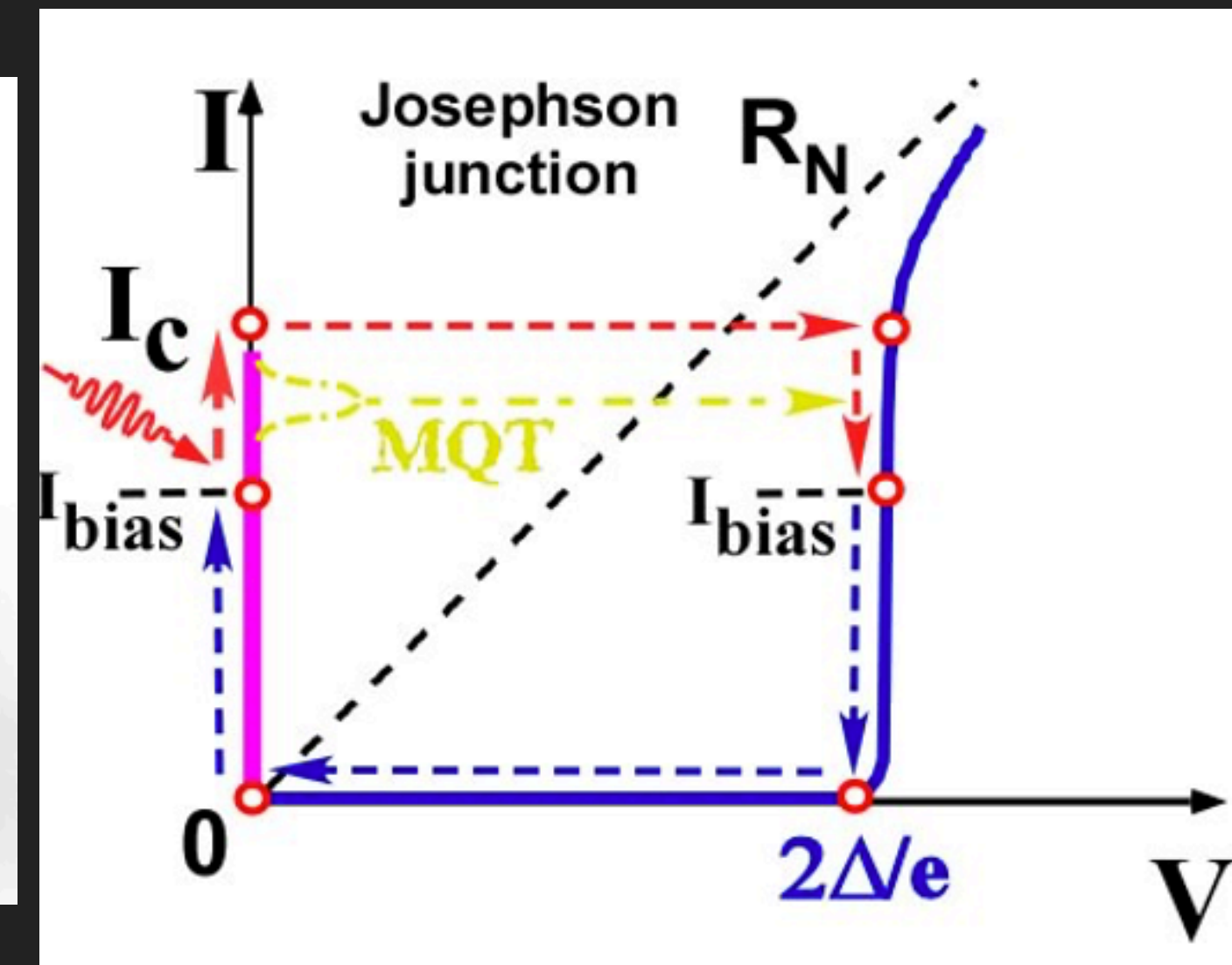
Phase I a began in June



Fridge and Magnet (12-14T)



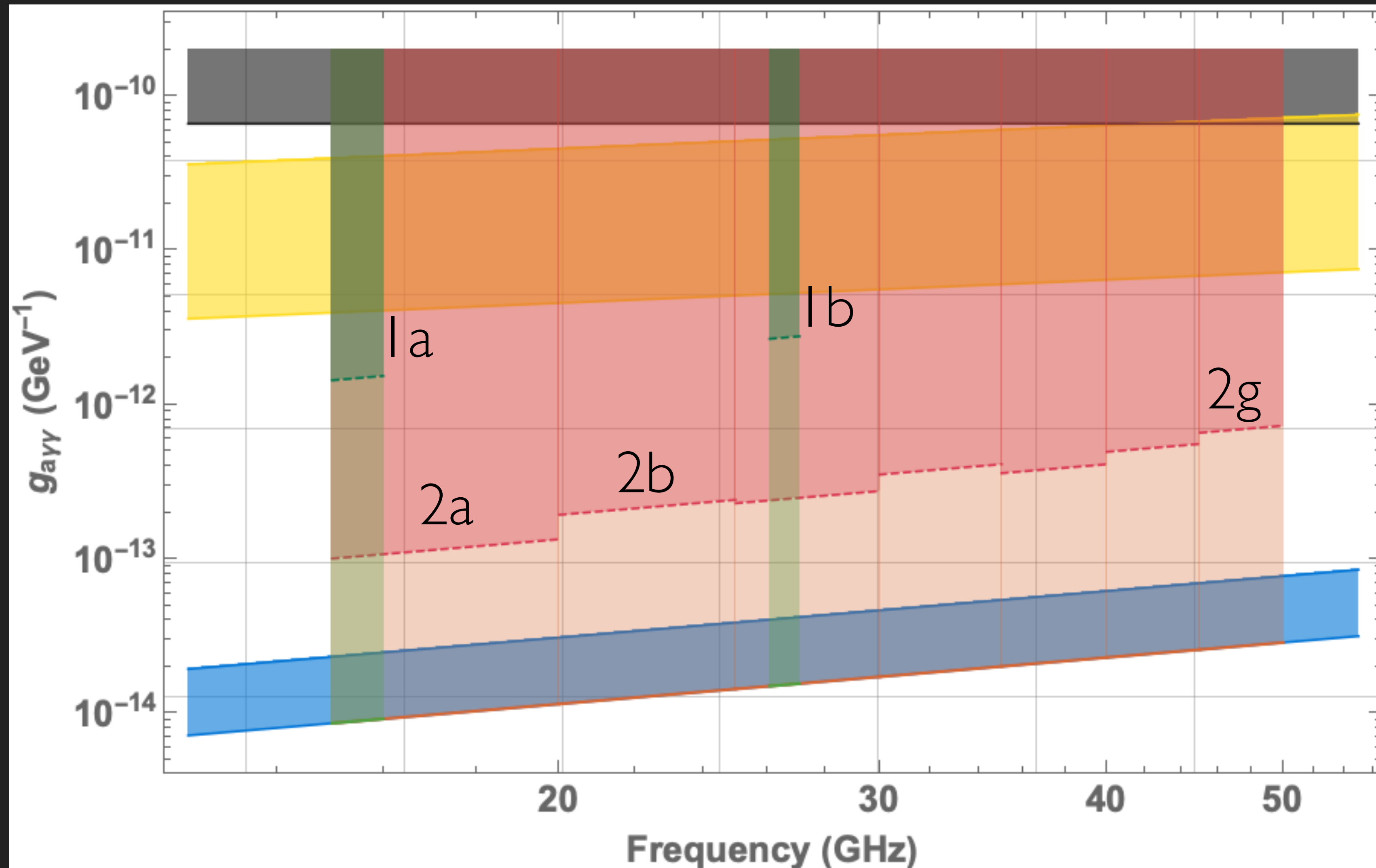
Cryogenic translation stage



Detection (GHz SPD)

ORGAN: Oscillating Resonant Group AxioN Experiment

ALP
Cogenesis



HEMPT or SQL
Limited Linear
amplifier

Efficient GHz
SPC



Upconversion Loop Oscillator Axion Detection Experiment: A Precision Frequency Interferometric Axion Dark Matter Search with a Cylindrical Microwave Cavity

Catriona A. Thomson^{✉*}, Ben T. McAllister, Maxim Goryachev, Eugene N. Ivanov, and Michael E. Tobar[†]
ARC Centre of Excellence for Engineered Quantum Systems and ARC Centre of Excellence for Dark Matter Particle Physics,
Department of Physics, University of Western Australia, 35 Stirling Highway, Crawley, Western Australia 6009, Australia

Physics of the Dark Universe 26 (2019) 100345

Contents lists available at ScienceDirect

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journal homepage: www.elsevier.com/locate/dark

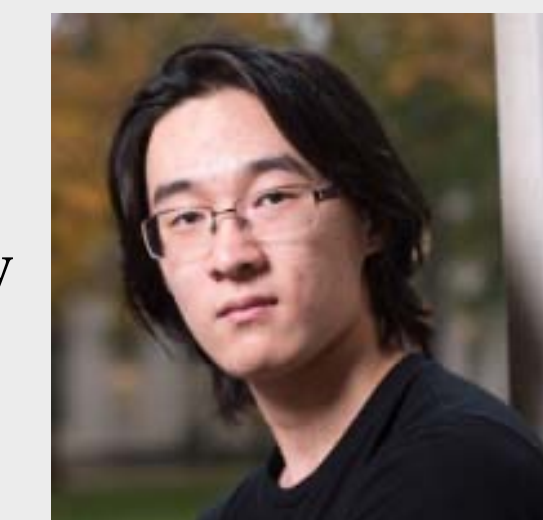


First Proposed in 2018
[arXiv:1806.07141 \[physics.ins-det\]](https://arxiv.org/abs/1806.07141)
Proposal Published in 2019

Erratum: UPconversion Loop Oscillator Axion Detection experiment: A precision frequency interferometric axion dark matter search with a Cylindrical Microwave Cavity

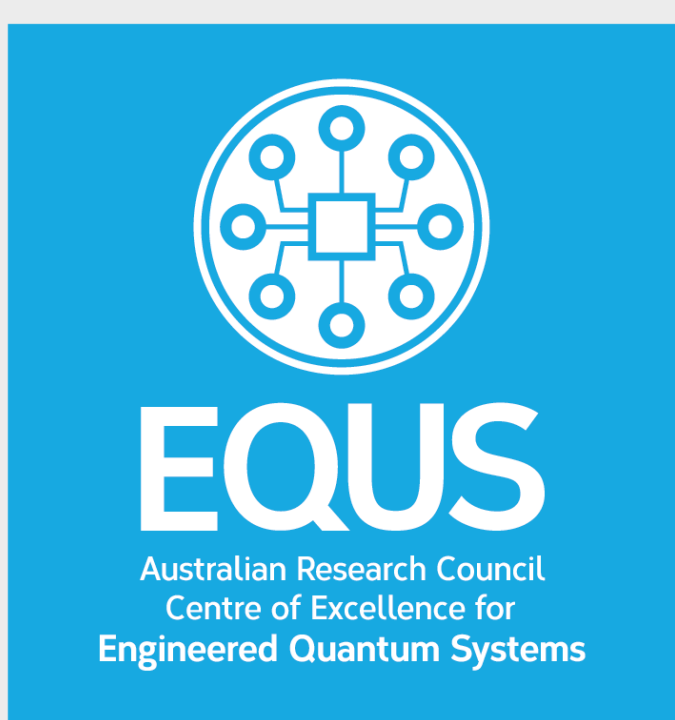
Catriona A. Thomson,^{*} Ben T. McAllister, Maxim Goryachev, Eugene N. Ivanov, and Michael E. Tobar[†]
ARC Centre of Excellence for Engineered Quantum Systems and ARC Centre of Excellence for Dark Matter Particle Physics,
Department of Physics, University of Western Australia,
35 Stirling Highway, Crawley, WA 6009, Australia.
(Dated: June 10, 2021)

We are very grateful to Kevin Zhou for meticulously going through our paper and finding our mistake, and reading this correction making sure of its validity. This work was funded by the ARC Centre of Excellence for Engineered Quantum Systems, CE170100009, and Dark Matter Particle Physics, CE200100008.



Heterodyne Detection of Axion Dark Matter in an RF Cavity

Sebastian Ellis



Axion detection with precision frequency metrology

Maxim Goryachev, Ben T. McAllister^{*}, Michael E. Tobar

ARC Centre of Excellence for Engineered Quantum Systems, School of Physics, University of Western Australia, 35 Stirling Highway, Crawley, WA 6009, Australia



Axion-Photon Coupling to Search for Axion

The Axion Haloscope Technique

$$\mathcal{L} \propto ag_{a\gamma\gamma} \vec{E}_{cavity} \cdot \vec{B}_{ext}$$

Lagrangian gives effective strength

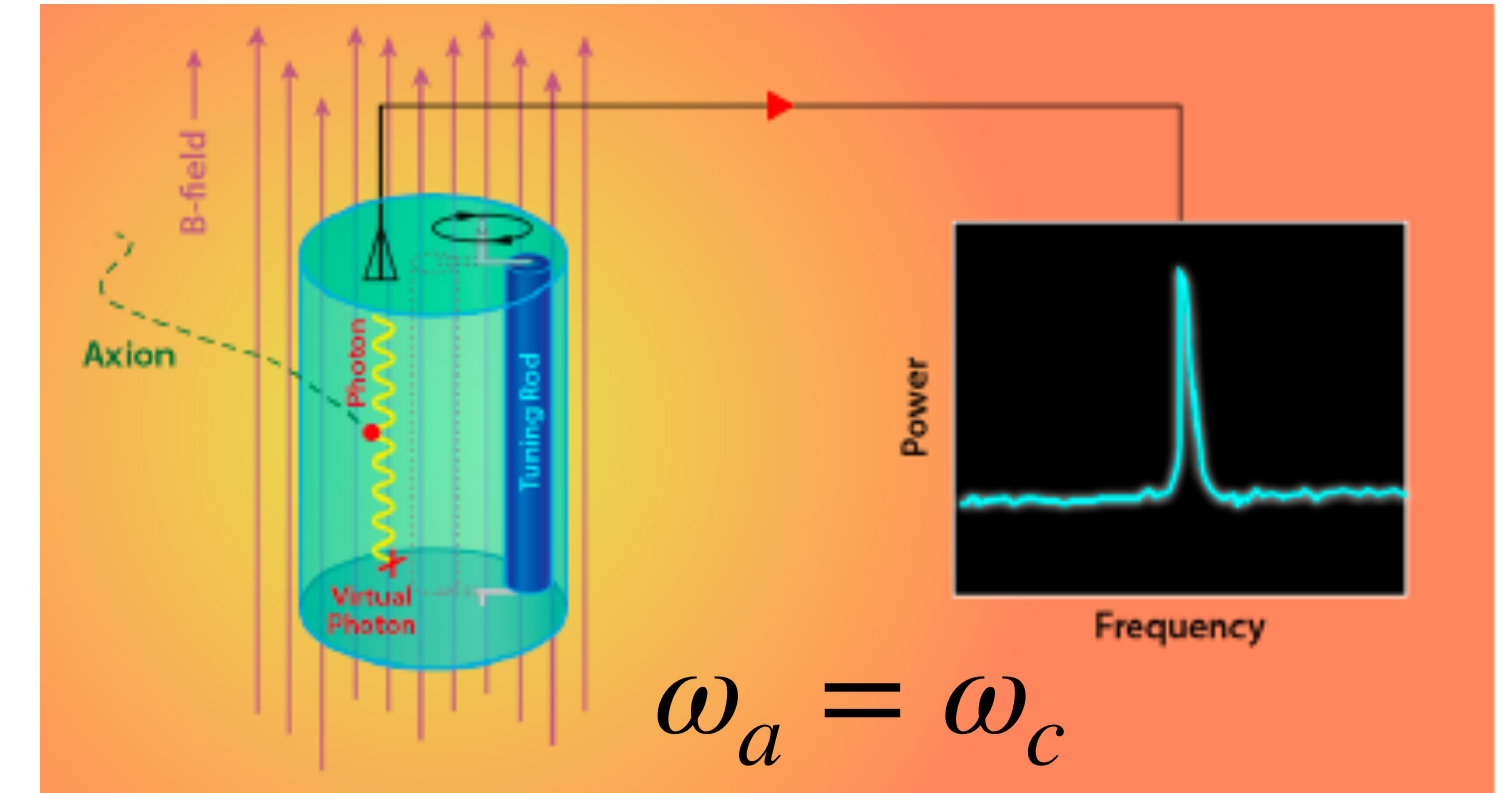
The AC Axion Haloscope Technique



a ————— a

$$\omega_a = \omega_1 \pm \omega_2$$

$g_{a\gamma\gamma}$



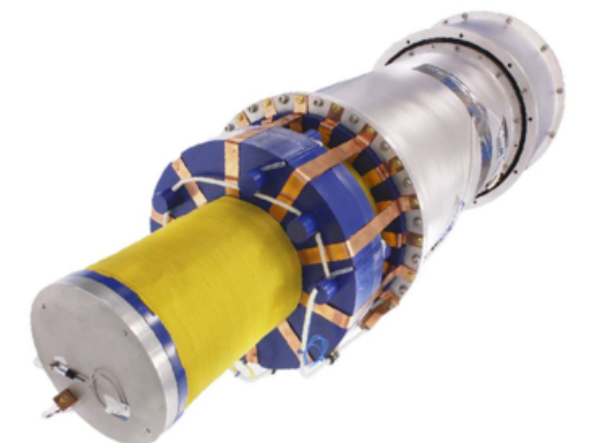
$$\omega_a = \omega_c$$



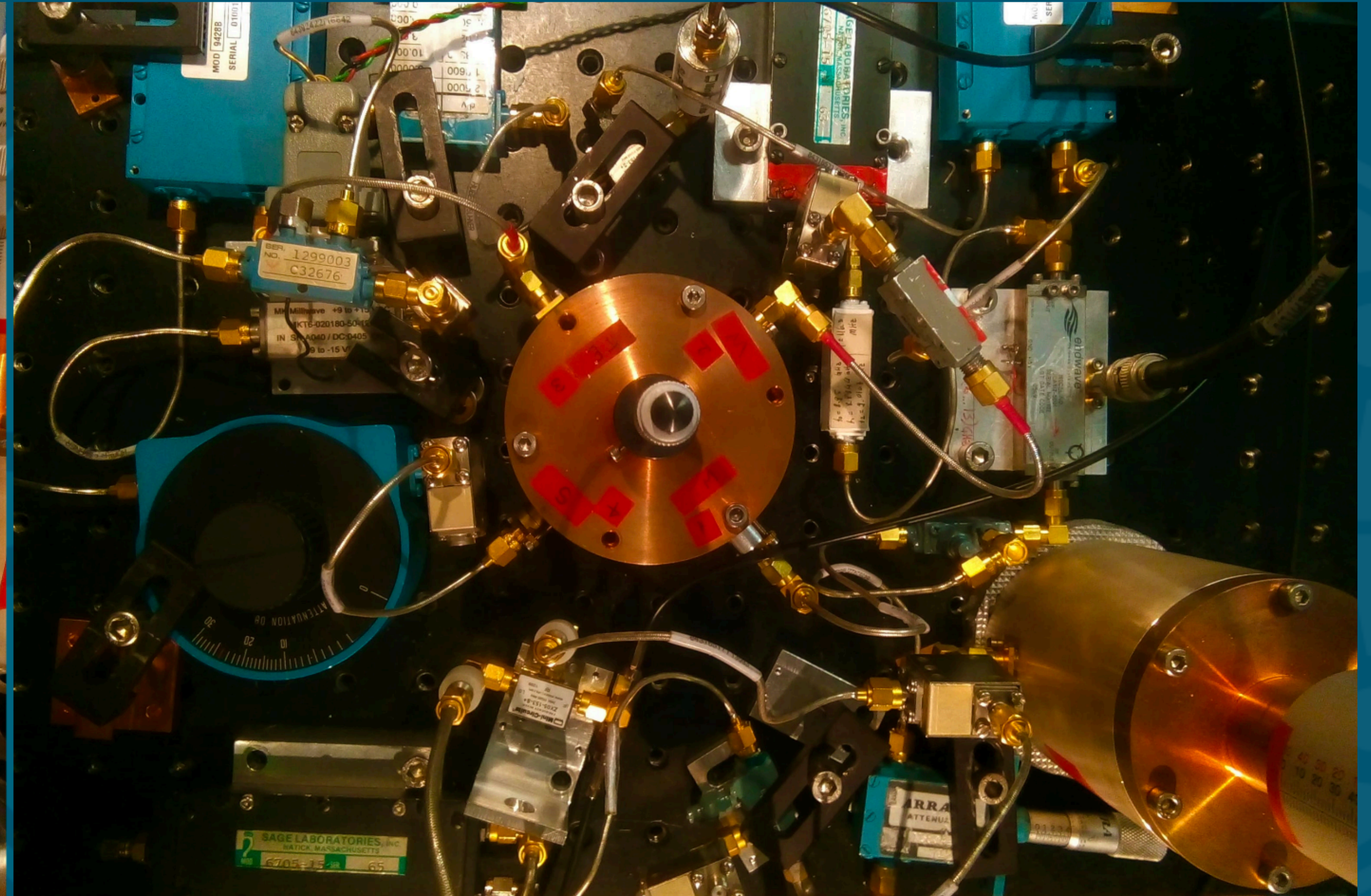
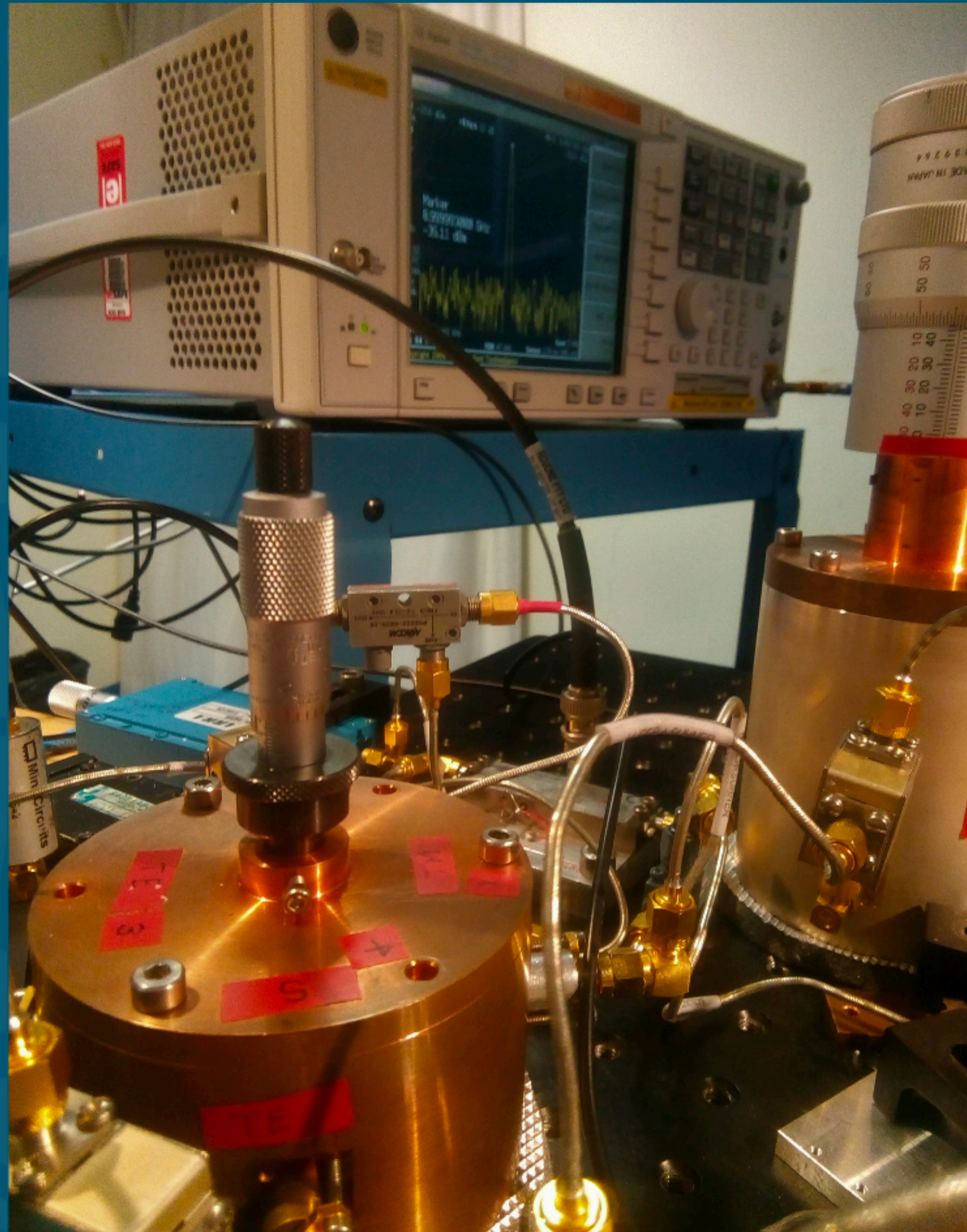
$g_{a\gamma\gamma}$

Microwave Cavity

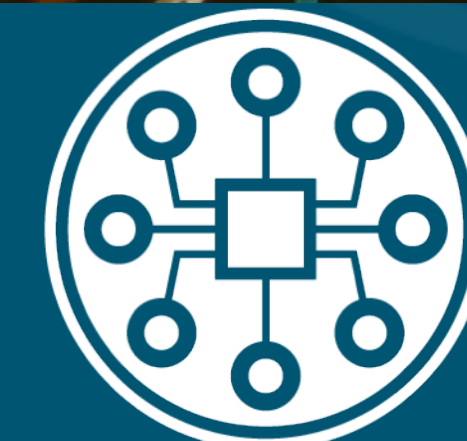
DC B-field



Construct Dual-Mode Loop Oscillator with Phase Noise Measurement System: What is the Signal to Noise Ratio?



