

# **The approach to equilibrium for idealized collisionless self-gravitating systems**

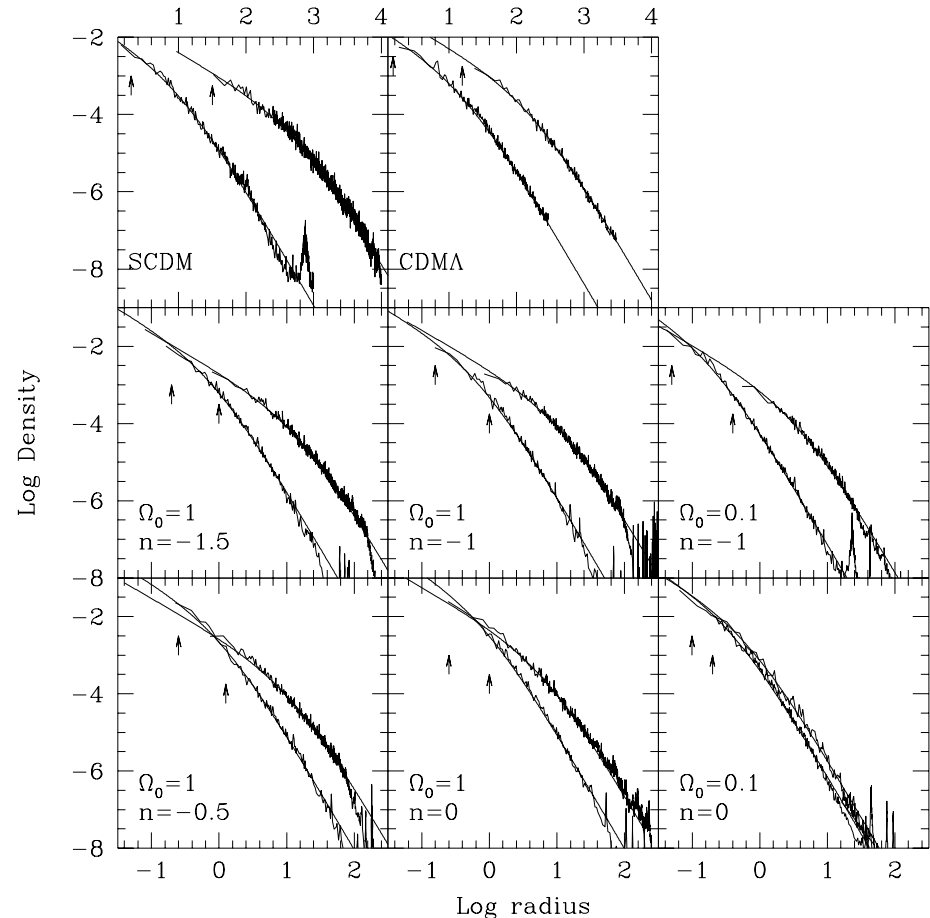
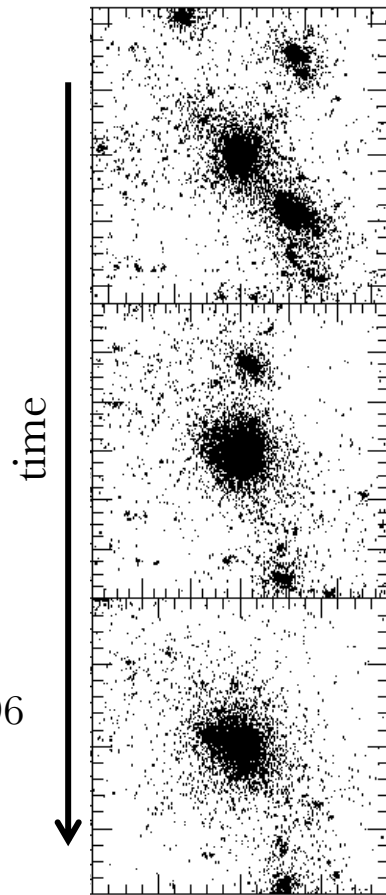
*Liliya L.R. Williams (University of Minnesota, US)*

*Jens Hjorth (University of Copenhagen, DK)*

# Introduction

Collisionless collapse of self-gravitating systems results in nearly universal mass density profiles of equilibrium halos

*NFW*



- Dubinski & Carlberg 1991
- Navarro, Frenk & White 1996
- Navarro et al. 2004
- Springel et al. 2008
- Boylan-Kolchin et al. 2009
- Vogelsberger et al. 2014

# Introduction

Theoretical explanation of this density profile shape and its universality have been more challenging.

