

Cosmic Strings in generalized hybrid metric-Palatini gravity

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Overview

1. Hybrid metric-Palatini
2. Cosmic Strings
3. Solutions for different scalar field potentials
4. Conclusions

Addressing the imbalance - Motivation for Modified Gravity

- Adding a previously unaccounted source. (e.g Λ ,...)
- Changing the governing equations.(e.g Gauss-Bonnet gravity, $f(R)$ gravity,...)

Hybrid metric-Palatini

Both metric (arXiv:astro-ph/0306438) and Palatini(arXiv:1101.3864) approach are successful in reproducing acceleration rate of the Universe, but are flawed, particularly changing the solar system dynamics. So, a new approach was proposed(arXiv:1508.04641), the hybrid metric-Palatini:

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} f(R, \mathcal{R}) + S_m, \quad (1)$$

$$\mathcal{R}_{\mu\nu} \equiv \hat{\Gamma}_{\mu\nu,\alpha}^{\alpha} - \hat{\Gamma}_{\mu\alpha,\nu}^{\alpha} + \hat{\Gamma}_{\alpha\lambda}^{\alpha} \hat{\Gamma}_{\mu\nu}^{\lambda} - \hat{\Gamma}_{\mu\lambda}^{\alpha} \hat{\Gamma}_{\alpha\nu}^{\lambda}. \quad (2)$$

Dynamic equations(in scalar-tensor representation)

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left[(\varphi - \psi)R + \frac{3}{2\phi} \partial_\mu \psi \partial^\mu \psi - V(\varphi, \psi) \right] + S_m . \quad (3)$$

Performing the variation of the action with respect to the metric and the scalar fields, and redefining φ as $\xi^2 = \varphi - \psi$ yields the field equations

$$\xi^2 G_{\mu\nu} = \kappa^2 T_{\mu\nu} + \nabla_\mu \nabla_\nu \xi^2 + \frac{3}{2\psi} \partial_\mu \psi \partial_\nu \psi - \left(\square \xi^2 + \frac{1}{2} \bar{V} + \frac{3}{4\psi} \partial^\alpha \psi \partial_\alpha \psi \right) g_{\nu\mu} \quad (4)$$

and

$$\square \psi - \frac{1}{2\psi} \partial_\mu \psi \partial^\mu \psi - \frac{\psi}{3} \left[\frac{1}{2\xi} \bar{V}_{,\xi} + \bar{V}_{,\psi} \right] = 0 \quad (5)$$

$$\square \xi^2 + \frac{1}{2\psi} \partial_\mu \psi \partial^\mu \psi + \frac{1}{6} [4\bar{V} - \xi \bar{V}_{,\xi}] = \frac{\kappa^2}{3} T , \quad (6)$$

Cosmic strings

- Topological defects formed at phase transitions after SSB.
- Featured on several Grand Unified scenarios.
- Possible impact on CMB anisotropies, small scale structure formation, Gamma Ray Bursts, stochastic background GW spectra.

Straight infinite Cosmic String

$$T_t^t = T_z^z = -\sigma(r), \quad (7)$$

where σ is the string tension. We consider a general cylindrically symmetric static metric

$$ds^2 = -e^{2(K-U)} dt^2 + e^{2(K-U)} dr^2 + e^{-2U} W^2 d\theta^2 + e^{2U} dz^2, \quad (8)$$

where t , r , θ and z denote the time, radial, angular and axial cylindrical coordinates, respectively, and K , U and W are functions of r alone.

Full set of equations

In terms of two dynamical variables defined as $u = \varphi'$ and $v = \psi'$

$$\frac{dW}{dr} = \frac{1}{u-v} \left(\frac{3v^2}{4\psi} - \frac{V}{2} \right) W \quad (9)$$

$$\frac{du}{dr} = -\frac{v}{u-v} \left(\frac{3v^2}{4\psi} - \frac{V}{2} \right) - \frac{1}{4} \frac{v^2}{\psi} + \frac{\psi}{3} (V_\psi + V_\varphi) - \frac{V}{2} \quad (10)$$

$$\frac{dv}{dr} = -\frac{v}{u-v} \left(\frac{3v^2}{4\psi} - \frac{V}{2} \right) + \frac{1}{2} \frac{v^2}{\psi} + \frac{\psi}{3} (V_\psi + V_\varphi) \quad (11)$$

$$\kappa^2 \sigma = \frac{1}{2} \left[V(\varphi, \psi) + V_\varphi(\varphi, \psi) - \frac{3v^2}{2\psi} \right] \quad (12)$$

Potential function and parameter space

- We need to specify a potential function $V(\varphi, \psi)$.
- We need also a set of boundary values $(\varphi(0), \psi(0), W(0), u(0), v_0)$
- In addition, more parameters can "appear" in the form of the potential.

Exact Solutions

- $V(\varphi, \psi) = 0$

$$W(r) = \frac{3c_1(r + c_2)}{4w_0} \qquad -\kappa^2\sigma = \frac{4w_0^2}{3(c_2+r)^2} \qquad (13)$$

$$(14)$$

- $V(\varphi, \psi) = \Lambda = \text{const.}$

$$W(r) = W_0(r + 2c_1) \qquad \kappa^2\sigma = \frac{\Lambda}{2} - \frac{3c_2}{4(2c_1+r^2)} \qquad (15)$$

$$(16)$$

Numerical solutions

We studied numerical solutions for different types of potentials:

- $V(\varphi, \psi) = a\varphi^2 + b\psi^2$
- $V(\varphi, \psi) = ae^{-\lambda_1\varphi} + be^{-\lambda_2\psi}$
- $V(\varphi, \psi) = a\varphi^2\psi^2$
- $V(\varphi, \psi) = ae^{-\lambda_1\varphi - \lambda_2\psi}$

What we learned?

- Unlike in Brans-Dicke gravity, string solutions are found in GHMP.
- String tension not necessarily positive.
- Number of parameters of the models allow a wide range of solutions.
- Solutions very sensitive on the initial conditions.
- Not asymptotically flat.

Behaviour of the string tension

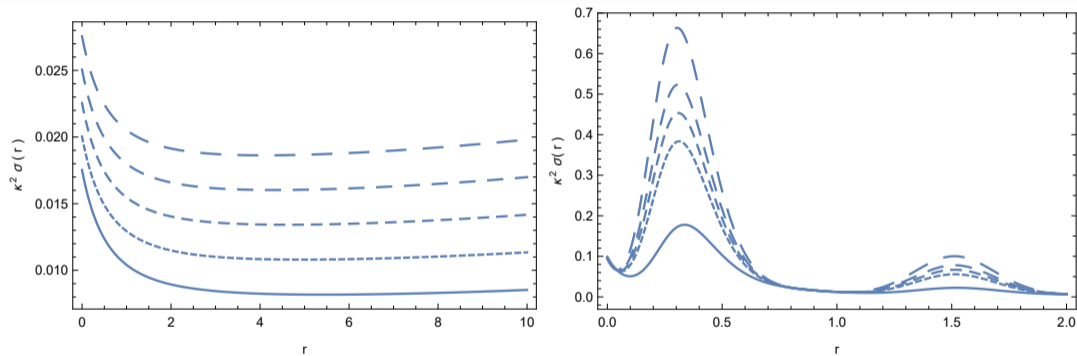


Figure: Variations of the string tension $\sigma(r)$ as a function of r for $V(\varphi, \psi) = a\varphi^2\psi^2$ (right panel) and for $V(\varphi, \psi) = ae^{-\lambda_1\varphi-\lambda_2\psi}$ (left panel) for different values of a .

Thank you.