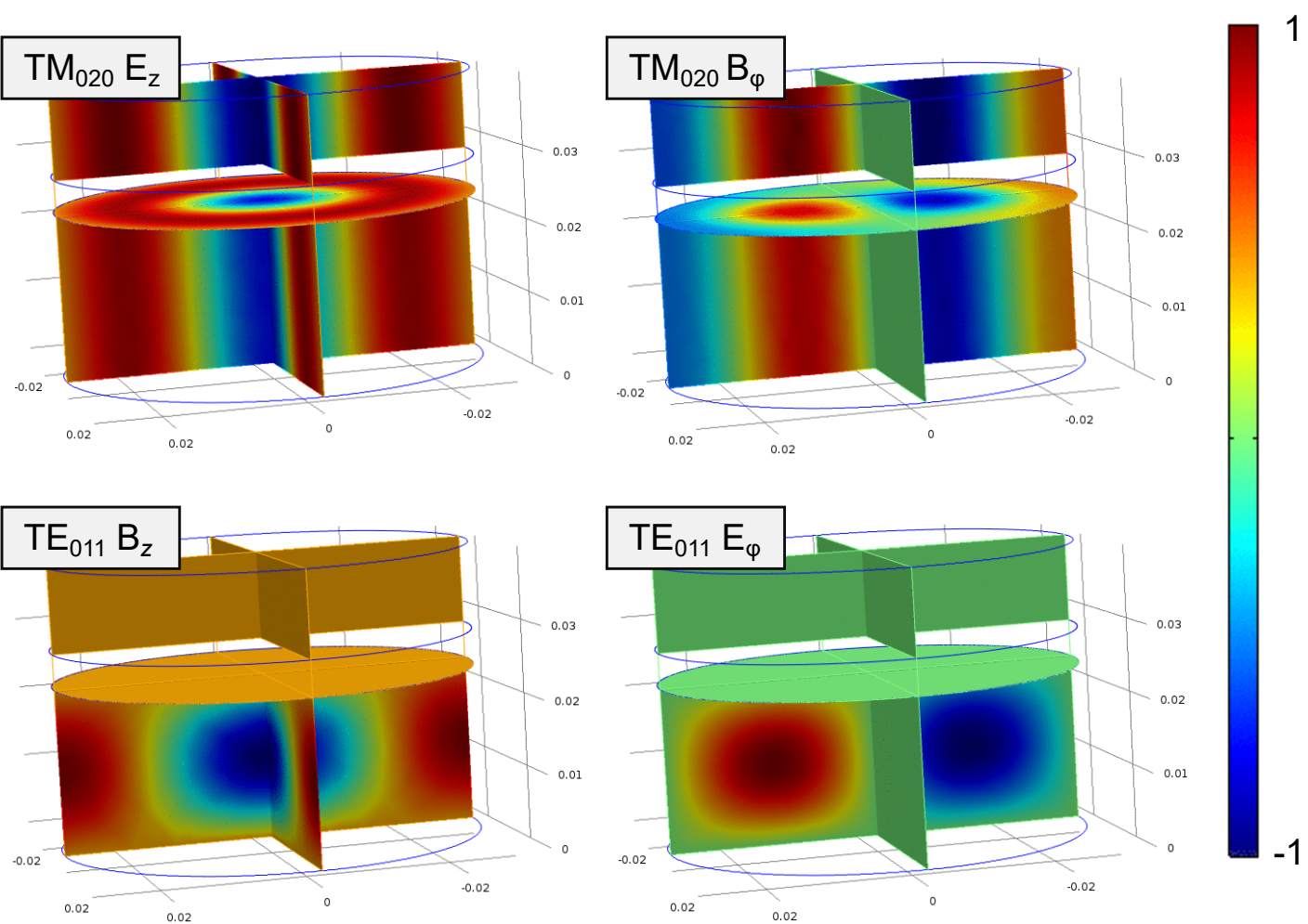
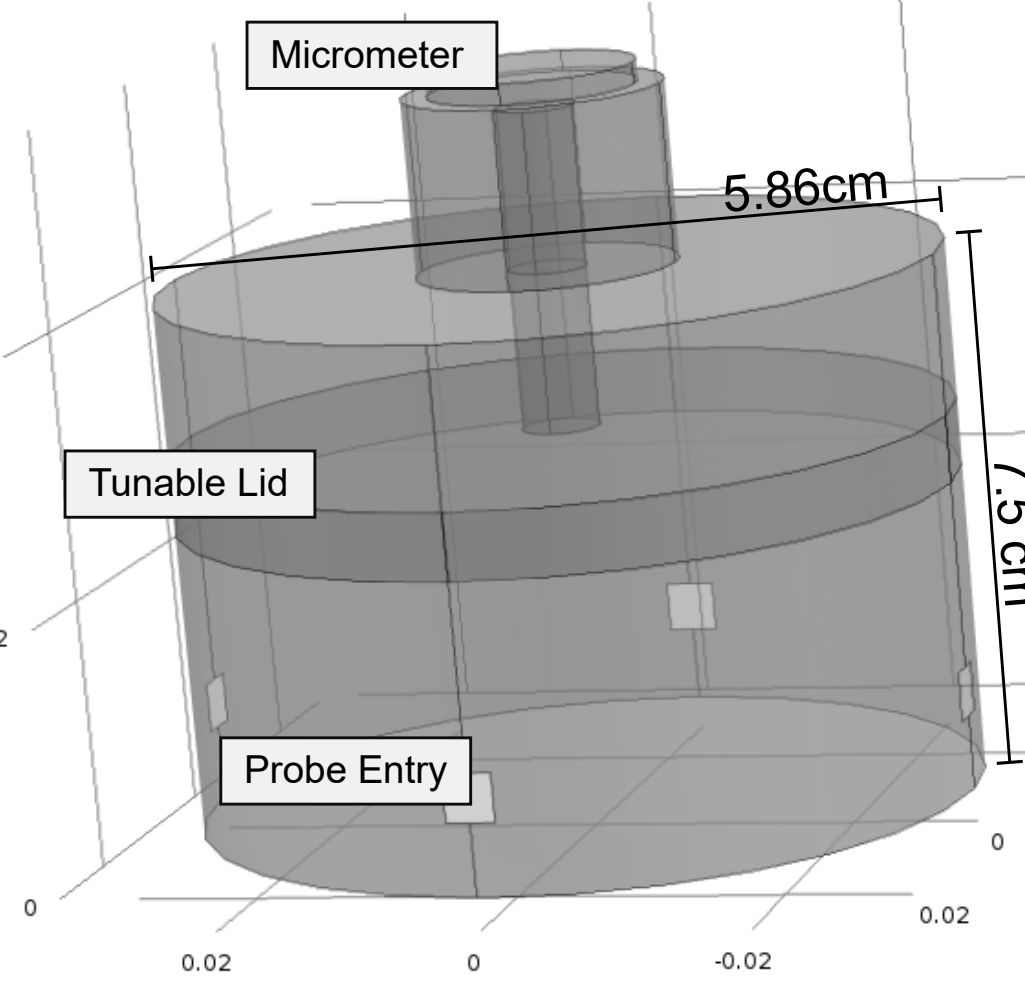
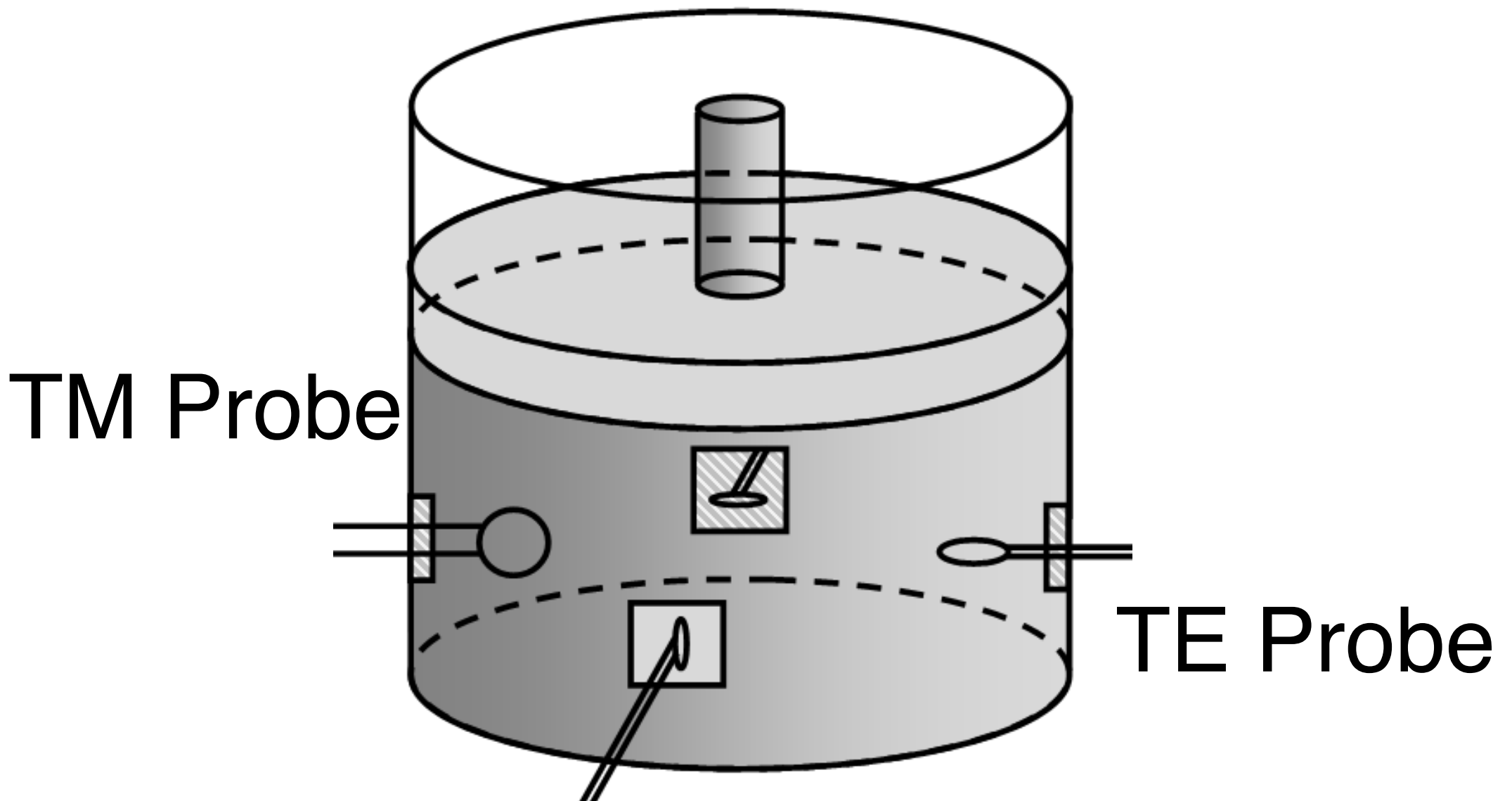
THE UNIVERSITY OF
**WESTERN
AUSTRALIA**



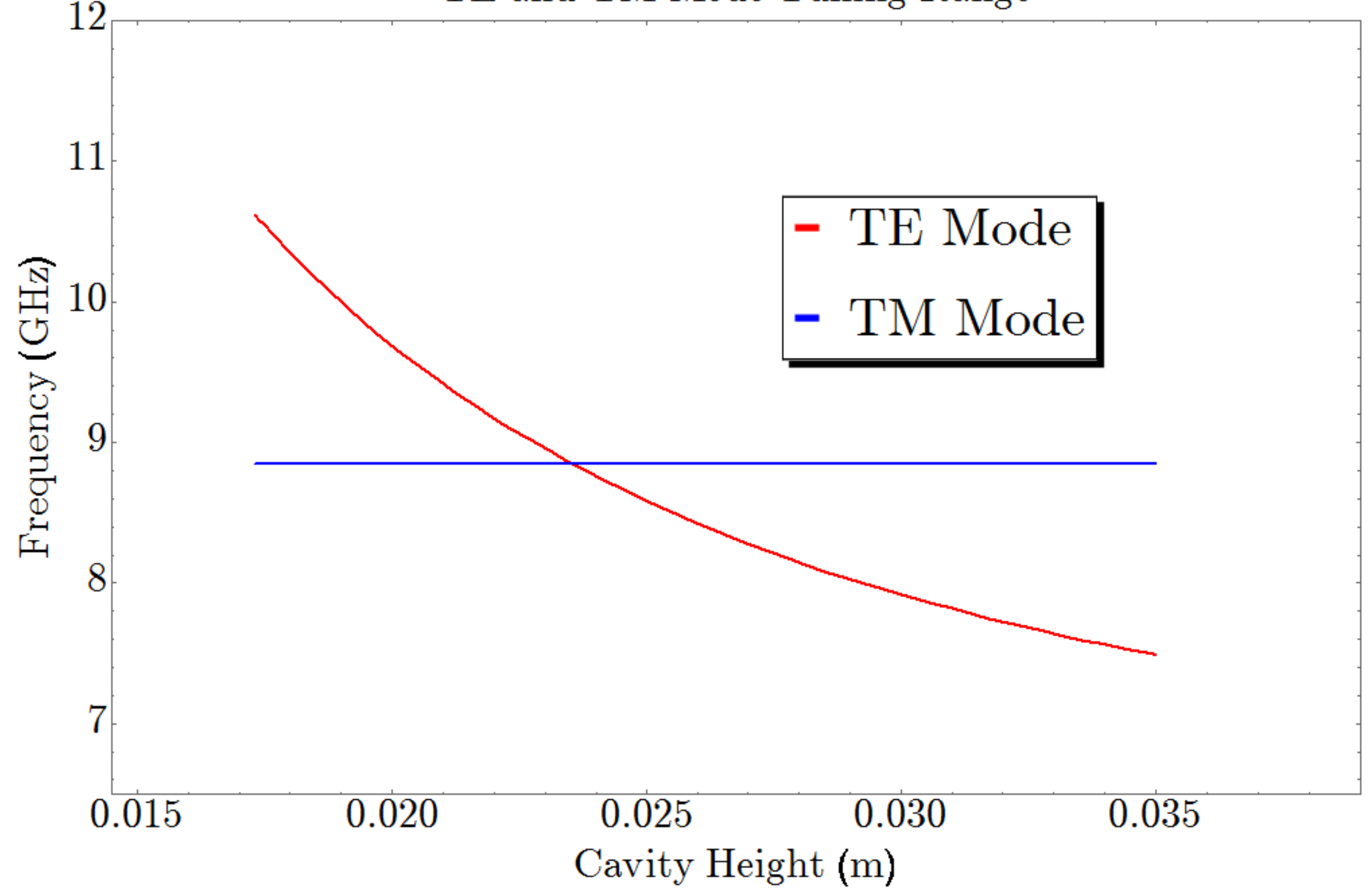
EQUUS
Australian Research Council
Centre of Excellence for
Engineered Quantum Systems

First Realisation: Cylindrical Microwave Cavity

- Tuneable cavity height (lid attached to micrometer)
- TM020 mode frequency fixed by cavity radius
- TE011 mode frequency tuned by cavity height



TE and TM Mode Tuning Range



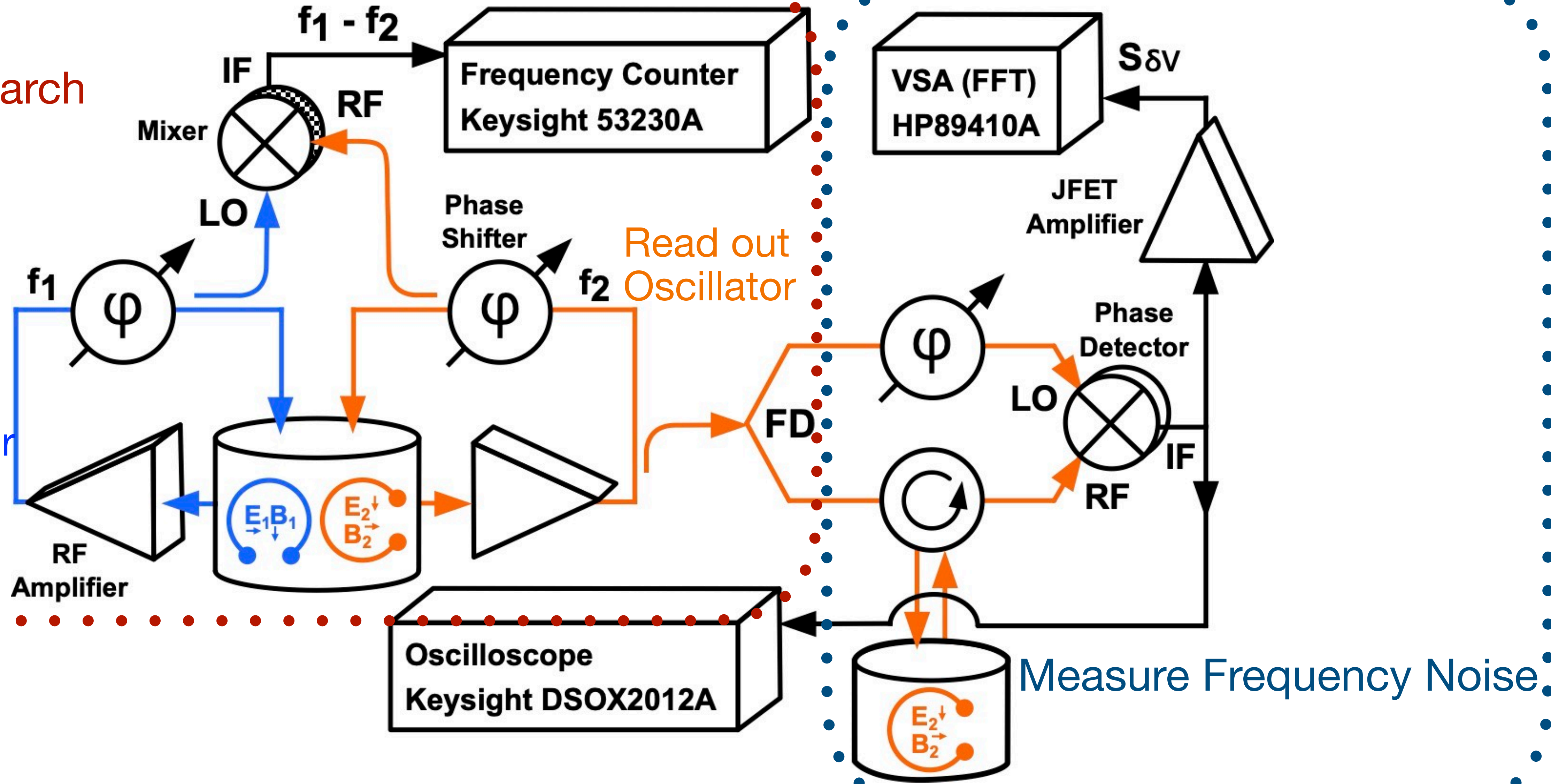
Down Conversion: $f_a \approx f_2 + f_1 \approx 18$ GHz

Up Conversion: $f_a \approx |f_2 - f_1| \rightarrow$ Low Mass

Schematic of the Experimental Setup

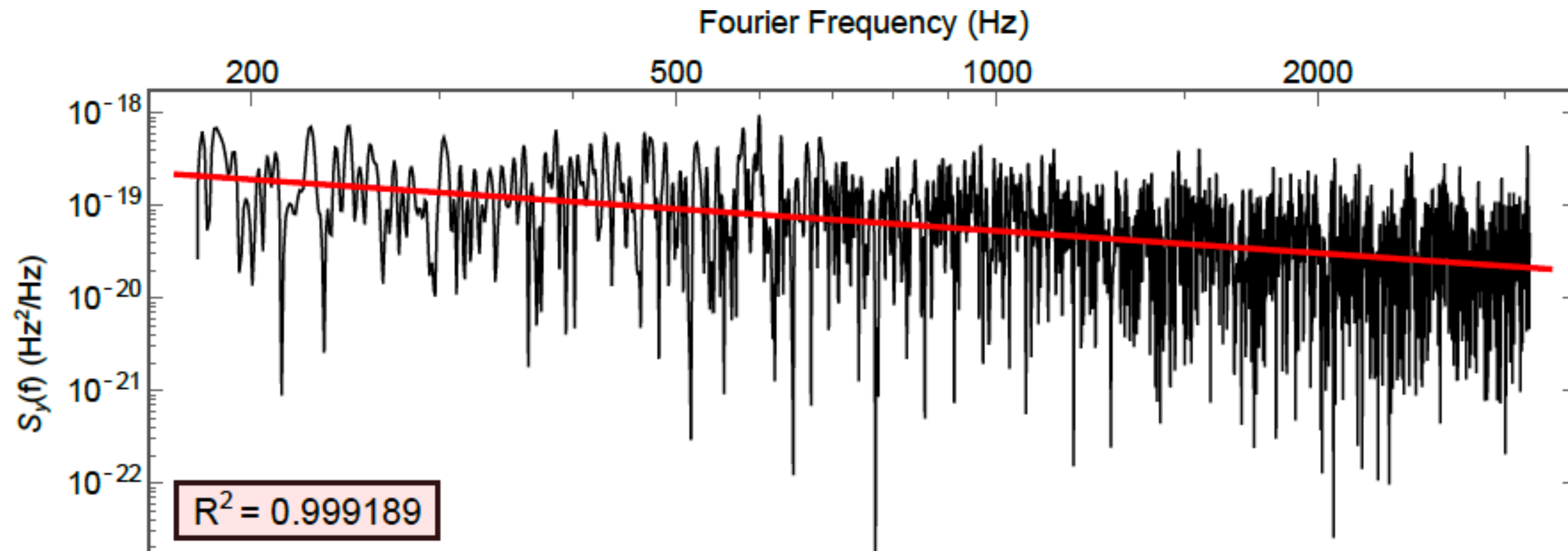
Axion Search
via Dual-
Mode

Pump
Oscillator



Oscillator Fractional Frequency Noise

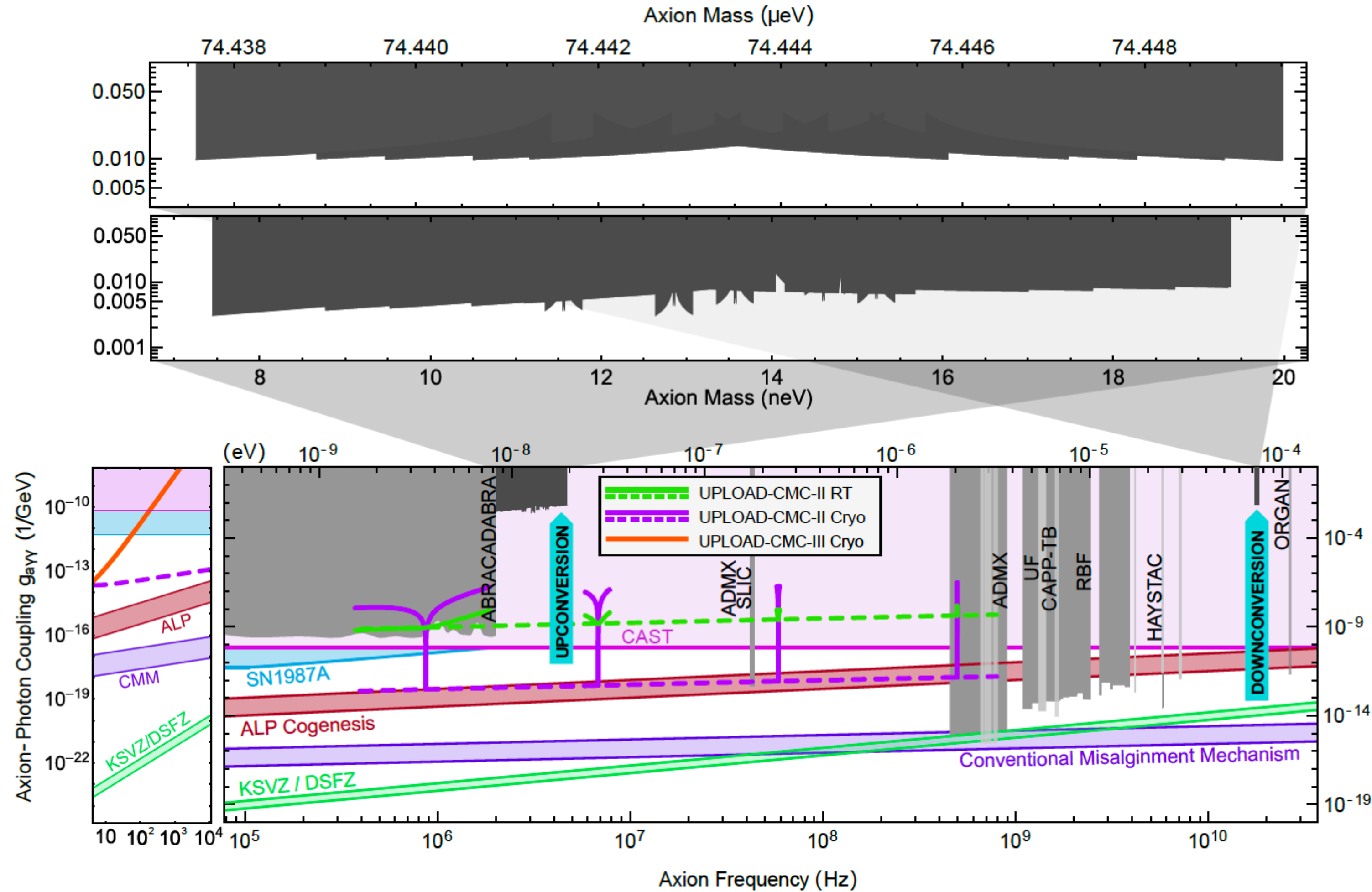
Search Fourier Frequency of read out oscillator for axions,
relation becomes $f_a = f_2 + f_1 \pm f$ or $f_a = |f_2 - f_1| \pm f$.