*Many attempts were made:*

Ogorodnikov 1957

Lynden-Bell 1967

Shu 1978

Madsen 1987

Stiavelli & Bertin 1987

Spergel & Hernquist 1992

Hjorth & Madsen 1991, 1993

Plastino & Plastino 1993

White & Narayan 1987

MacMillan et al. 2006

Kang & He 2011

Barnes & Williams 2012

Pontzen & Governato 2013

Beraldo e Silva et al. 2014, 2017

Levin & Pakter 2014

Chavanis et al. 2015

Wagner 2020

Many authors have noted that

- ❖ large particle number suggests the use of statistical mechanics
- ❖ universality of the equilibrium state may correspond to the most likely state  
→ use maximum entropy

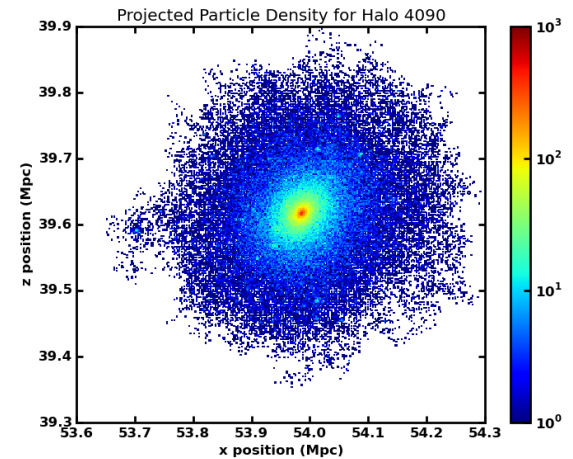
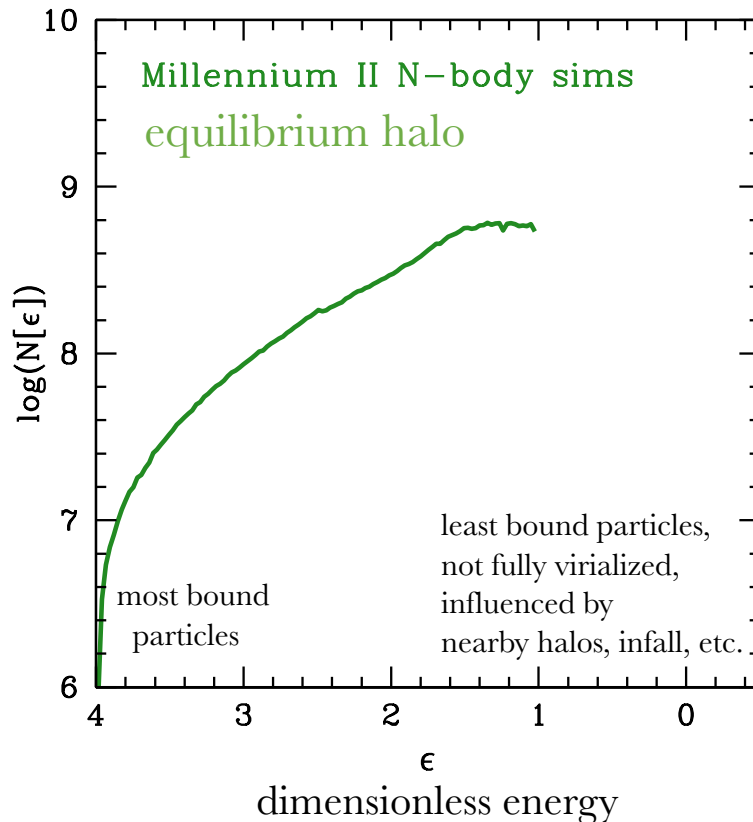
# Stat. mech. $\rightarrow$ energy distribution, $n(E)$

Density profiles are important,  
but in statistical mechanics we work with  $n(E)$ , where  $E = \text{kinetic} + \text{grav. potential}$

Why  $n(E)=f(E)g(E)$ , and not the distribution function  $f(E)$ ?