Searching at $\pm f$, Fourier frequencies at the same time

Next: What is the Signal that the Axion will imprint on the Phase Noise?

Sensitivity Limits



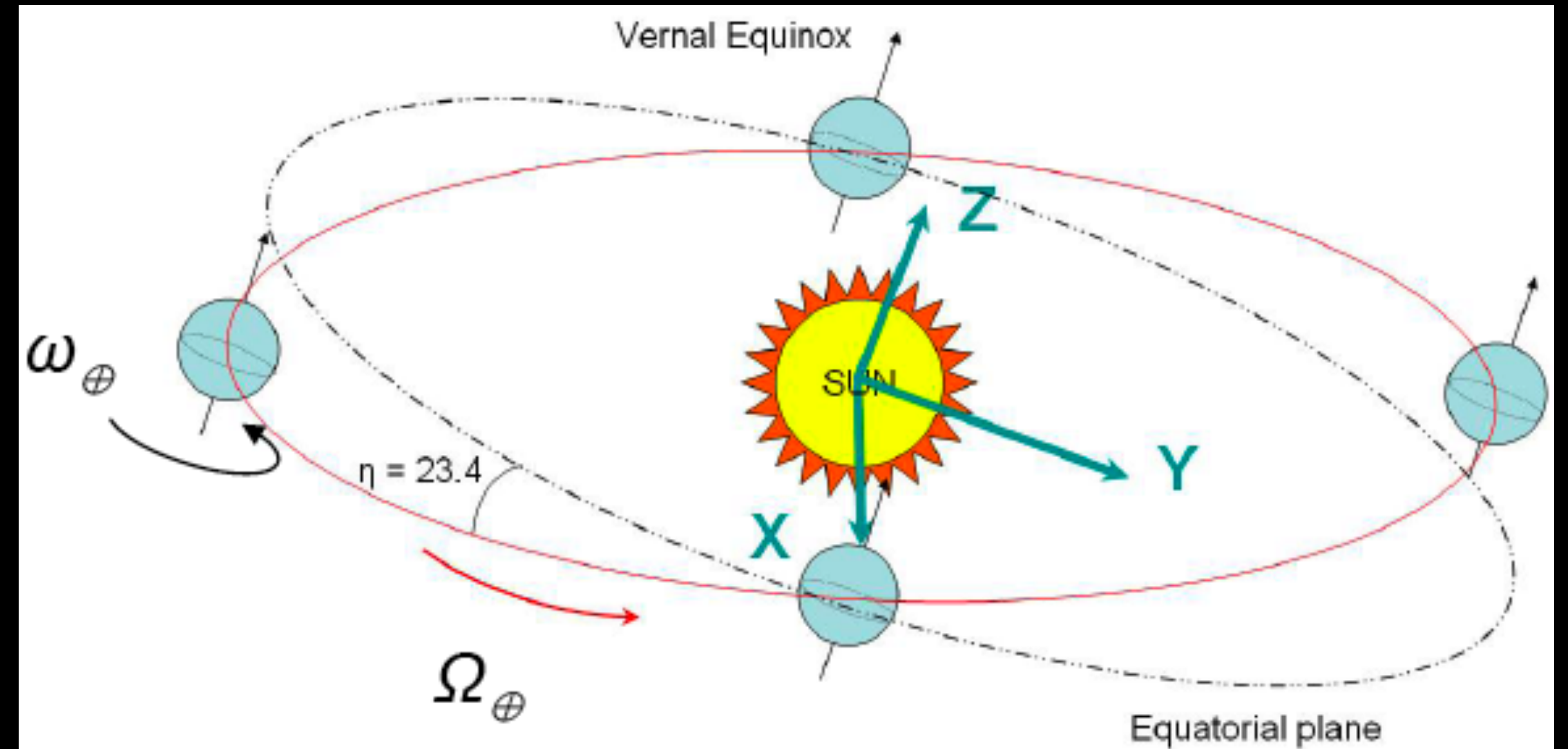
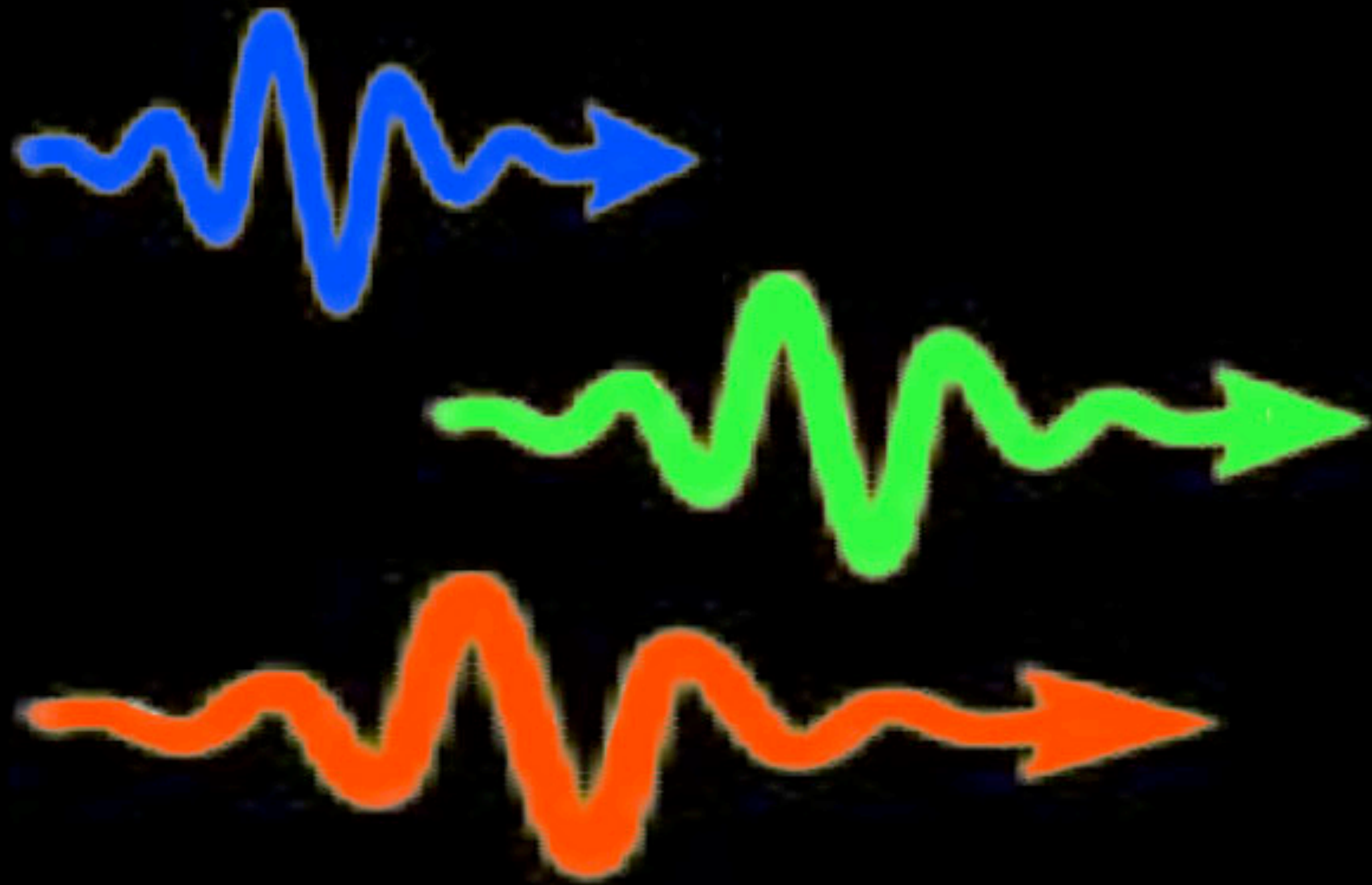
- Puts limits 7.44-19.38 neV with only 5 positions of cavity tuning
- New Room Temperature Version Under Construction
- Cryogenic Version under Design

• Looking at a few different schemes, injection locking etc. and power measurement schemes

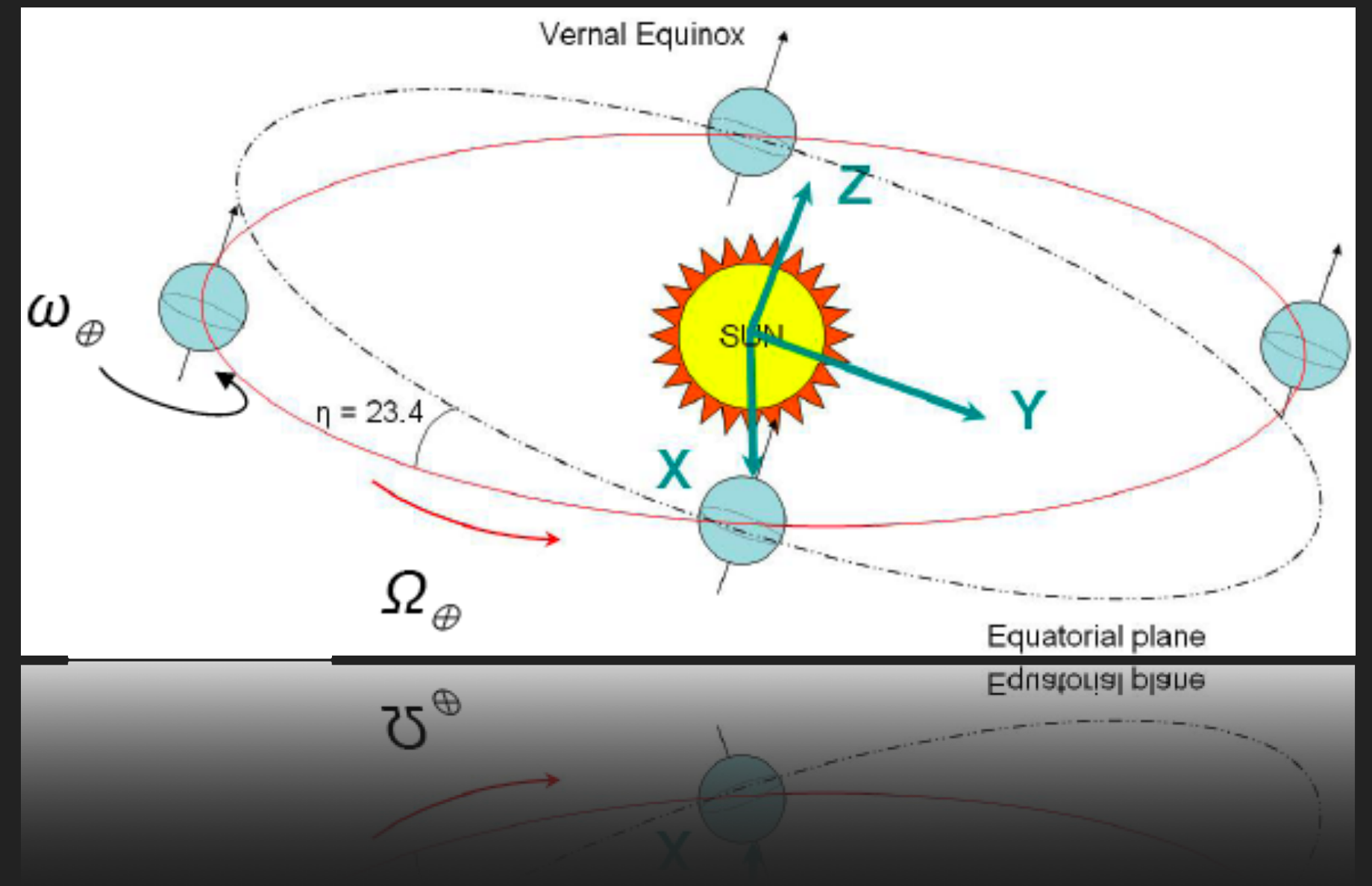
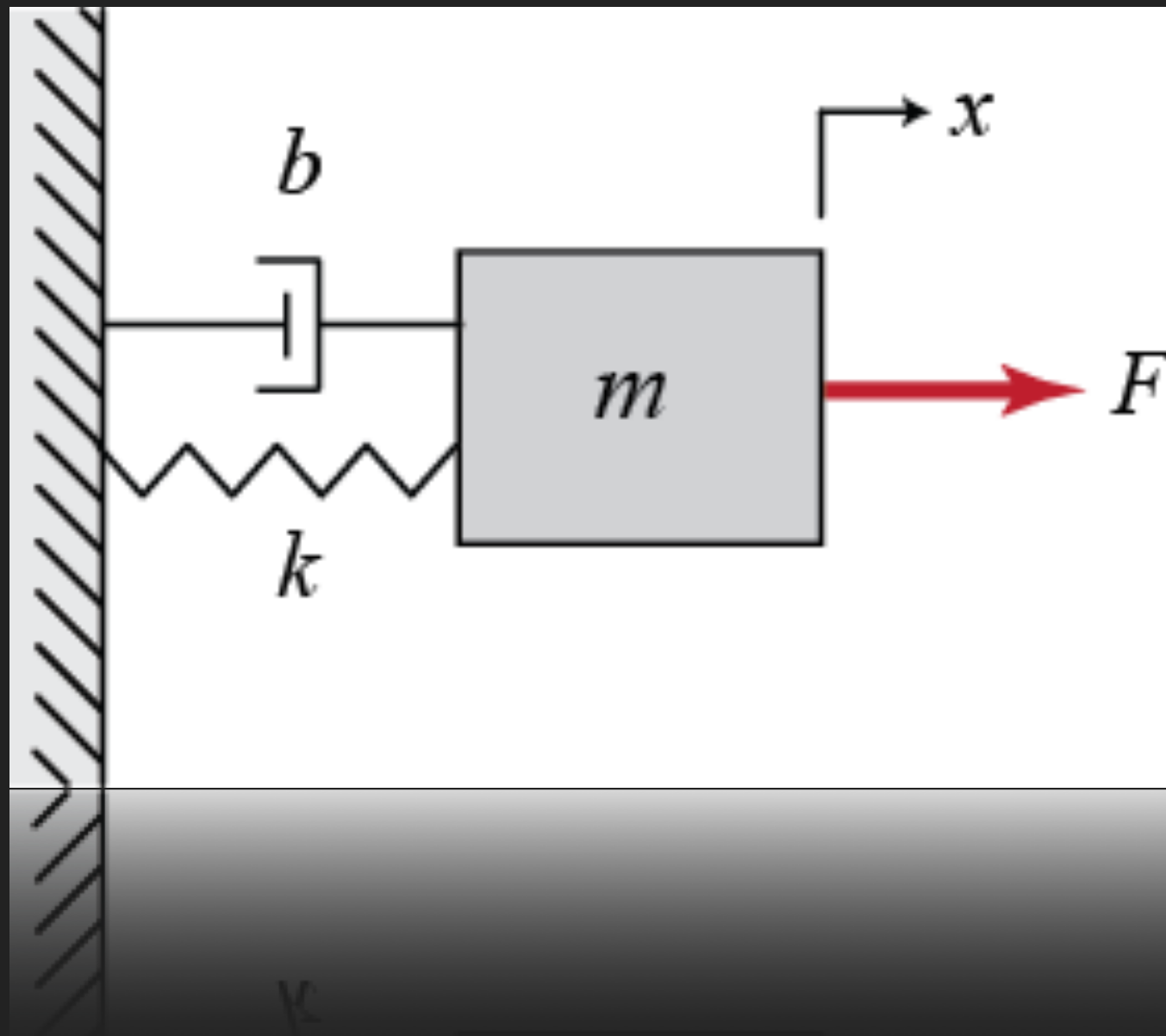
Heterodyne Detection of Axion Dark Matter in an RF Cavity

Sebastian Ellis

Phonon Sector SME Experiments at the QDM LAB-UWA



Lorentz Violations in Mass Sector



$$f(m) \rightarrow f(\hat{X}, \hat{Y}, \hat{Z})$$

$$\frac{\delta\nu}{\nu} = -\frac{1}{2} \frac{\delta m}{m} = -\frac{1}{2} \left(2c_{[xx]}^{\mathcal{Q}} + c_{00}^{\mathcal{Q}} \right)$$

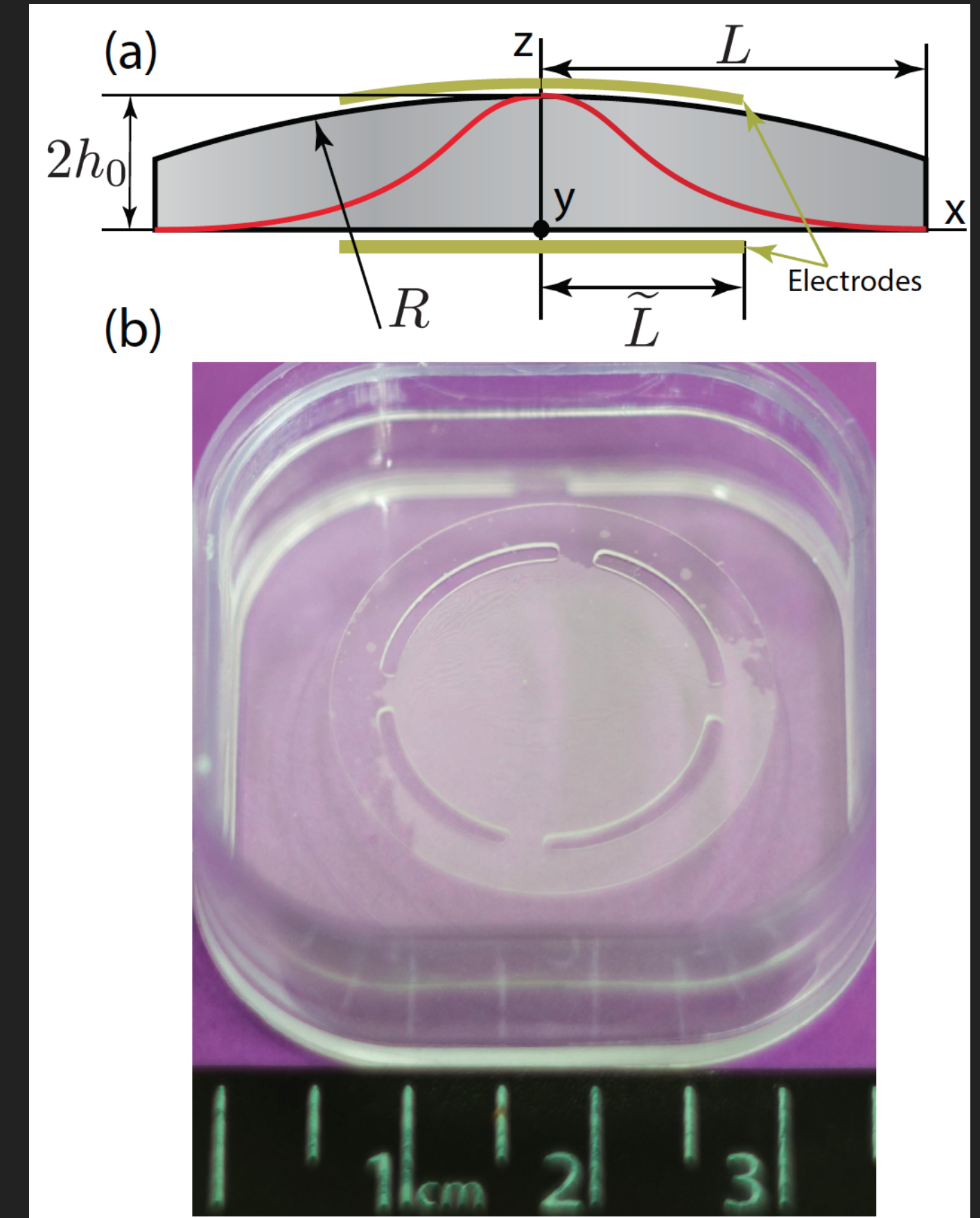
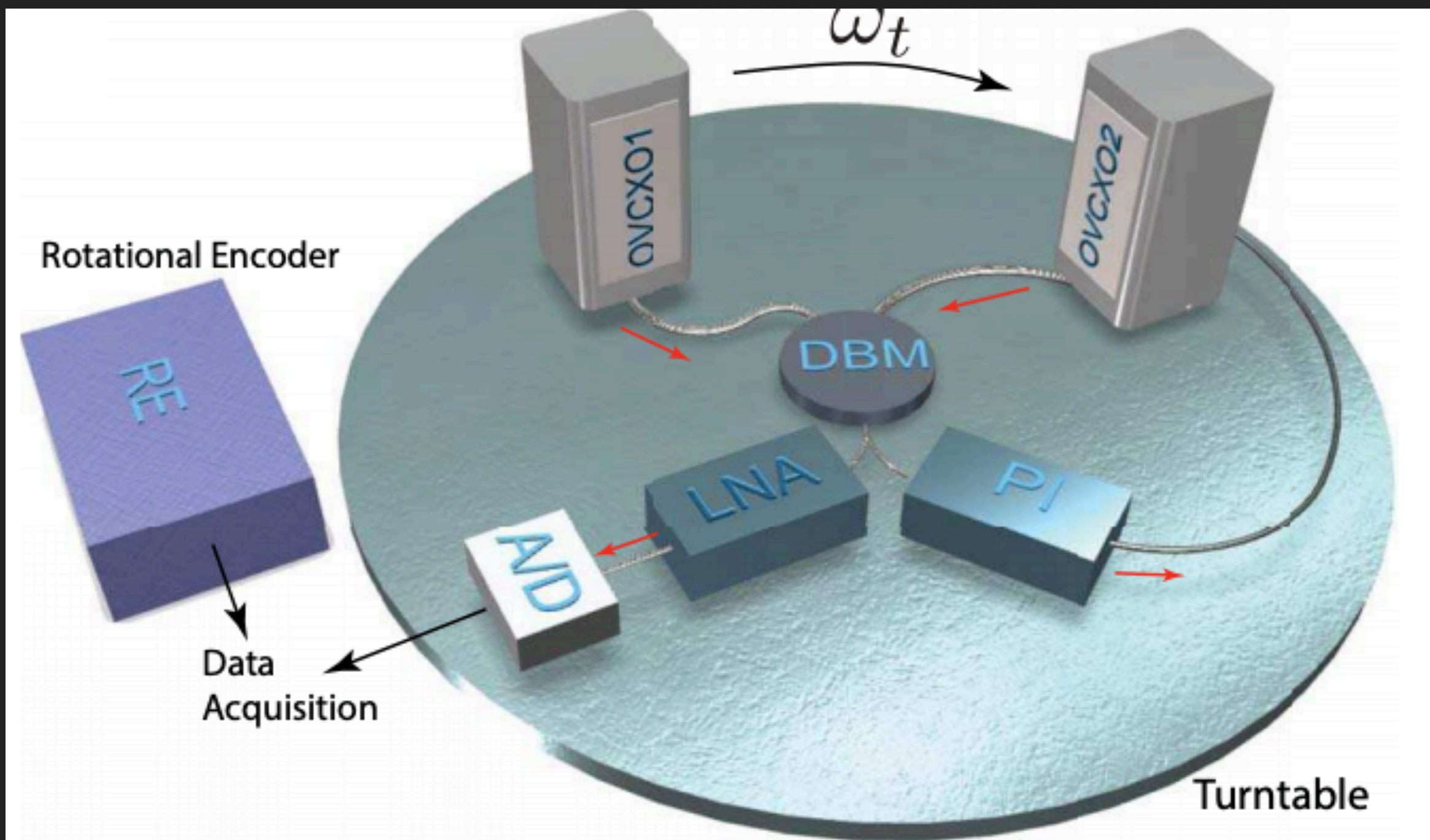
Lorentz Invariance

Acoustic Tests of Lorentz Symmetry Using Quartz Oscillators

Anthony Lo, Philipp Haslinger, Eli Mizrachi, Loïc Anderegg, Holger Müller, Michael Hohensee, Maxim Goryachev, and Michael E. Tobar

Phys. Rev. X **6**, 011018 – Published 24 February 2016

Phys. Rev. X **6**, 011018 – Published 24 February 2016



Acoustic analogue of a Fabry-Perot cavity, The most stable mechanical resonator



Lorentz Invariance

Acoustic Tests of Lorentz Symmetry Using Quartz Oscillators

Anthony Lo, Philipp Haslinger, Eli Mizrachi, Loïc Anderegg, Holger Müller, Michael Hohensee, Maxim Goryachev, and Michael E. Tobar

Phys. Rev. X **6**, 011018 – Published 24 February 2016

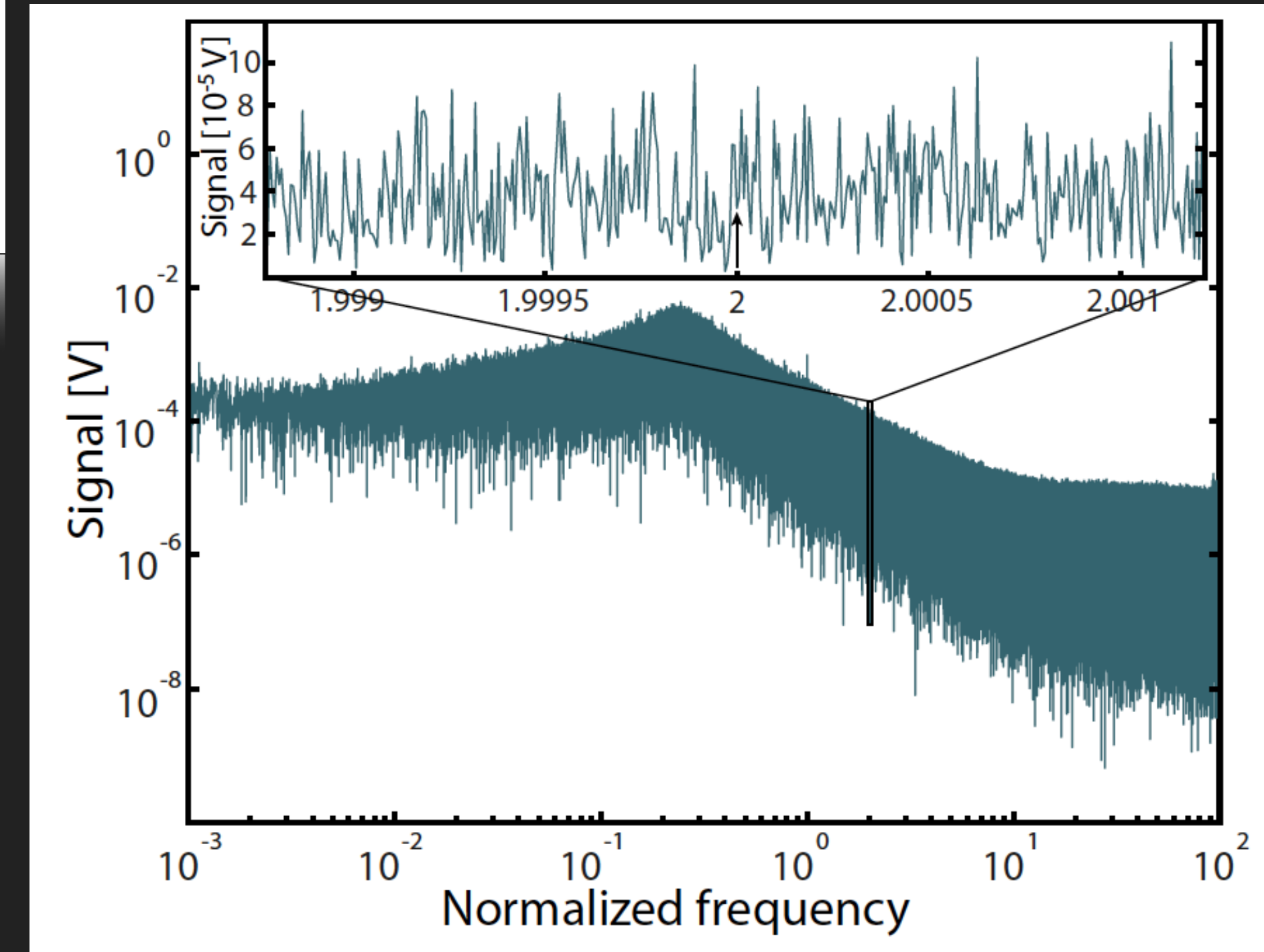
Phys. Rev. X **6**, 011018 – Published 24 February 2016

Level of Frequency Deviations:

$$\frac{\delta\nu}{\nu} = (-2 \pm 2.4) \times 10^{-15}$$

Neutron sector coefficient:

$$\tilde{c}_Q^n = (-1.8 \pm 2.2) \times 10^{-14} \text{GeV}$$



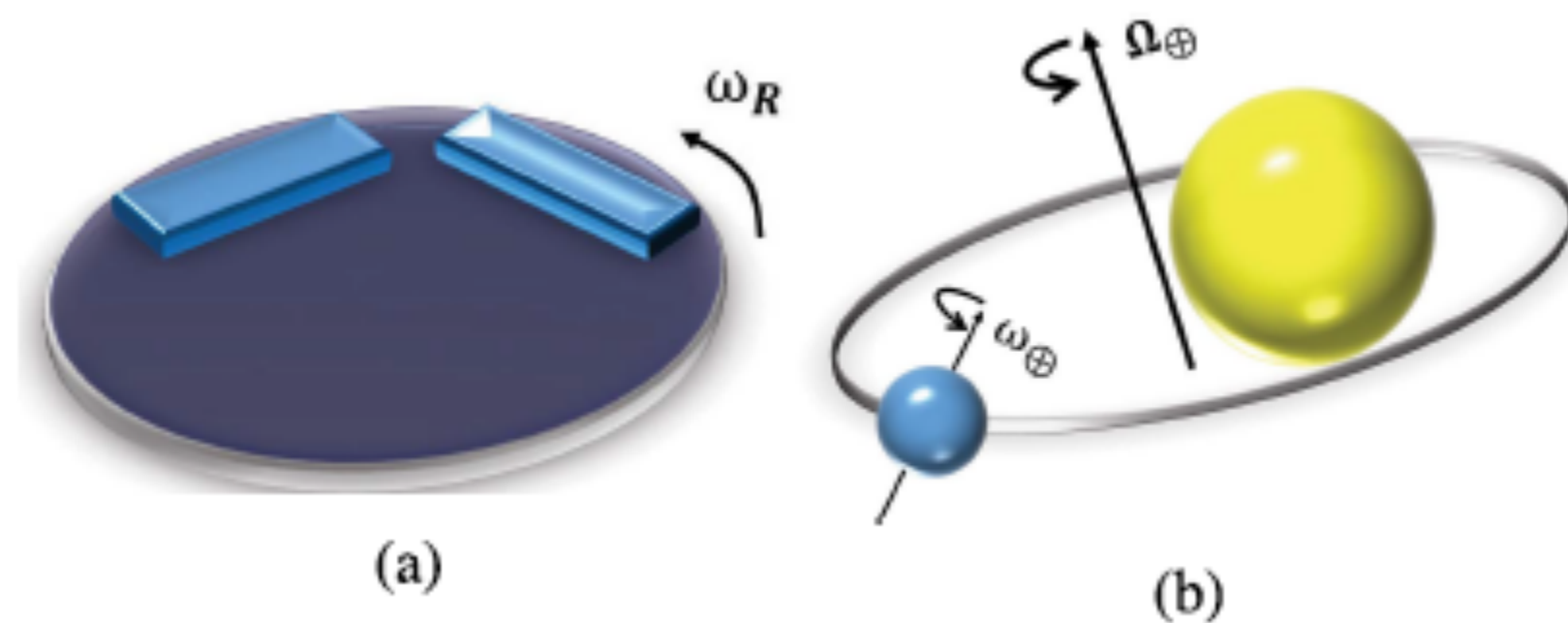
120 HOURS OF OBSERVATIONS

This rules out all possibilities for Lorentz violating anisotropies in the inertial mass of neutrons, protons, and electron at the $\sim 10^{-14}$ GeV level. A few orders of magnitude improvement over previous laboratory test and astrophysical bounds

Next Generation of Phonon Tests of Lorentz Invariance Using Quartz BAW Resonators

Maxim Goryachev, Zeyu Kuang, Eugene N. Ivanov, Philipp Haslinger, Holger Müller, and Michael E. Tobar¹, *Fellow, IEEE*

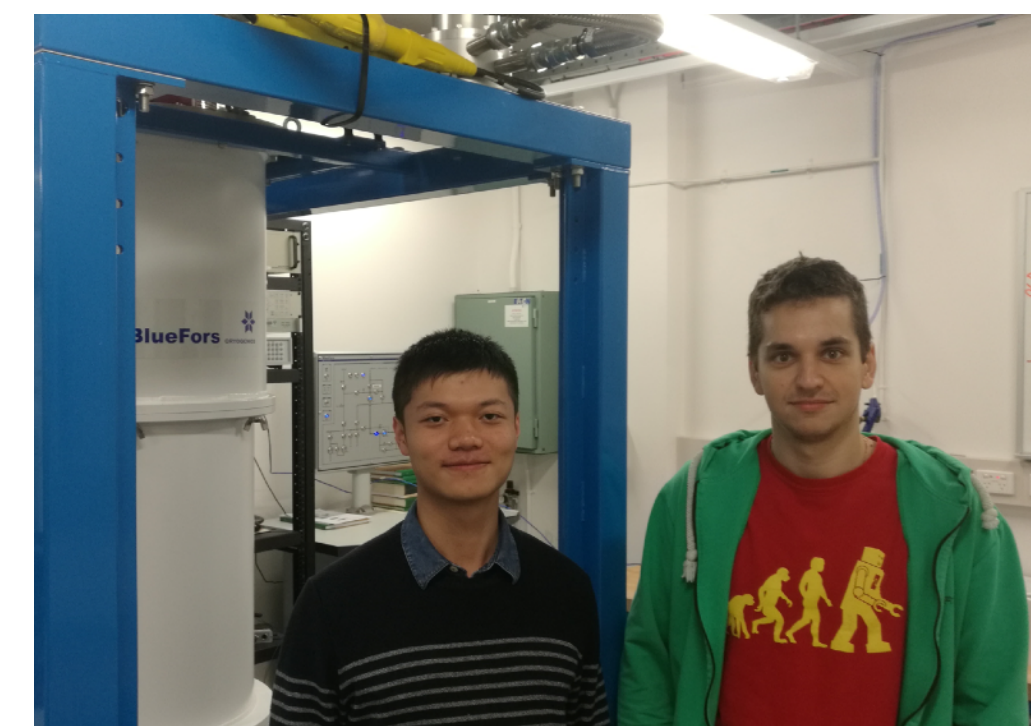
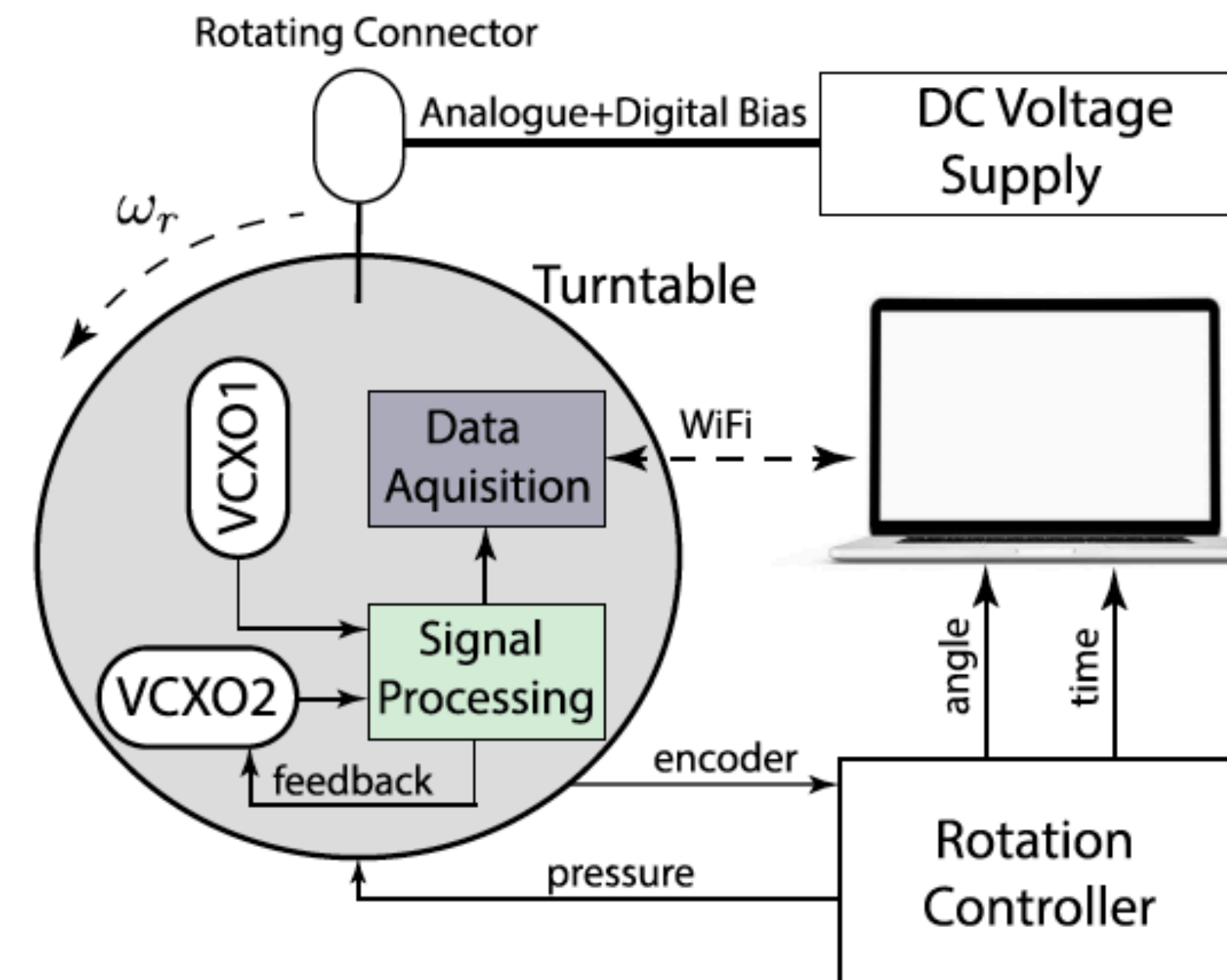
Abstract—We demonstrate technological improvements in phonon sector tests of the Lorentz invariance that implement quartz bulk acoustic wave oscillators. In this experiment, room temperature oscillators with state-of-the-art phase noise are continuously compared on a platform that rotates at a rate of order of a cycle per second. The discussion is focused on improvements in noise measurement techniques, data acquisition, and data processing. Preliminary results of the second generation of such tests are given, and indicate that standard model extension coefficients in the matter sector can be measured at a precision of order 10^{-16} GeV after taking a year's worth of data. This is equivalent to an improvement of **two orders** of magnitude over the prior acoustic phonon sector experiment.



Current Status: Data taking Finished

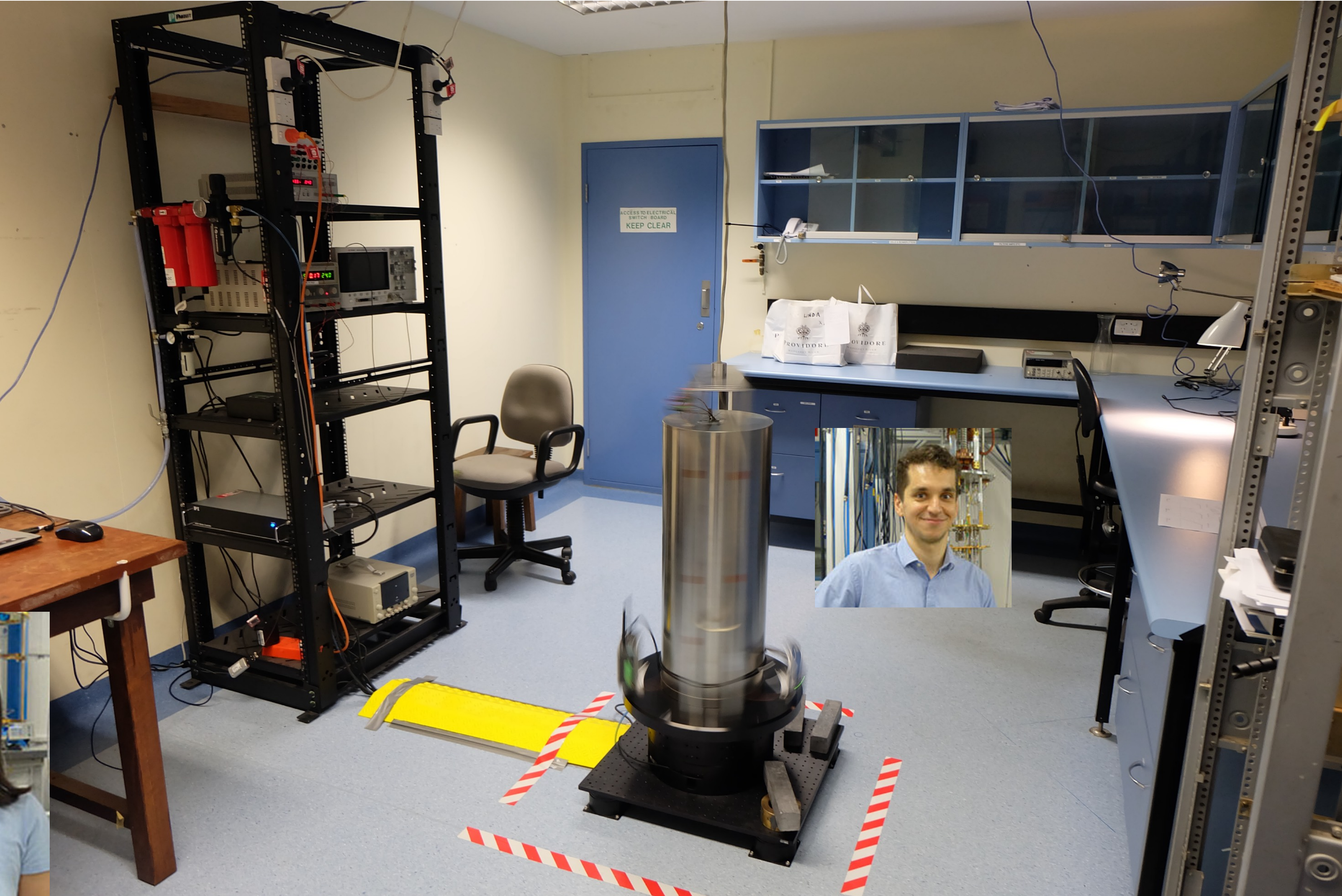
-> ~ 2 years of data

-> Multiple Coefficients, Higher Dimensions



Rotating Bulk Acoustic Wave Oscillators

Experiment Finished: Looking at Data ~2 years data (Higher Dimensions)



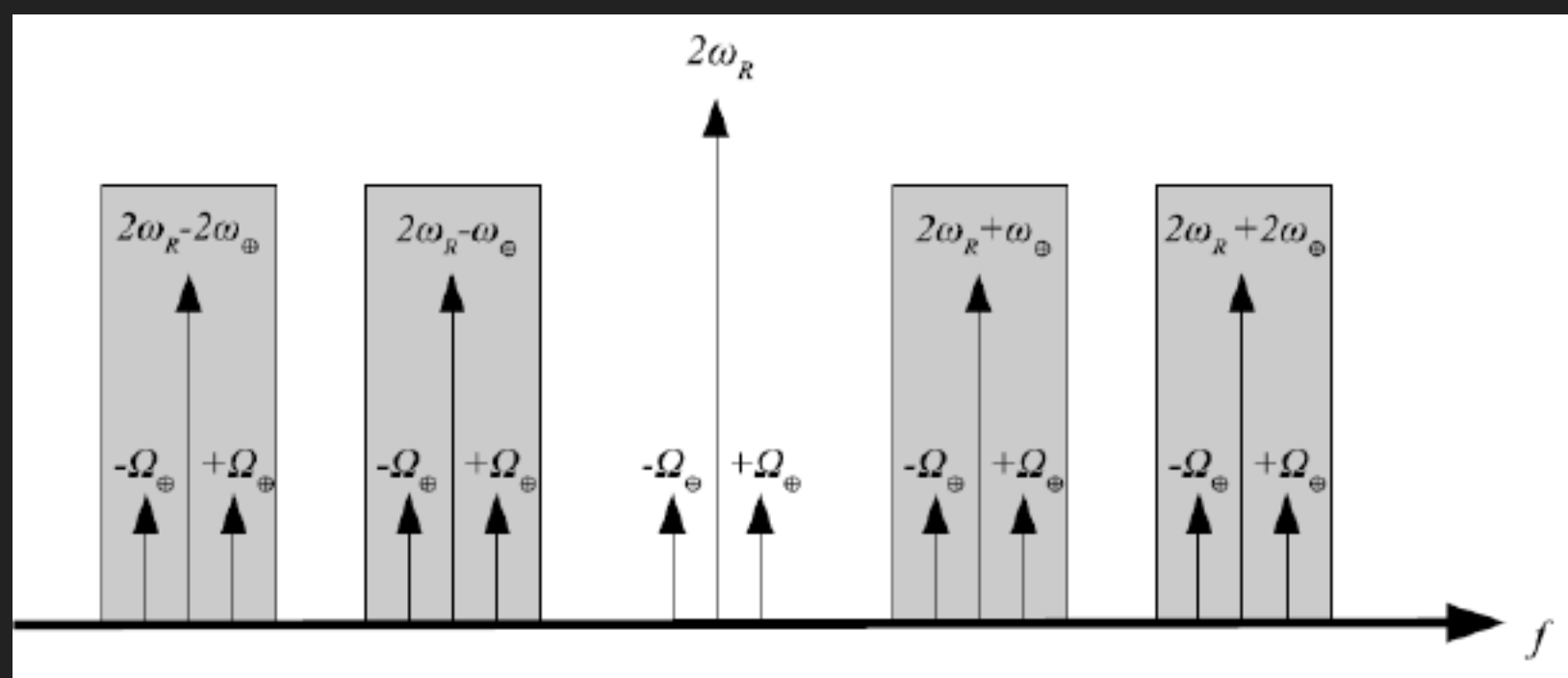
Data Spanning MJD 58211 -59047 = 2.3 Years 4/3/18 - 16/7/20