In an equilibrium collisionless system each particle has a fixed energy  
Differential energy distribution,  $n(E)$  gives the full description of a system

***A successful theory must reproduce  $n(E)$  of equilibrium  $N$ -body halos***



x-y projection of  
one equilibrium  
Millennium II halo

# Stat. mech. $\rightarrow$ energy distribution, $n(E)$

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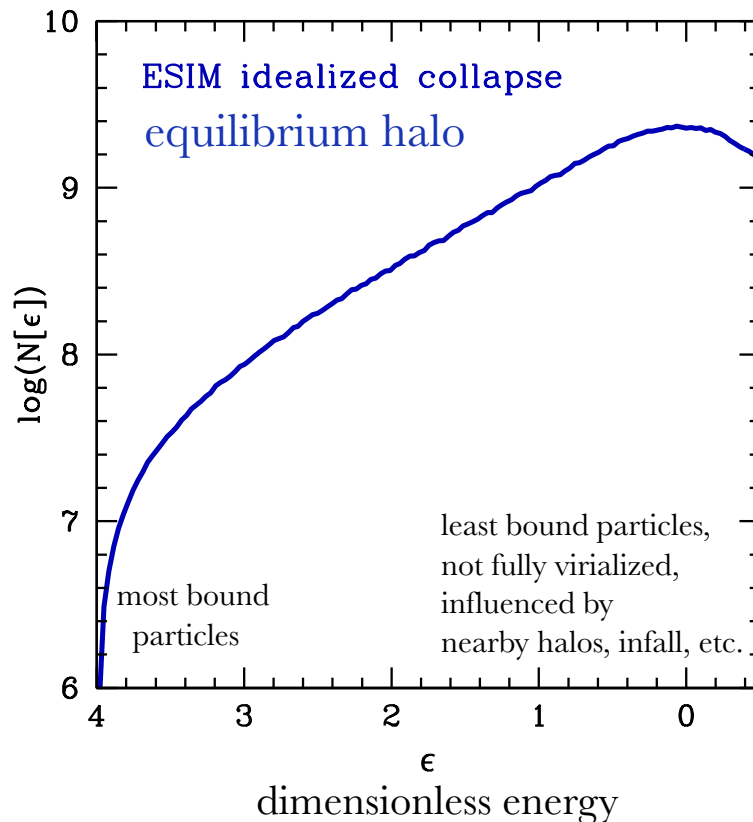
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***We use a simplified collapse model*** (*N-body simulations have complex dynamics*)



***Extended Secondary Infall Model***  
(**ESIM**, based on Ryden & Gunn (1987))

- mean-field model
- spherically symmetric
- radial forces only
- 3D phase-space (1 spatial + 2 velocity)
- radial, angular actions conserved