	Run 13	Run 14	Run 15	Run 16	Run 17	Run 18	stop time
Date	04/03/18	08/09/18	04/03/19	21/06/19	28/10/19	19/11/19	16/07/20
Time	6:01:18 am	10:08:12 am	1:12:30 am	8:18:28 am	03:49:21 am	1:50:29 am	
MJD	58211.25	58339.42	58576.05	58655.35	58784.16	58806.08	59046.08
delta t (s)	33092409.00	44166423.00	64611081.00	71462239.00	82591692.00	84485360.00	
t_0 (s)	33096464.88	44170478.88	64615136.88	71466294.88	82595747.88	84489415.88	
duration (days)	128.17	236.63	79.30	128.81	21.92	240.00	
used	y	y	n	n	y	y	

$$\frac{\Delta f}{f} = \frac{1}{8} (A + S(T_{\oplus}) \sin(2\omega_R T_{\oplus}) + C(T_{\oplus}) \cos(2\omega_R T_{\oplus})).$$

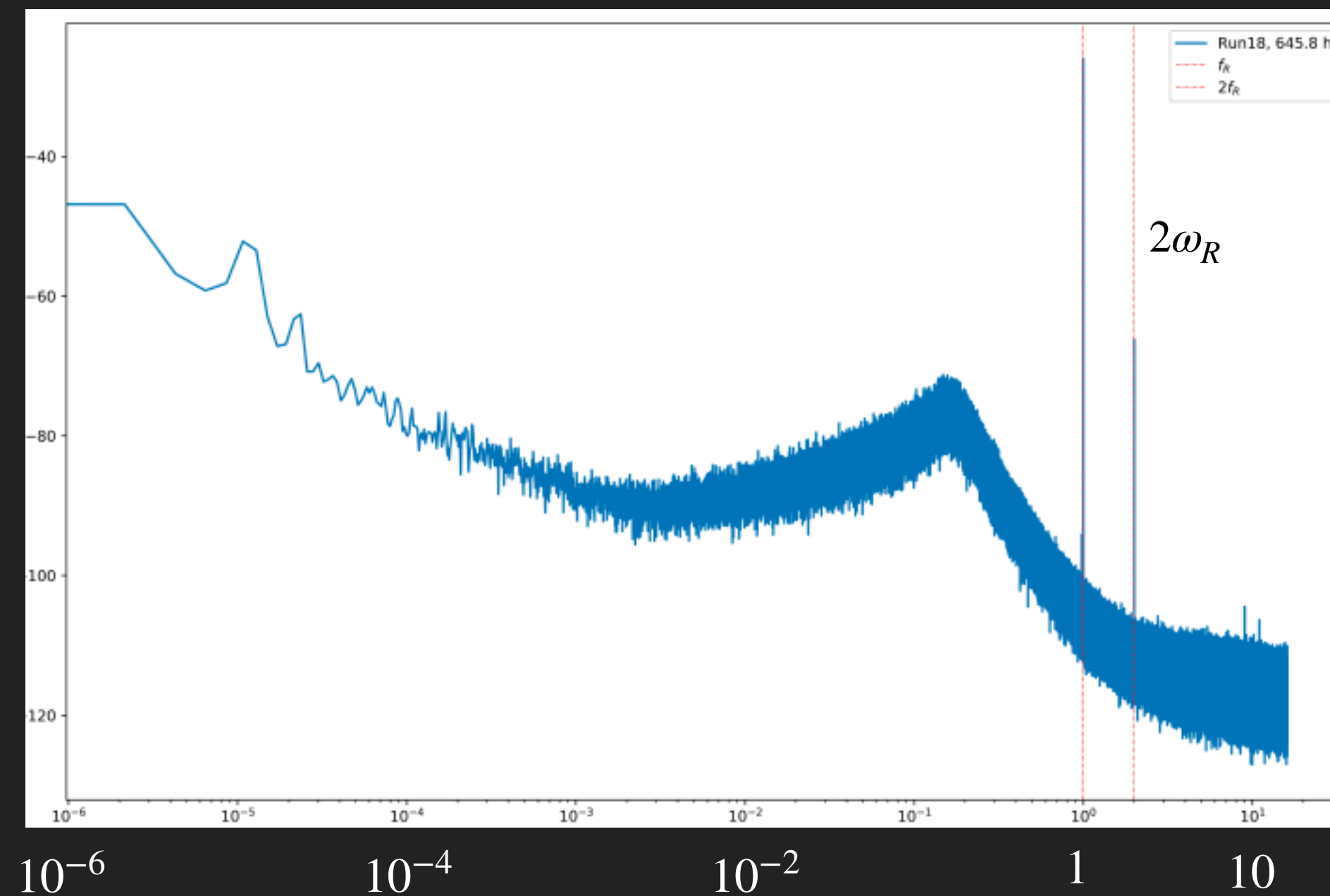
$$S(T_{\oplus}) = S_0 + \sum_i S_{s,i} \sin(\omega_i T_{\oplus}) + S_{c,i} \cos(\omega_i T_{\oplus})$$

$$C(T_{\oplus}) = C_0 + \sum_i C_{s,i} \sin(\omega_i T_{\oplus}) + C_{c,i} \cos(\omega_i T_{\oplus}).$$



ω_i (offset from $2\omega_R$)	C_{ω_i}	S_{ω_i}
DC (A)	$-4\tilde{c}_Q^T \cos(2\theta)$	0
0	$-4\sin^2 \chi \tilde{c}_Q^T$	0
$+\Omega_{\oplus}$	$2\sin^2 \chi (\cos \eta \tilde{c}_{TY}^T - 2\tilde{c}_{TZ}^T \sin \eta) \beta_{\oplus}$	$-2\sin^2 \chi \tilde{c}_{TX}^T \beta_{\oplus}$
$-\Omega_{\oplus}$	$2\sin^2 \chi (\cos \eta \tilde{c}_{TY}^T - 2\tilde{c}_{TZ}^T \sin \eta) \beta_{\oplus}$	$2\sin^2 \chi \tilde{c}_{TX}^T \beta_{\oplus}$
$+\omega_{\oplus} - \Omega_{\oplus}$	$2(1 + \cos \chi) \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_{\oplus}$	$2(1 + \cos \chi) \sin \chi \beta_{\oplus} \times$ $(\tilde{c}_{TZ}^T + \cos \eta \tilde{c}_{TZ}^T + \tilde{c}_{TY}^T \sin \eta)$
$-\omega_{\oplus} + \Omega_{\oplus}$	$2(\cos \chi - 1) \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_{\oplus}$	$2(\cos \chi - 1) \sin \chi \beta_{\oplus} \times$ $(\tilde{c}_{TZ}^T + \cos \eta \tilde{c}_{TZ}^T + \tilde{c}_{TY}^T \sin \eta)$
$+\omega_{\oplus}$	$-4(1 + \cos \chi) \tilde{c}_Y^T \sin \chi$	$-4(1 + \cos \chi) \tilde{c}_X^T \sin \chi$
$-\omega_{\oplus}$	$4(\cos \chi - 1) \tilde{c}_Y^T \sin \chi$	$4(\cos \chi - 1) \tilde{c}_X^T \sin \chi$
$+\omega_{\oplus} + \Omega_{\oplus}$	$2(1 + \cos \chi) \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_{\oplus}$	$2(1 + \cos \chi) \sin \chi \beta_{\oplus} \times$ $[(-1 + \cos \eta) \tilde{c}_{TZ}^T + \sin \eta \tilde{c}_{TY}^T]$
$-\omega_{\oplus} - \Omega_{\oplus}$	$2(\cos \chi - 1) \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_{\oplus}$	$2(1 - \cos \chi) \sin \chi \beta_{\oplus} \times$ $[(-1 + \cos \eta) \tilde{c}_{TZ}^T + \sin \eta \tilde{c}_{TY}^T]$
$+2\omega_{\oplus} - \Omega_{\oplus}$	$(1 + \cos \eta)(1 + \cos \chi)^2 \tilde{c}_{TY}^T \beta_{\oplus}$	$-(1 + \cos \eta)(1 + \cos \chi)^2 \tilde{c}_{TX}^T \beta_{\oplus}$
$-2\omega_{\oplus} + \Omega_{\oplus}$	$(1 + \cos \eta)(-1 + \cos \chi)^2 \tilde{c}_{TY}^T \beta_{\oplus}$	$(1 + \cos \eta)(-1 + \cos \chi)^2 \tilde{c}_{TX}^T \beta_{\oplus}$
$+2\omega_{\oplus}$	$2\tilde{c}_-^T (1 + \cos \chi)^2$	$2(1 + \cos \chi)^2 \tilde{c}_Z^T$
$-2\omega_{\oplus}$	$2(\cos \chi - 1)^2 \tilde{c}_-^T$	$-2(\cos \chi - 1)^2 \tilde{c}_Z^T$
$+2\omega_{\oplus} + \Omega_{\oplus}$	$(\cos \eta - 1)(1 + \cos \chi)^2 \tilde{c}_{TY}^T \beta_{\oplus}$	$(1 - \cos \eta)(1 + \cos \chi)^2 \tilde{c}_{TX}^T \beta_{\oplus}$
$-2\omega_{\oplus} - \Omega_{\oplus}$	$(\cos \eta - 1)(-1 + \cos \chi)^2 \tilde{c}_{TY}^T \beta_{\oplus}$	$(\cos \eta - 1)(-1 + \cos \chi)^2 \tilde{c}_{TX}^T \beta_{\oplus}$

dBV/√Hz

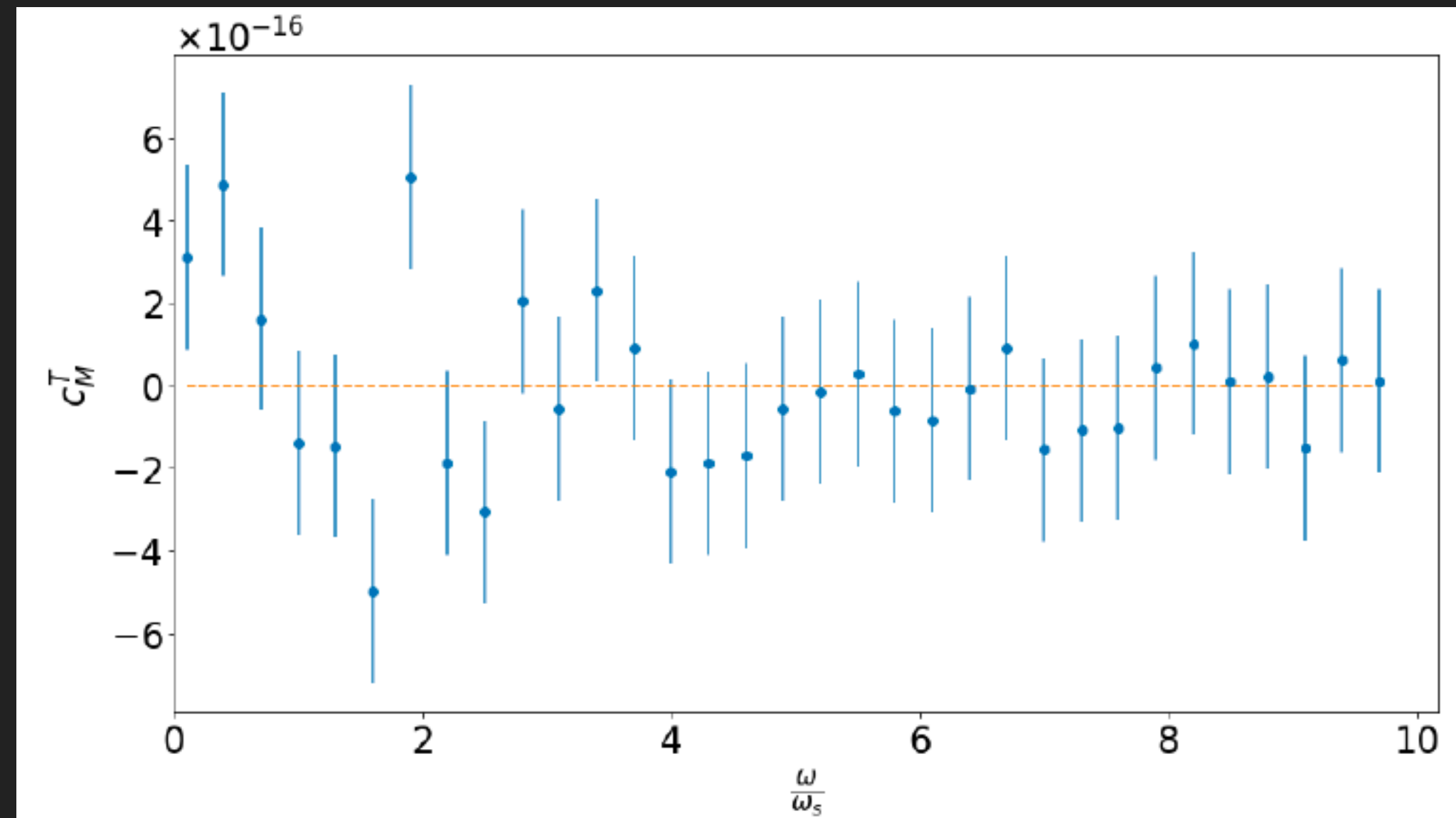
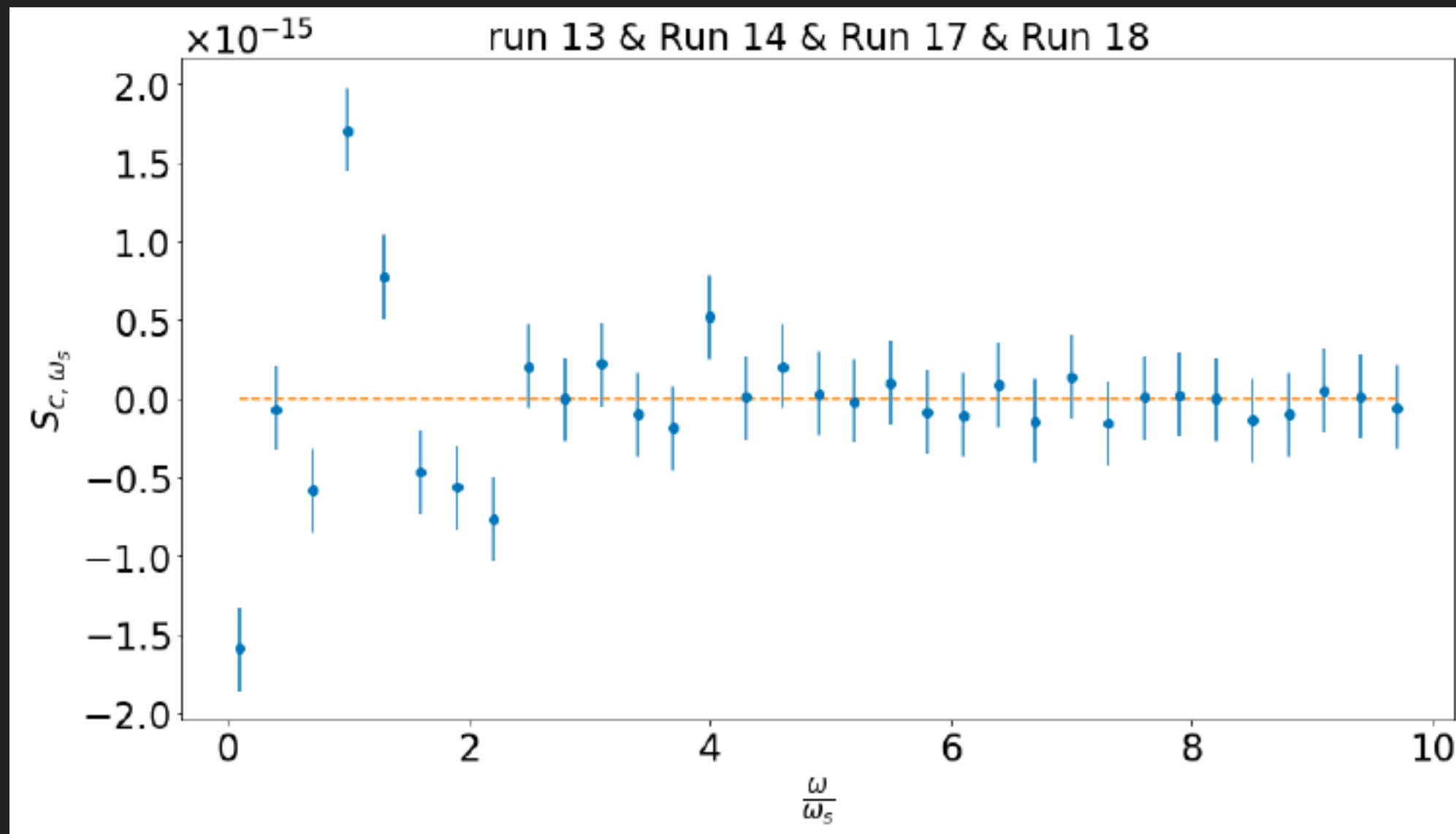
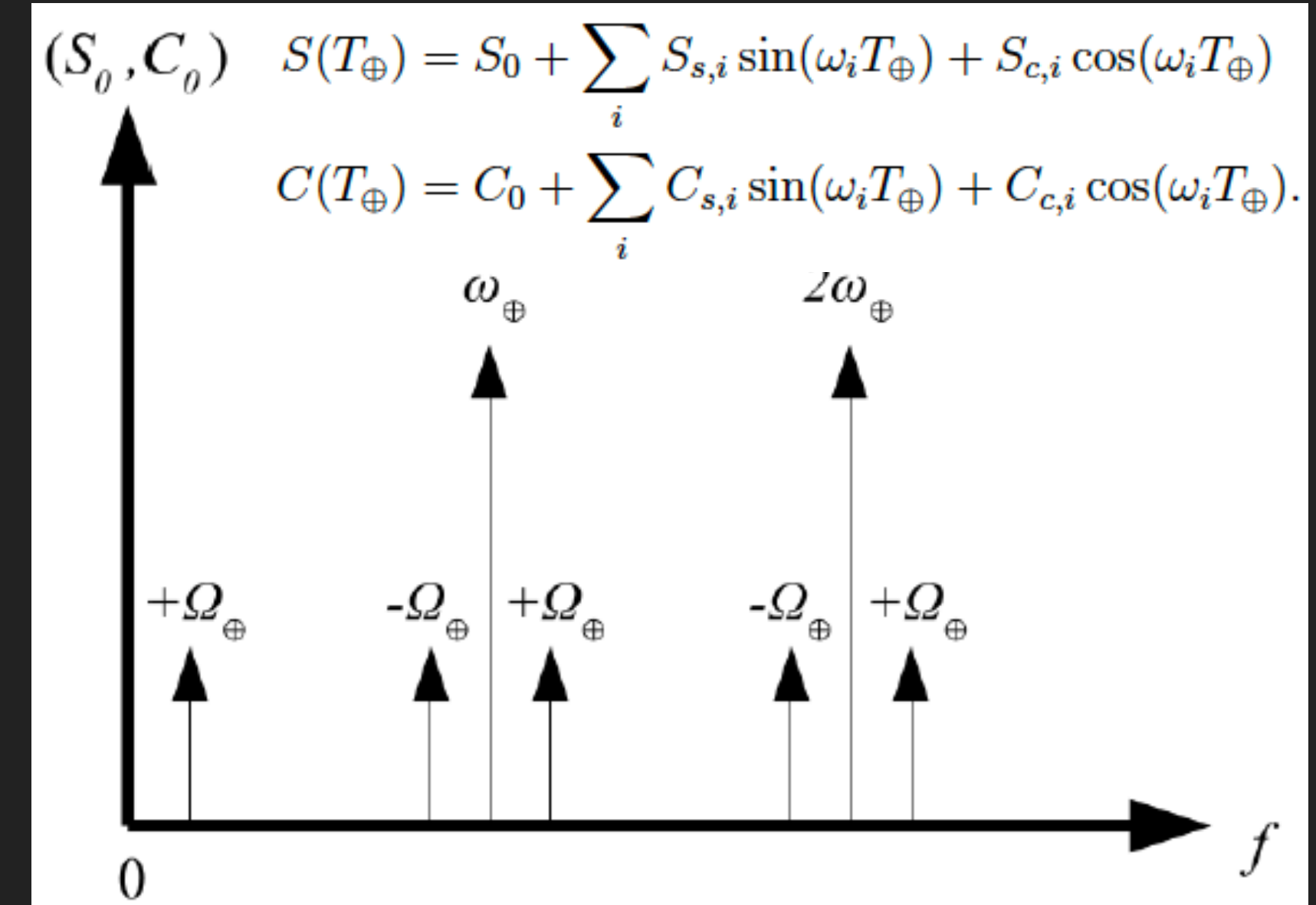


Fourier Frequency Hz

Demodulate, by averaging over an optimum number of rotations


ω_i	C_{C,ω_i}	C_{S,ω_i}
0	$-4 \sin^2 \chi \tilde{c}_Q^T$	0
Ω_\oplus	$4 \sin^2 \chi (\cos \eta \tilde{c}_{TY}^T - 2 \tilde{c}_{TZ}^T \sin \eta) \beta_\oplus$	$-4 \sin^2 \chi \tilde{c}_{TX}^T \beta_\oplus$
$\omega_\oplus - \Omega_\oplus$	$4 \cos \chi \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_\oplus$	$4 \cos \chi (\tilde{c}_{TZ}^T + \cos \eta \tilde{c}_{TZ}^T + \tilde{c}_{TY}^T \sin \eta) \times \sin \chi \beta_\oplus$
ω_\oplus	$-8 \tilde{c}_Y^T \sin \chi$	$-8 \cos \chi \tilde{c}_X^T \sin \chi$
$\omega_\oplus + \Omega_\oplus$	$4 \cos \chi \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_\oplus$	$4 [(-1 + \cos \eta) \tilde{c}_{TZ}^T + \cos \chi \tilde{c}_{TY}^T \sin \eta] \times \sin \chi \beta_\oplus$
$2\omega_\oplus - \Omega_\oplus$	$2(1 + \cos \eta)(1 + \cos^2 \chi) \tilde{c}_{TY}^T \beta_\oplus$	$-2(1 + \cos^2 \chi)(1 + \cos \eta) \tilde{c}_{TX}^T \beta_\oplus$
$2\omega_\oplus$	$4 \tilde{c}_-^T (\cos^2 \chi + 1)$	$4(1 + \cos^2 \chi) \tilde{c}_Z^T$
$2\omega_\oplus + \Omega_\oplus$	$2(1 + \cos^2 \chi)(\cos \eta - 1) \tilde{c}_{TY}^T \beta_\oplus$	$2(1 - \cos \eta)(1 + \cos^2 \chi) \tilde{c}_{TX}^T \beta_\oplus$

ω_i	S_{C,ω_i}	S_{S,ω_i}
0	0	0
Ω_\oplus	0	0
$\omega_\oplus - \Omega_\oplus$	$4(\tilde{c}_{TZ}^T + \cos \eta \tilde{c}_{TZ}^T + \tilde{c}_{TY}^T \sin \eta) \sin \chi \beta_\oplus$	$-4 \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_\oplus$
ω_\oplus	$-8 \tilde{c}_X^T \sin \chi$	$8 \cos \chi \tilde{c}_Y^T \sin \chi$
$\omega_\oplus + \Omega_\oplus$	$4[\cos \chi (-1 + \cos \eta) \tilde{c}_{TZ}^T + \tilde{c}_{TY}^T \sin \eta] \sin \chi \beta_\oplus$	$-4 \tilde{c}_{TX}^T \sin \eta \sin \chi \beta_\oplus$
$2\omega_\oplus - \Omega_\oplus$	$-4(1 + \cos \eta) \cos \chi \tilde{c}_{TX}^T \beta_\oplus$	$-4(1 + \cos \eta) \cos \chi \tilde{c}_{TY}^T \beta_\oplus$
$2\omega_\oplus$	$8 \tilde{c}_Z^T \cos \chi$	$-8 \tilde{c}_-^T \cos \chi$
$2\omega_\oplus + \Omega_\oplus$	$4(1 - \cos \eta) \cos \chi \tilde{c}_{TX}^T \beta_\oplus$	$-4 \cos \chi (\cos \eta - 1) \tilde{c}_{TY}^T \beta_\oplus$



Article

Non-Minimal Lorentz Violation in Macroscopic Matter

Matthew Mewes 

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Abstract: The effects of Lorentz and CPT violations on macroscopic objects are explored. Effective composite coefficients for Lorentz violation are derived in terms of coefficients for electrons, protons, and neutrons in the Standard-Model Extension, including all minimal and non-minimal violations. The hamiltonian and modified Newton's second law for a test body are derived. The framework is applied to free-fall and torsion-balance tests of the weak equivalence principle and to orbital motion. The effects on continuous media are studied, and the frequency shifts in acoustic resonators are calculated.

3.3. Acoustic Resonators

$$\frac{\delta\omega}{\omega_0} \approx -\frac{1}{2} \frac{\int_V d^3x (u_0^{[-1]} \cdot C^{[0]} \cdot u_0^{[1]} + u_0^{[1]} \cdot C^{[-2]} \cdot u_0^{[1]})}{\int_V d^3x u_0^{[1]} \cdot u_0^{[-1]}}$$

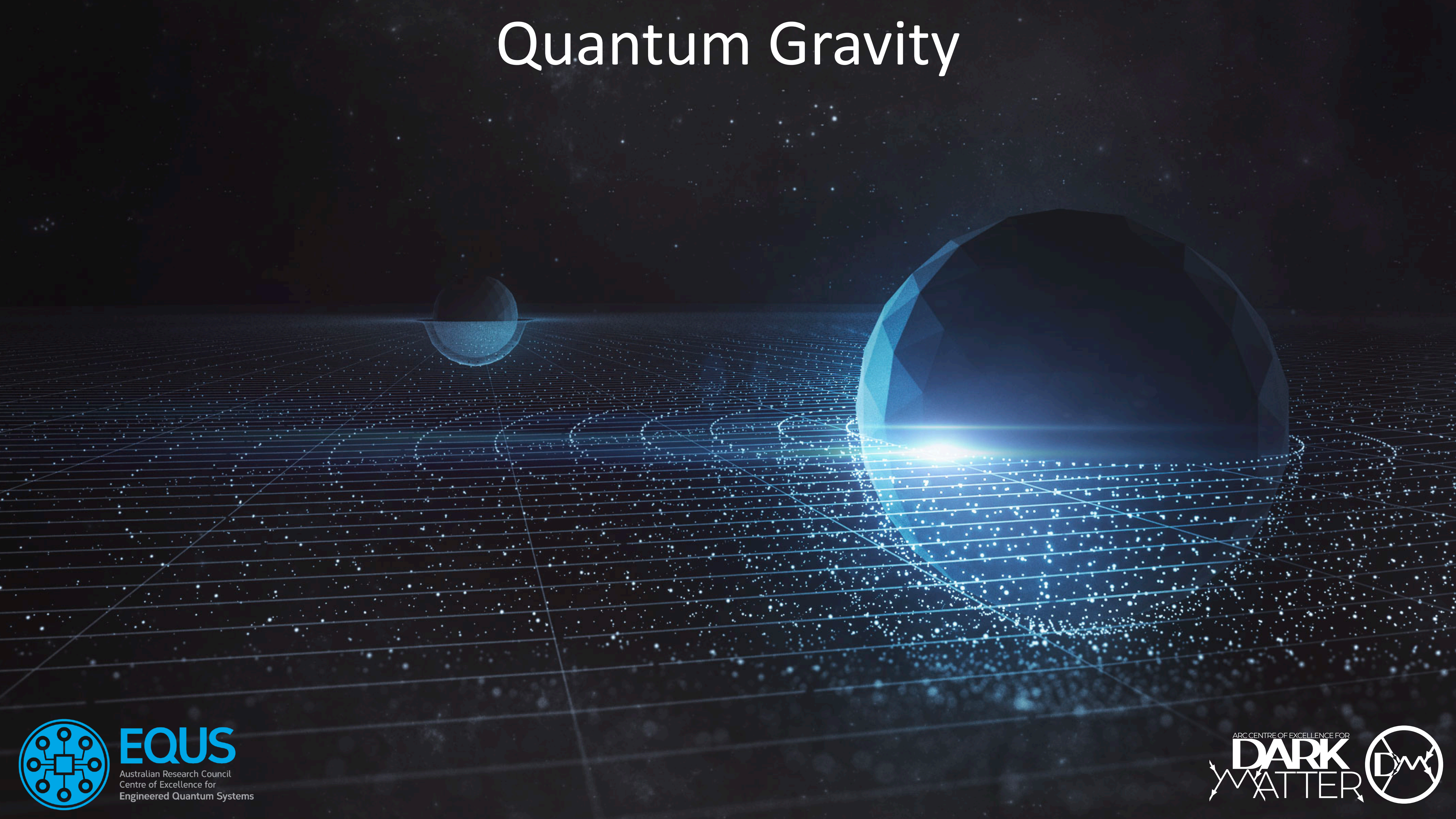
$$= -\frac{1}{2} \frac{\int_V d^3x \langle u_0 \cdot C \cdot u_0 \rangle_t}{\int_V d^3x \langle u_0 \cdot u_0 \rangle_t},$$

$$\frac{\delta\omega}{\omega_0} = \sum_{m_r m_s m_a} A_{m_r m_s m_a} e^{im_r \varphi + im_s \omega_{\oplus} T_{\oplus} + im_a \Omega_{\oplus} T},$$



Searches for Lorentz violation in quartz resonators have demonstrated sensitivities on the order of parts in 10^{14} to $d = 4$ violations [63], and are expected to improve by two orders of magnitude [103]. We therefore expect sensitivities near $10^{-16} \text{ GeV}^{4-d}$ to the dimension- d combinations $a'_{npe2jm}^{(d)}$ and $c'_{npe2jm}^{(d)}$.

Quantum Gravity



Quantum Gravity

Testing the generalized uncertainty principle with macroscopic mechanical oscillators and pendulums

P. A. Bushev, J. Bourhill, M. Goryachev, N. Kukharchyk, E. Ivanov, S. Galliou, M. E. Tobar, and S. Danilishin
Phys. Rev. D **100**, 066020 – Published 20 September 2019

Recent progress in observing and manipulating mechanical oscillators at quantum regime provides new opportunities of studying fundamental physics, for example to search for low energy signatures of quantum gravity. For example, it was recently proposed that such devices can be used to test quantum gravity effects, by detecting the change in the $[\hat{x}, \hat{p}]$ commutation relation that could result from quantum gravity corrections. We show that such a correction results in a dependence of a resonant frequency of a mechanical oscillator on its amplitude, which is known as the amplitude-frequency effect. By implementing this new method we measure the amplitude-frequency effect for a 0.3 kg ultra-high-Q sapphire split-bar mechanical resonator and for an $\sim 10^{-5}$ kg quartz bulk acoustic wave resonator. Our experiments with a sapphire resonator have established the upper limit on a quantum gravity correction constant of β_0 to not exceed 5.2×10^6 , which is a factor of 6 better than previously measured. The reasonable estimates of β_0 from experiments with quartz resonators yields $\beta_0 < 4 \times 10^4$. The datasets of 1936 measurements of a physical pendulum period by Atkinson [E. C. Atkinson, *Proc. Phys. Soc. London* **48**, 606 (1936).] could potentially lead to significantly stronger limitations on $\beta_0 \ll 1$. Yet, due to the lack of proper pendulum frequency stability measurement in these experiments the exact upper bound on β_0 cannot be reliably established. Moreover, pendulum based systems only allow one to test a specific form of the modified commutator that depends on the mean value of momentum. The electromechanical oscillators to the contrary enable testing of any form of generalized uncertainty principle directly due to a much higher stability and a higher degree of control.

control

form of generalized uncertainty principle directly due to a much higher stability and a higher degree of mean value of momentum. The electromechanical oscillators to the contrary enable testing of any

Quantum Gravity

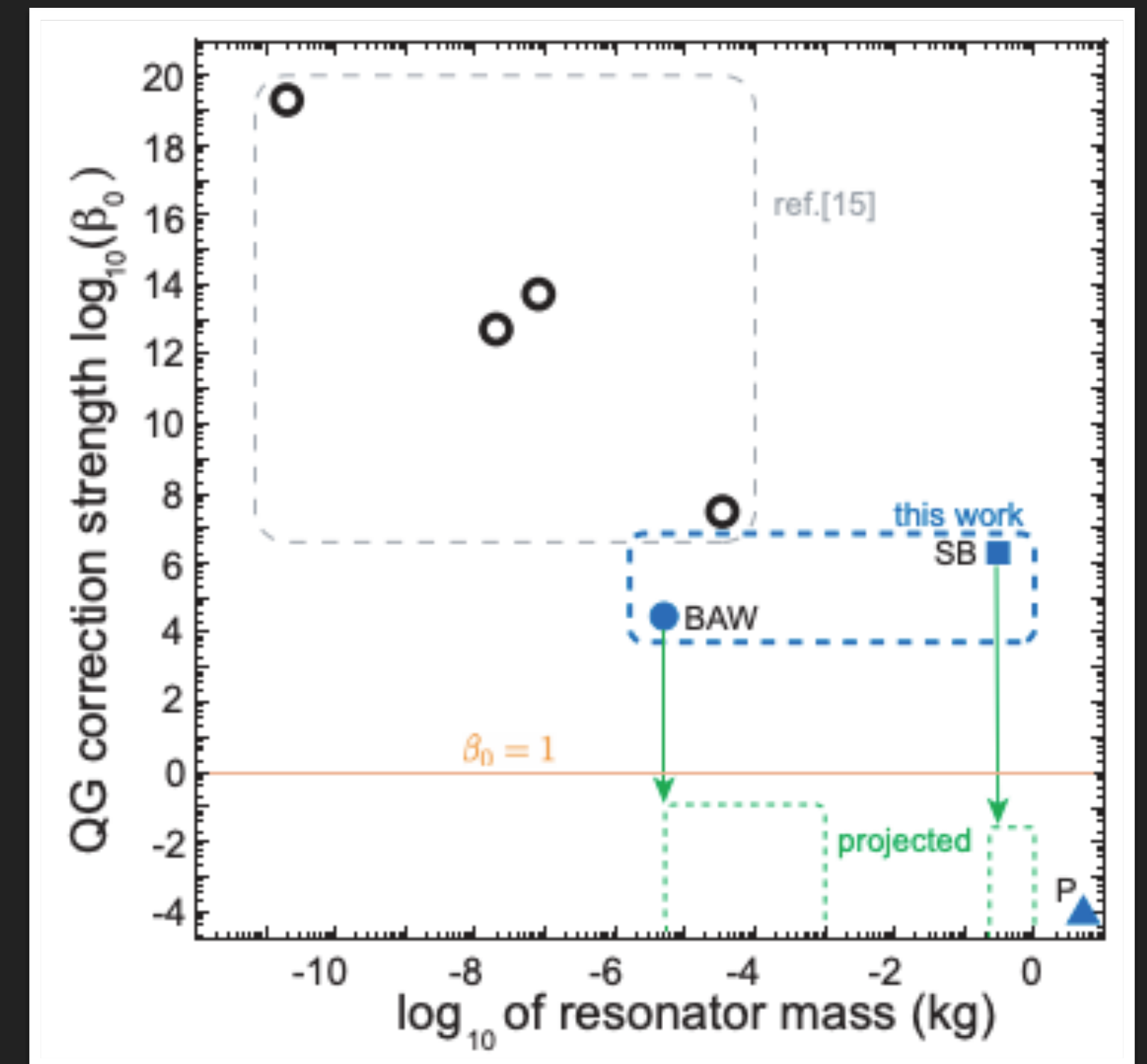
Testing the generalized uncertainty principle with macroscopic mechanical oscillators and pendulums

P. A. Bushev, J. Bourhill, M. Goryachev, N. Kukharchyk, E. Ivanov, S. Galliou, M. E. Tobar, and S. Danilishin
Phys. Rev. D **100**, 066020 – Published 20 September 2019

$$[\hat{x}, \hat{p}]_{\beta_0} = i\hbar \left[1 + \beta_0 \left(\frac{\hat{p}}{M_p c} \right)^2 \right],$$

$$\Delta x \Delta p \geq \frac{\hbar}{2} \left[1 + \beta_0 \frac{\Delta p^2 + \langle p \rangle^2}{M_p^2 c^2} \right]$$

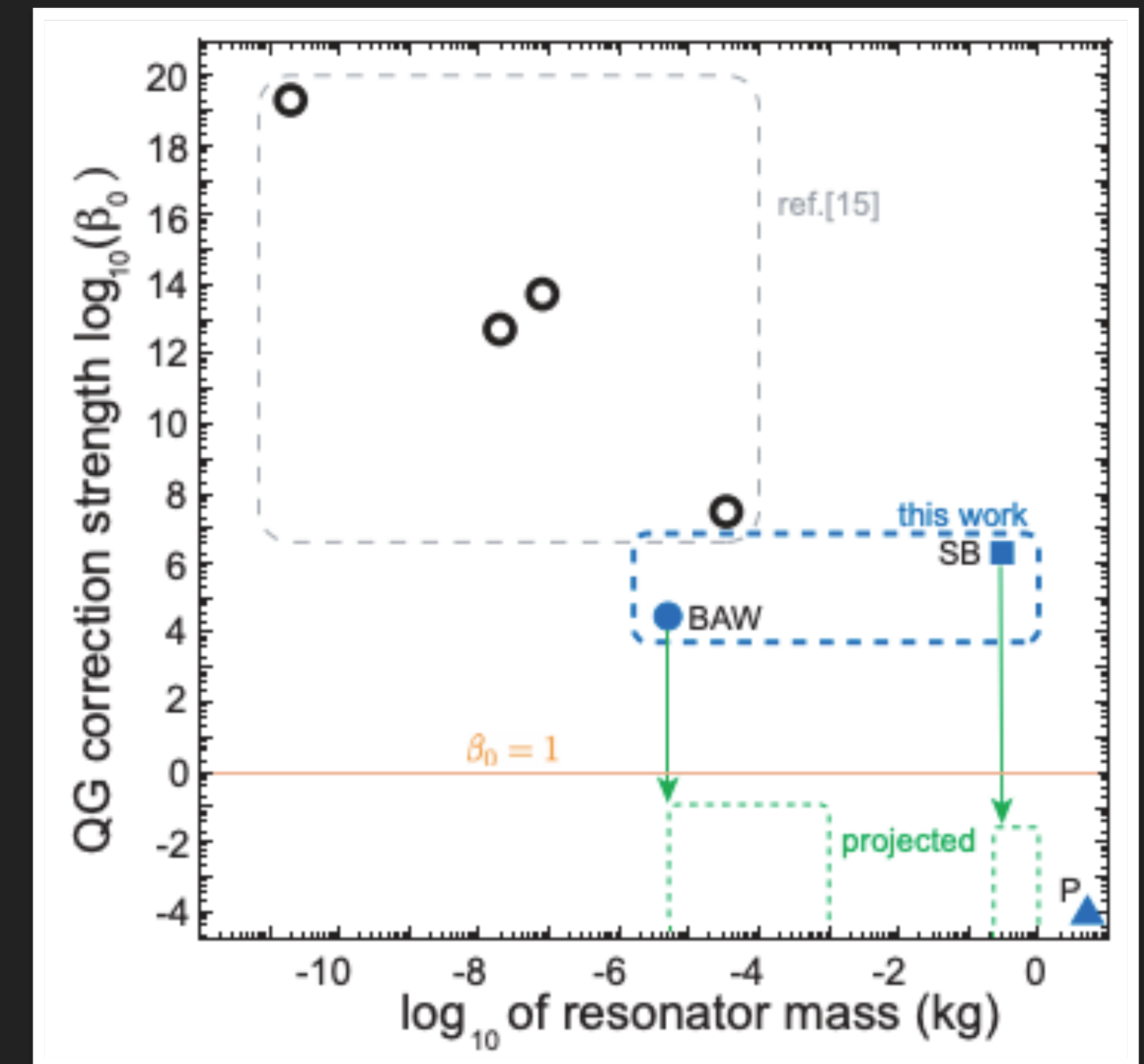
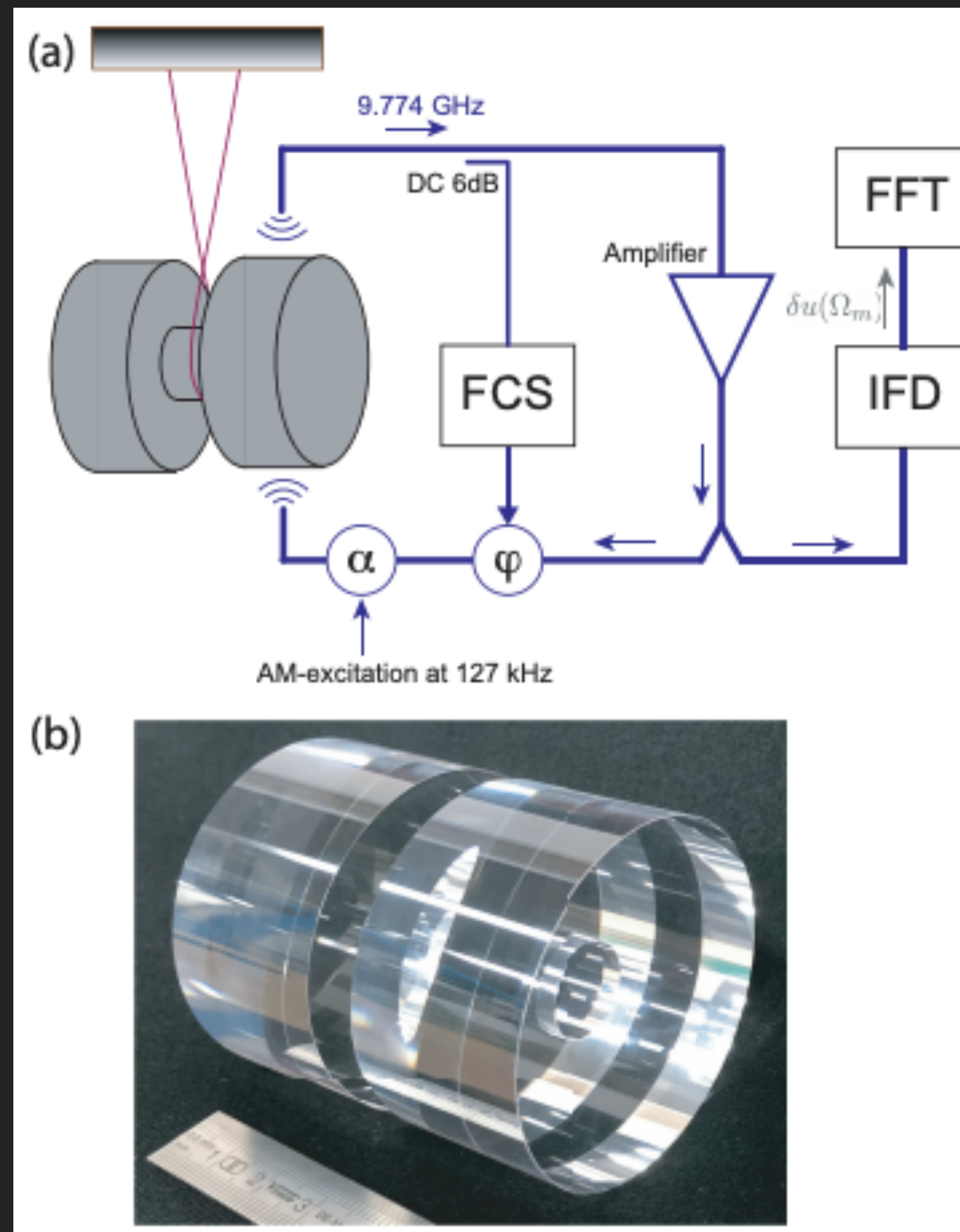
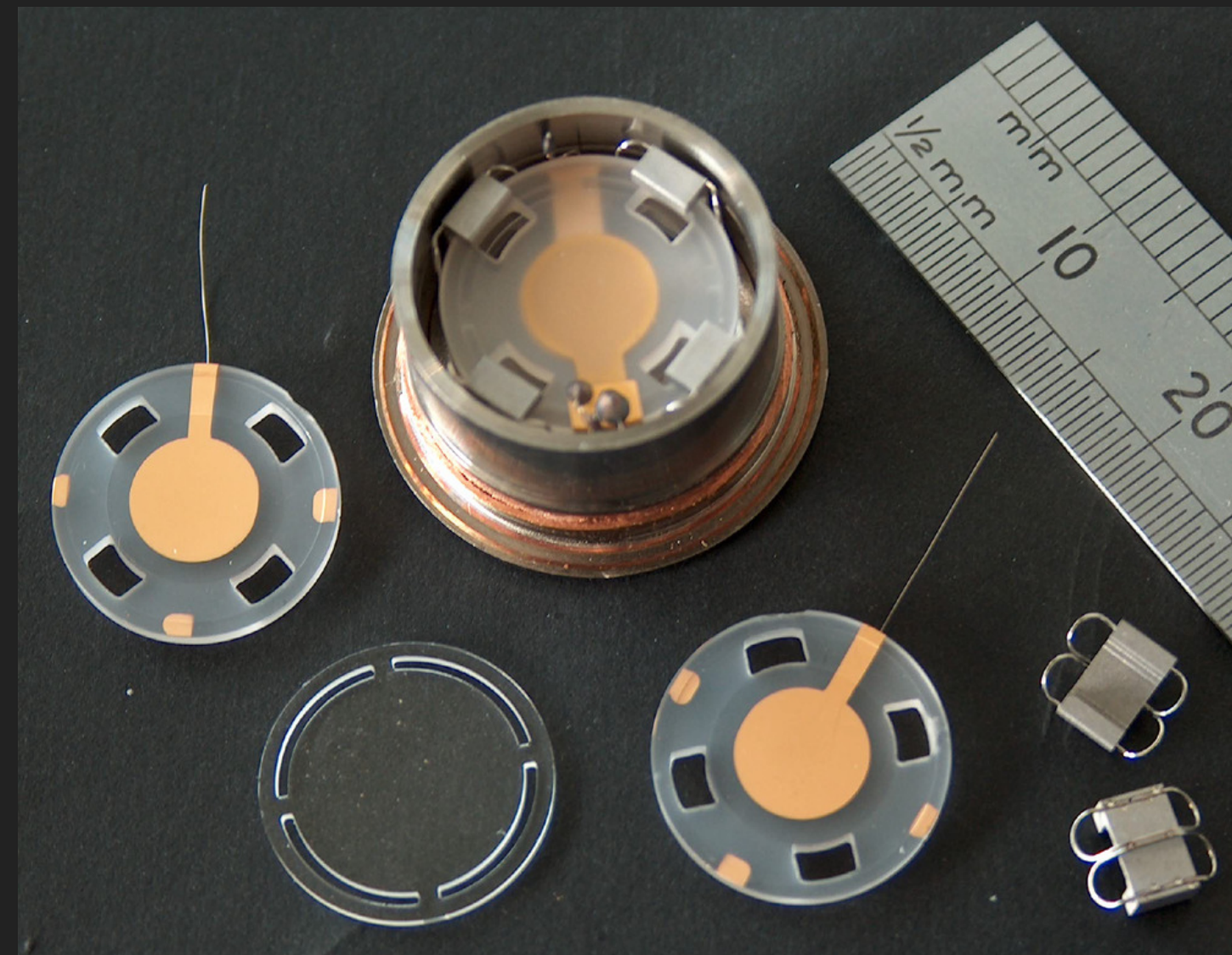
$$\hat{H} \rightarrow \hat{H}_0 + \Delta\hat{H} = \left(\hat{p}^2 / 2m + m\Omega_0^2 \hat{x}^2 / 2 \right) + \beta_0 \hat{p}^4 / (3m(M_p c)^2)$$



Quantum Gravity

Testing the generalized uncertainty principle with macroscopic mechanical oscillators and pendulums

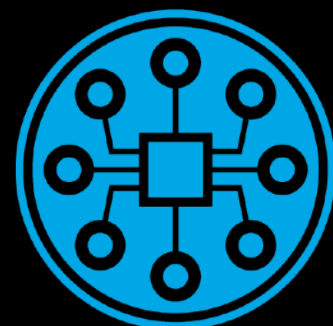
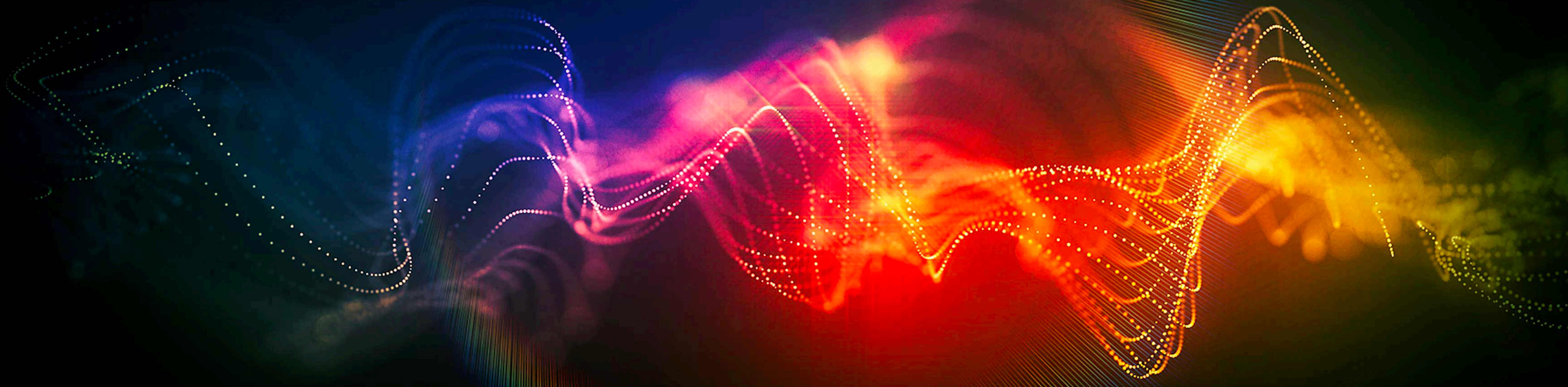
P. A. Bushev, J. Bourhill, M. Goryachev, N. Kukharchyk, E. Ivanov, S. Galliou, M. E. Tobar, and S. Danilishin
Phys. Rev. D **100**, 066020 – Published 20 September 2019



High Frequency Gravitational Waves

QDM

Quantum Technologies and Dark Matter Research Lab



EQUS

Australian Research Council
Centre of Excellence for
Engineered Quantum Systems



High Frequency Gravitational Waves

Gravitational wave detection with high frequency phonon trapping acoustic cavities

Maxim Goryachev and Michael E. Tobar

Phys. Rev. D **90**, 102005 – Published 24 November 2014

Phys. Rev. D **90**, 102005 – Published 24 November 2014

Maxim Goryachev and Michael E. Tobar

There are a number of theoretical predictions for astrophysical and cosmological objects, which emit high frequency ($10^6 - 10^9$ Hz) gravitation waves (GW) or contribute somehow to the stochastic high frequency GW background. Here we propose a new sensitive detector in this frequency band, which is based on existing cryogenic ultrahigh quality factor quartz bulk acoustic wave cavity technology, coupled to near-quantum-limited SQUID amplifiers at 20 mK. We show that spectral strain sensitivities reaching 10^{-22} per $\sqrt{\text{Hz}}$ per mode is possible, which in principle can cover the frequency range with multiple (> 100) modes with quality factors varying between 10^6 and 10^{10} allowing wide bandwidth detection. Due to its compactness and well-established manufacturing process, the system is easily scalable into arrays and distributed networks that can also impact the overall sensitivity and introduce coincidence analysis to ensure no false detections.

High Frequency Gravitational Waves

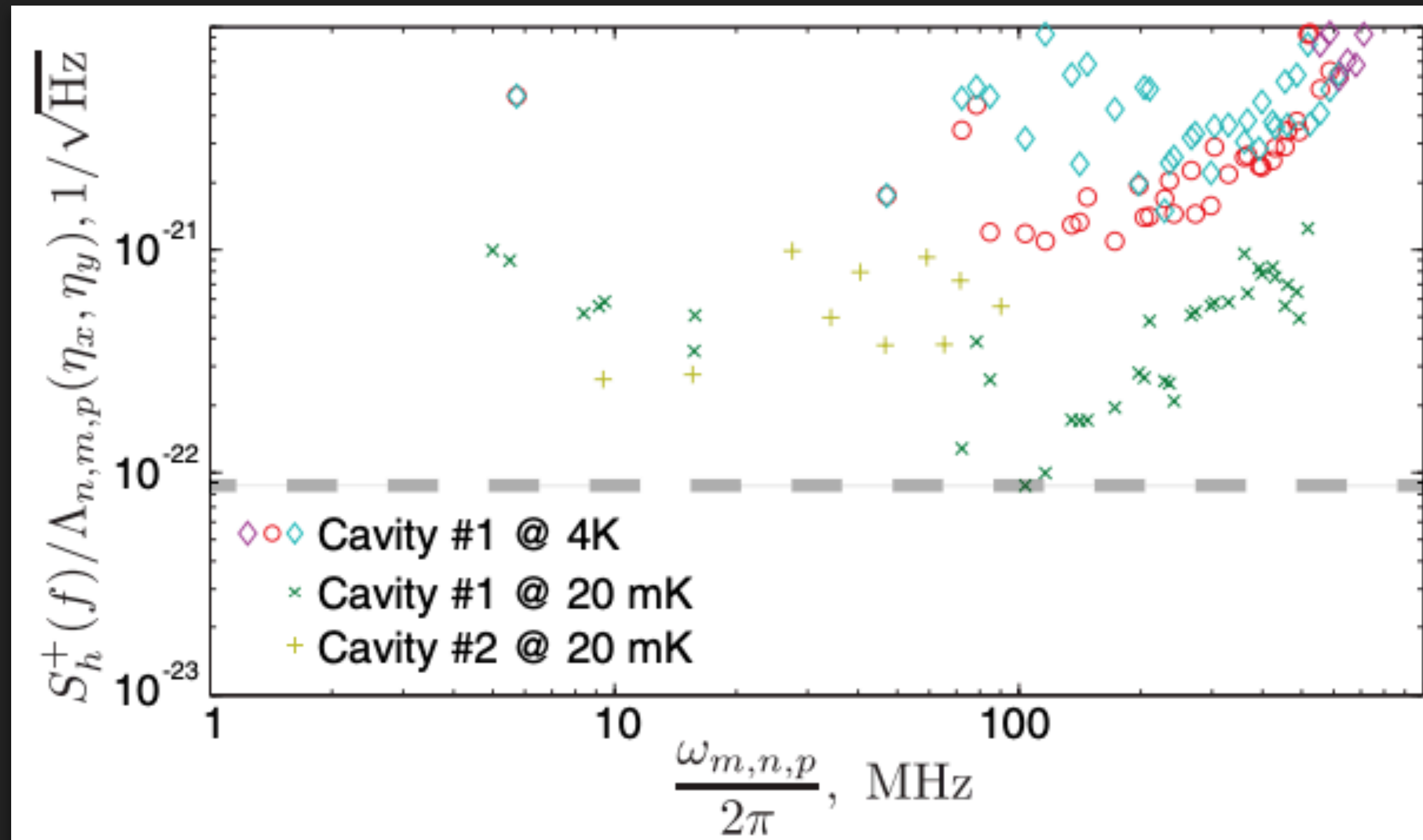
Gravitational wave detection with high frequency phonon trapping acoustic cavities

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MAXIM GORYACHEV AND MICHAEL E. TOBAR



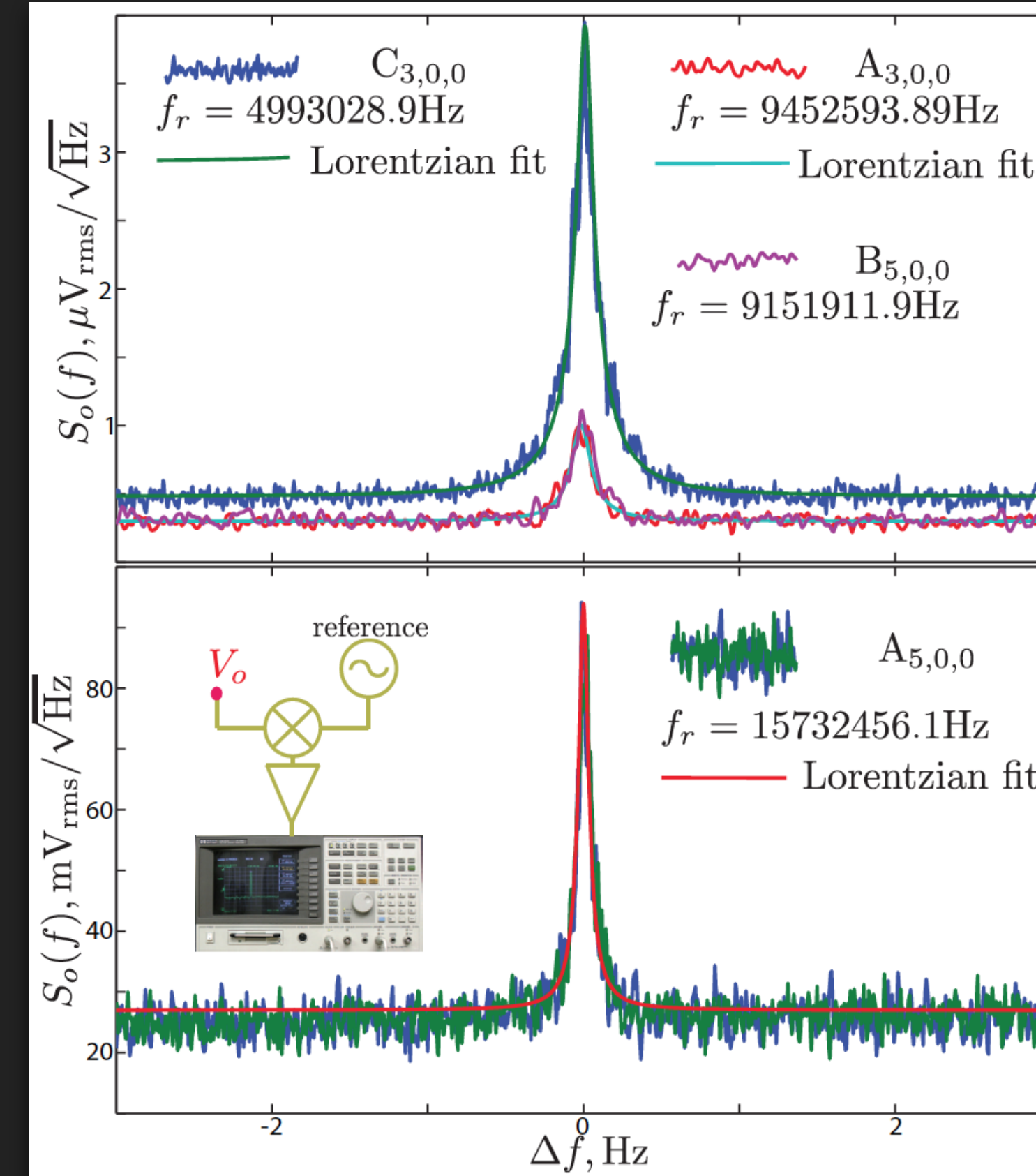
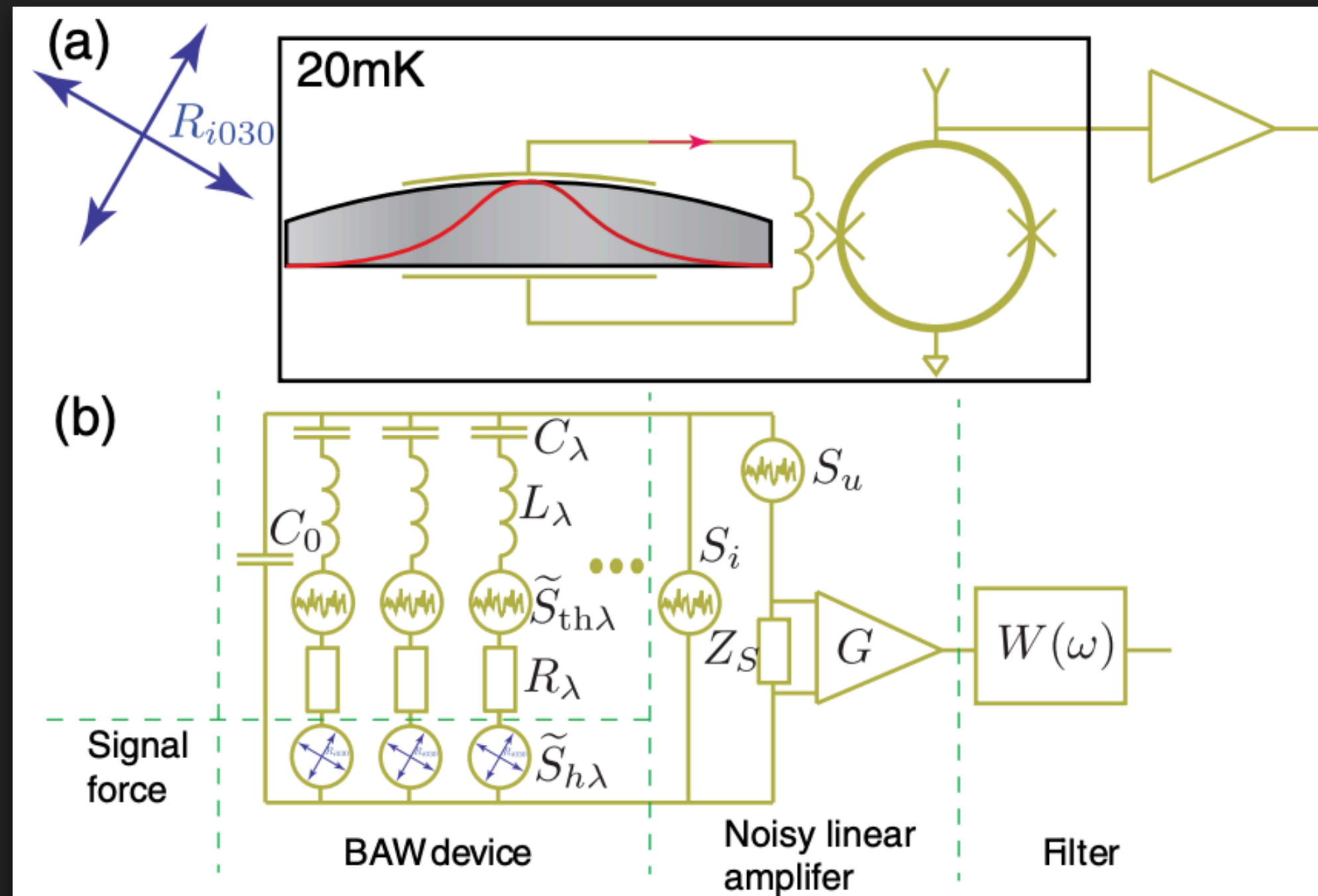
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 MAXIM GORYACHEV AND MICHAEL E. TOBAR



High Frequency Gravitational Waves

arXiv.org > gr-qc > arXiv:2011.12414

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General Relativity and Quantum Cosmology

[Submitted on 24 Nov 2020]

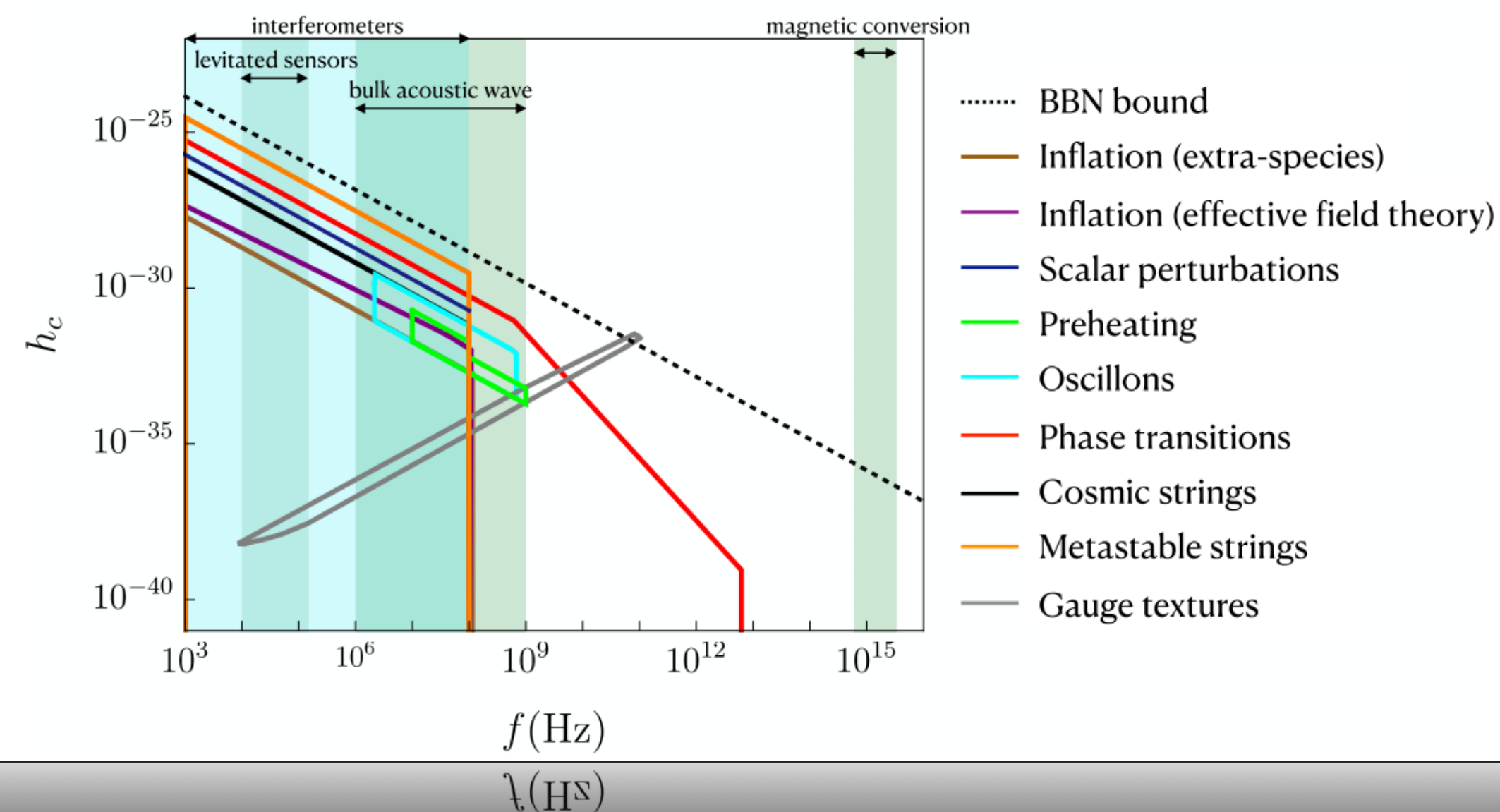
Challenges and Opportunities of Gravitational Wave Searches at MHz to GHz Frequencies

N. Aggarwal, O. D. Aguiar, A. Bauswein, G. Cella, S. Clesse, A. M. Cruise, V. Domcke, D. G. Figueroa, A. Geraci, M. Goryachev, H. Grote, M. Hindmarsh, F. Muia, N. Mukund, D. Ottaway, M. Peloso, F. Quevedo, A. Ricciardone, J. Steinlechner, S. Steinlechner, S. Sun, M. E. Tobar, F. Torrenti, C. Unal, G. White

The first direct measurement of gravitational waves by the LIGO and Virgo collaborations has opened up new avenues to explore our Universe. This white paper outlines the challenges and gains expected in gravitational wave searches at frequencies above the LIGO/Virgo band, with a particular focus on the MHz and GHz range. The absence of known astrophysical sources in this frequency range provides a unique opportunity to discover physics beyond the Standard Model operating both in the early and late Universe, and we highlight some of the most promising gravitational sources. We review several detector concepts which have been proposed to take up this challenge, and compare their expected sensitivity with the signal strength predicted in various models. This report is the summary of the workshop "Challenges and opportunities of high-frequency gravitational wave detection" held at ICTP Trieste, Italy in October 2019.

workshop "Challenges and opportunities of high-frequency gravitational wave detection" held at ICTP Trieste, Italy in October 2019.

detector
purpose of
gravitational

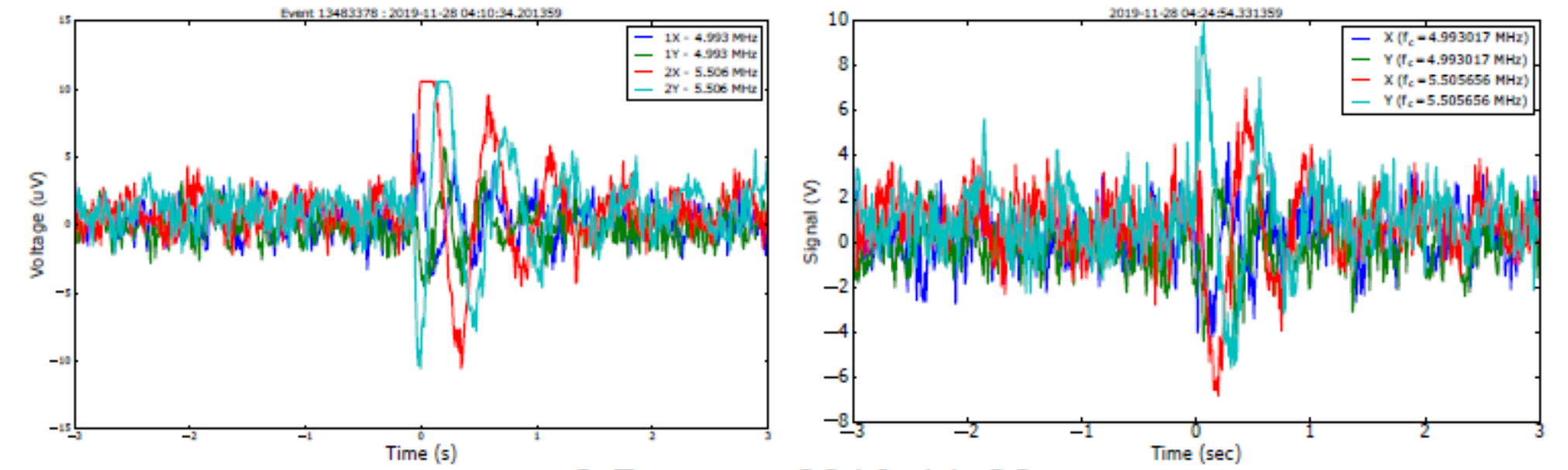
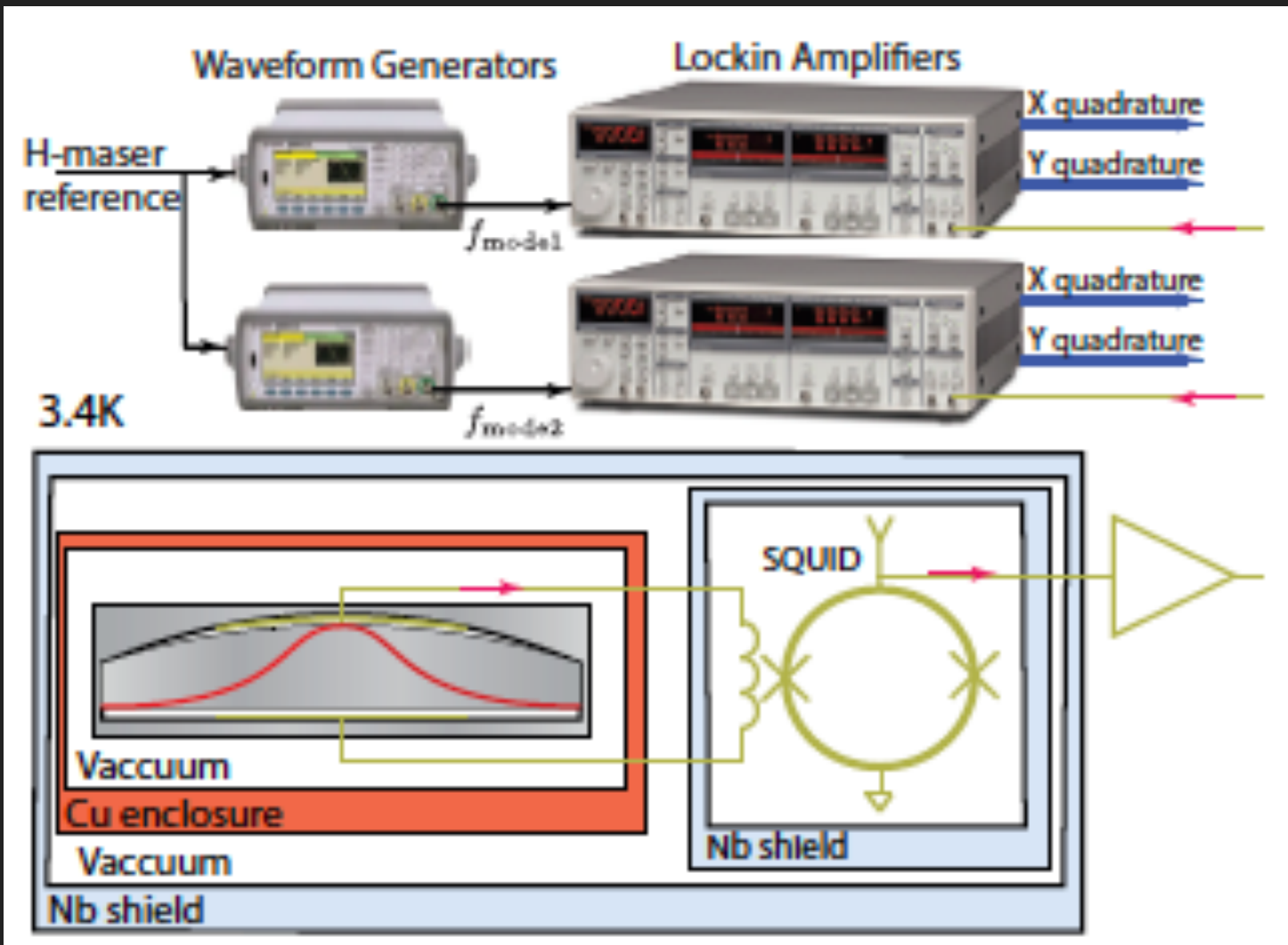


strength predicted in various models. This report is the summary of the
highlight some of the most promising gravitational sources. We review several
absence of known astrophysical sources in this frequency range provides a

- Neutron star mergers
- Light primordial black hole mergers
- Exotic compact objects
- Black hole superradiance
- Inflation
- Preheating
- Phase transitions
- Topological defects
- Evaporating primordial black holes

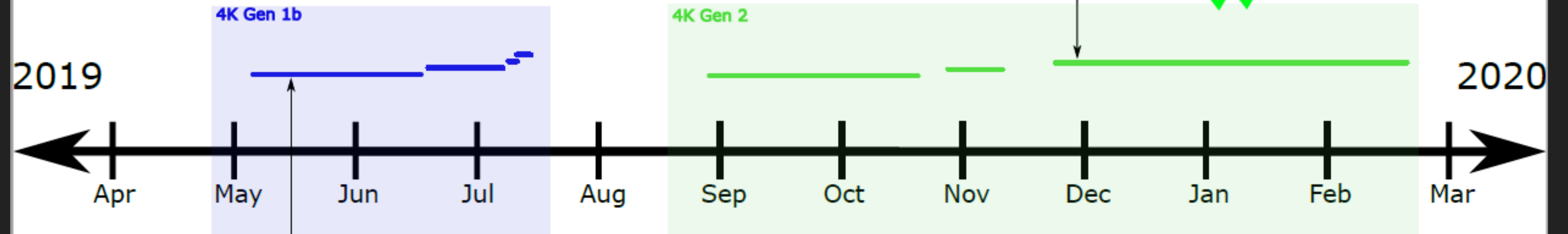
There is a number of theories predicting GW cosmic sources for the frequency range 1-1000 MHz, other theories predict GW from the early Universe.

First Detection?



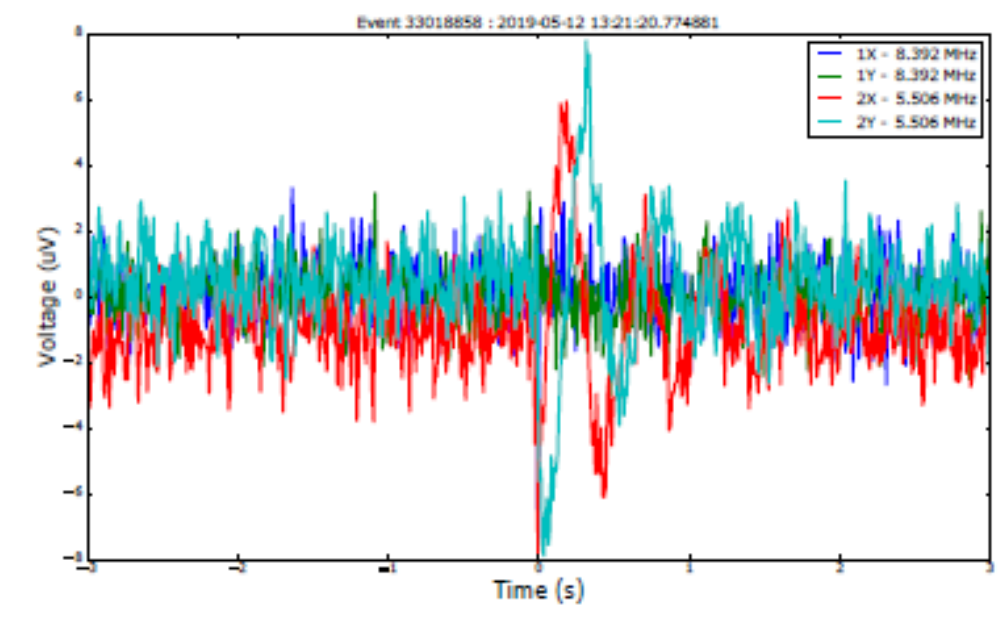
2 Events : 2019-11-28

Mode 1: ✓
Mode 2: ✓✓



Mode 1: 5C - 8.392 MHz
Mode 2: 3B - 5.506 MHz

Mode 1: 3C - 4.993 MHz
Mode 2: 3B - 5.506 MHz



Mode 1: ✗
Mode 2: ✓✓

Event : 2019-05-12

153 days of observation



First Detection?

arXiv.org > gr-qc > arXiv:2102.05859

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General Relativity and Quantum Cosmology

[Submitted on 11 Feb 2021]

Rare Events Detected with a Bulk Acoustic Wave High Frequency Gravitational Wave Antenna

Maxim Goryachev, William M. Campbell, Ik Siong Heng, Serge Galliou, Eugene N. Ivanov, Michael E. Tobar

This work describes the operation of a High Frequency Gravitational Wave detector based on a cryogenic Bulk Acoustic Wave (BAW) cavity and reports observation of rare events during 153 days of operation over two separate experimental runs (Run 1 and Run 2). In both Run 1 and Run 2 two modes were simultaneously monitored. Across both runs, the 3rd overtone of the fast shear mode (3B) operating at 5.506 MHz was monitored, while in Run 1 the second mode was chosen to be the 5th OT of the slow shear mode (5C) operating at 8.392 MHz. However, in Run 2 the second mode was selected to be closer in frequency to the first mode, and chosen to be the 3rd overtone of the slow shear mode (3C) operating at 4.993 MHz. Two strong events were observed as transients responding to energy deposition within acoustic modes of the cavity. The first event occurred during Run 1 on the 12/05/2019 (UTC), and was observed in the 5.506 MHz mode, while the second mode at 8.392 MHz observed no event. During Run 2, a second event occurred on the 27/11/2019(UTC) and was observed by both modes. Timing of the events were checked against available environmental observations as well as data from other detectors. Various possibilities explaining the origins of the events are discussed.

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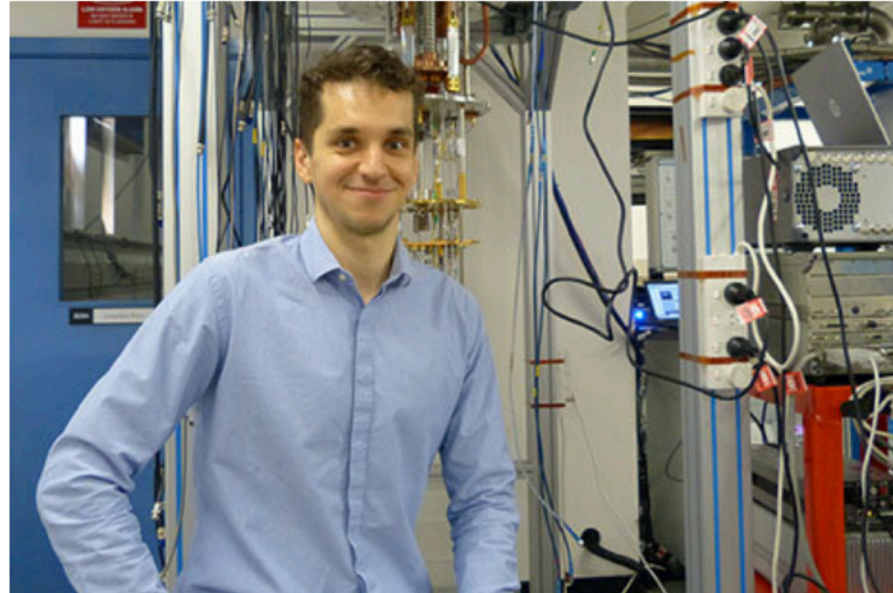
LIGO/VIRGO event catalogue, weather perturbations, earthquakes, meteor events / cosmic showers, FRBs

Possible sources:

Internal solid state processes, internal radioactive events, cosmic ray events, HFGW sources, domain walls, WIMPs, dark matter



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Director



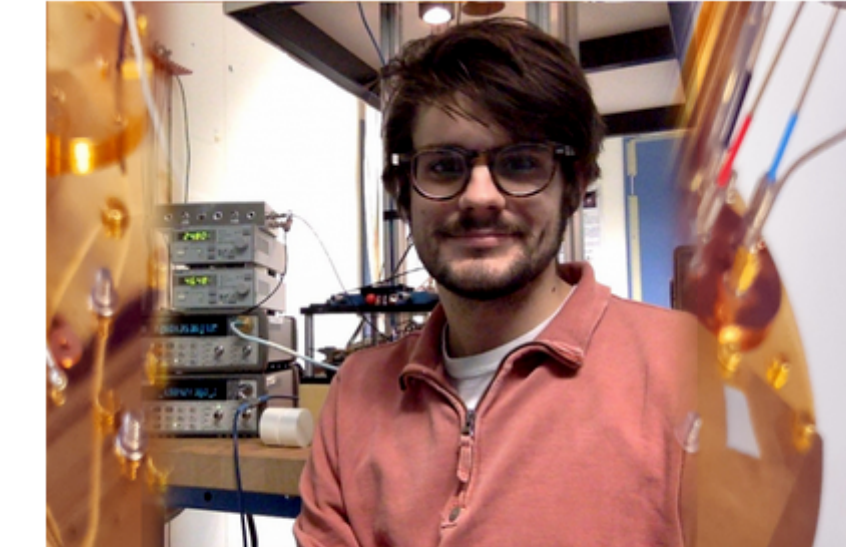
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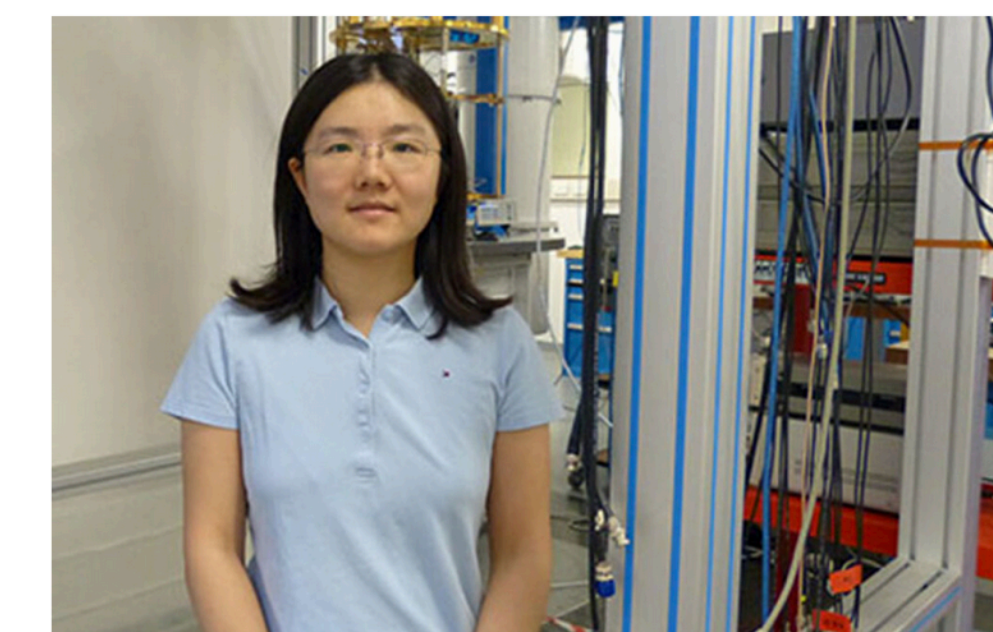
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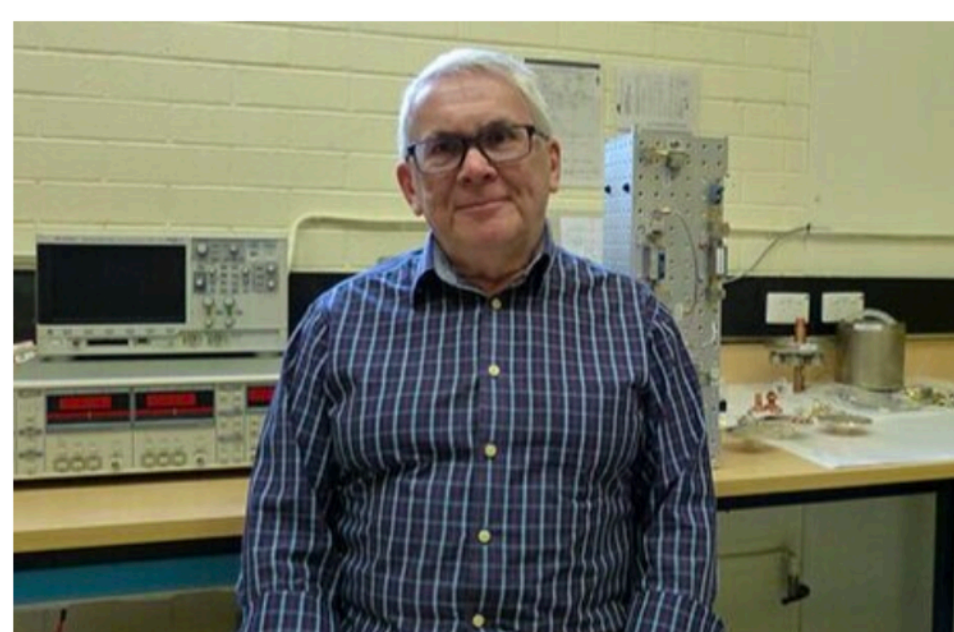
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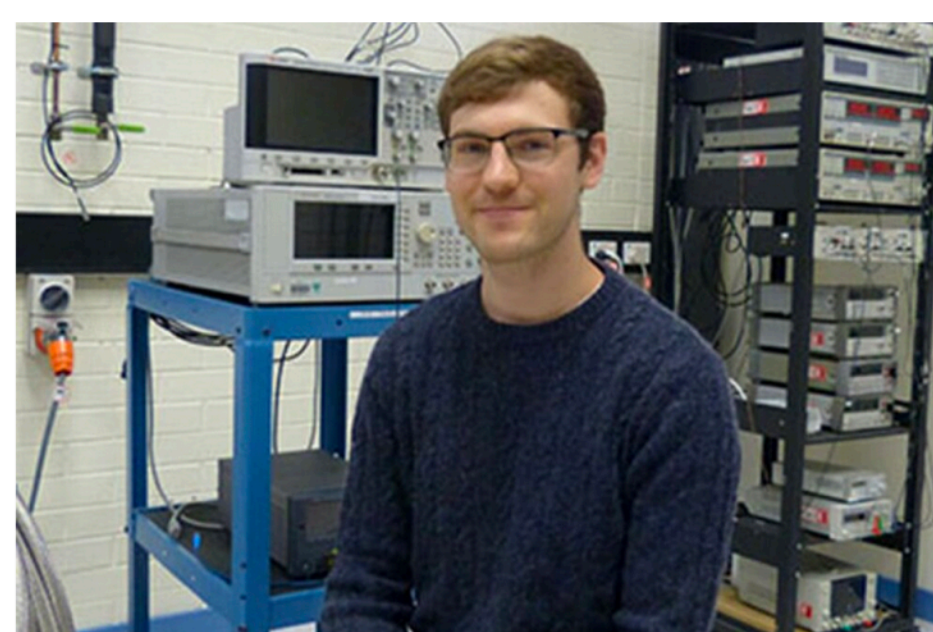
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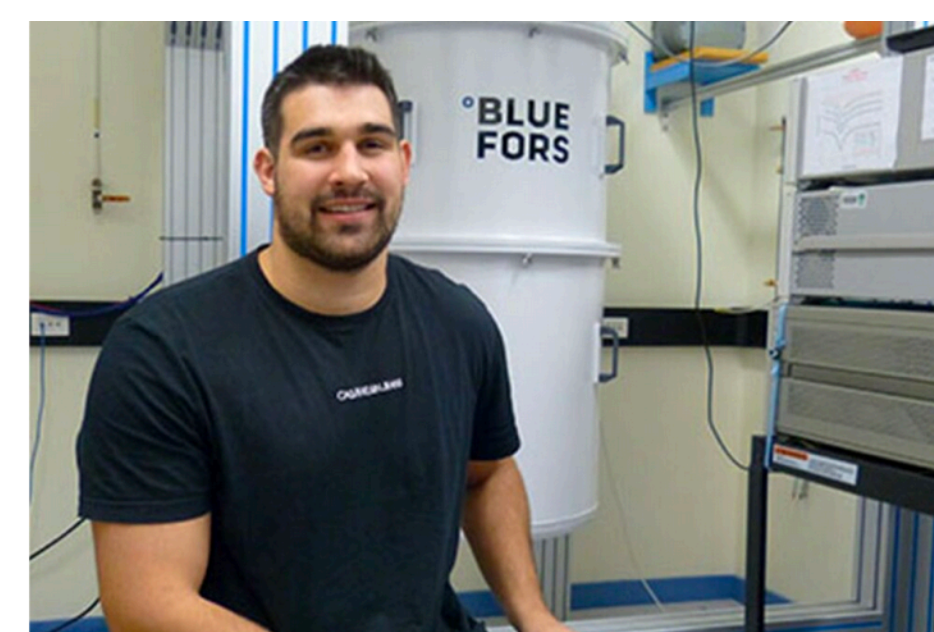
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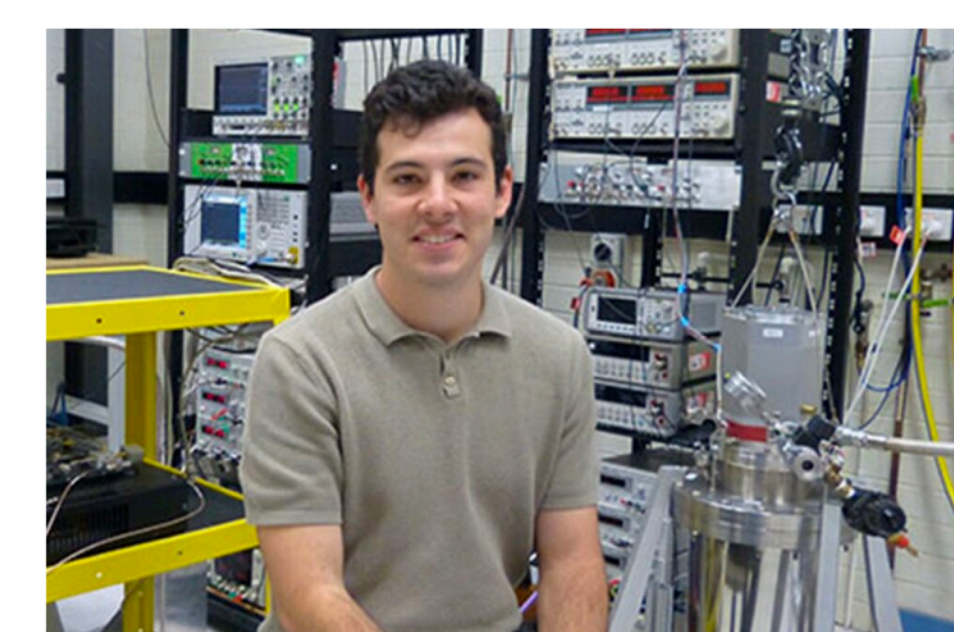
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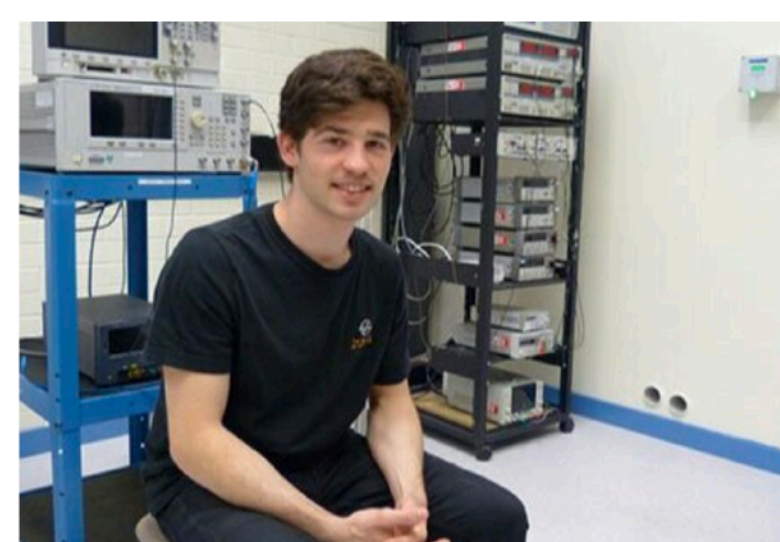
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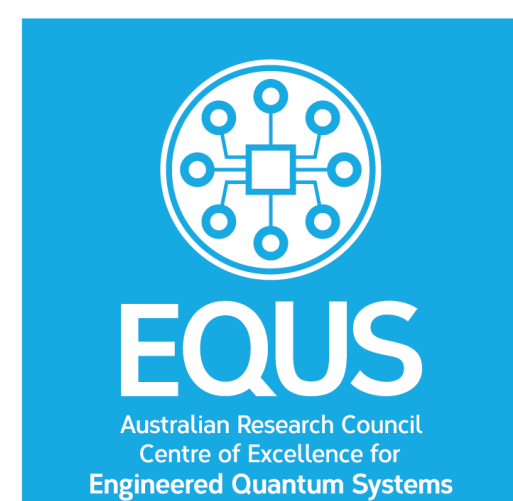
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