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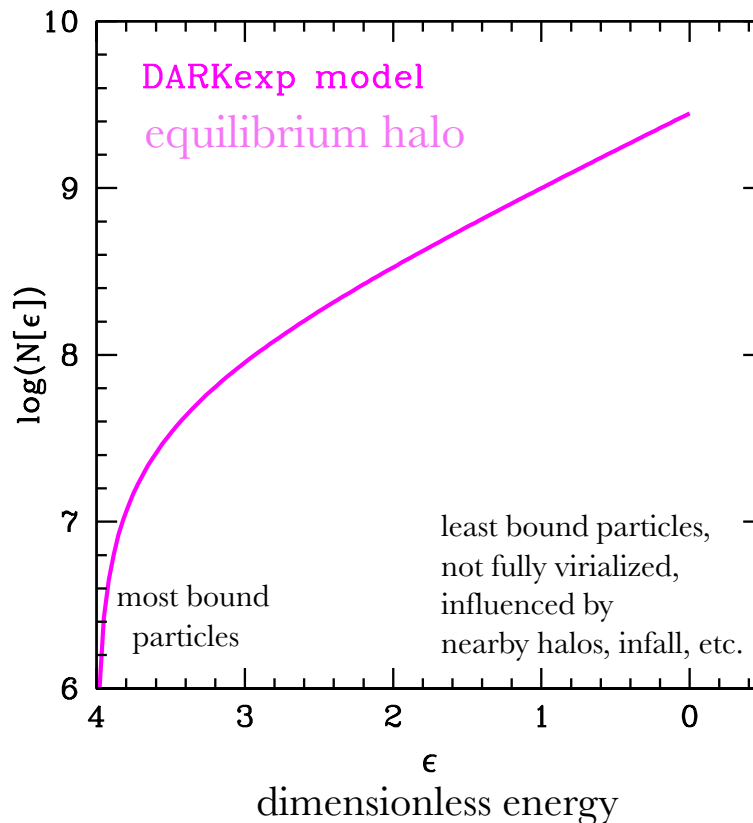
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## $n(E)$ of our theoretical model, *DARKexp*



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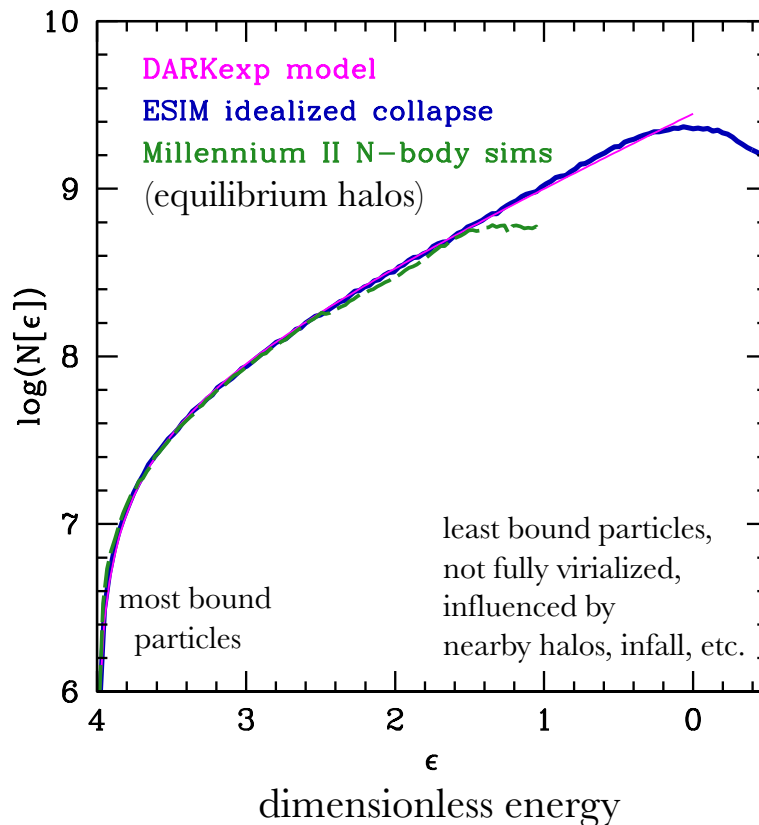
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**All 3: *N*-body sims, ESIM simplified collapse, DARKexp model**



DARKexp fits equilibrium states of *N*-body simulated halos and ESIM very well (also density and velocity dispersion profiles)

# DARKexp entropy $S_D$ derived

Entropy is related to the # microstates corresponding to a given macrostate

$$W = N! \prod_j \frac{g_j^{n_j}}{n_j!}$$

fully captured by  $n(E)$  want to derive its shape

$n$  – occupation # of cells in energy space  
 $N$  – total # of particles  
 $g$  – degeneracy / cell size  
 $j$  – energy level label

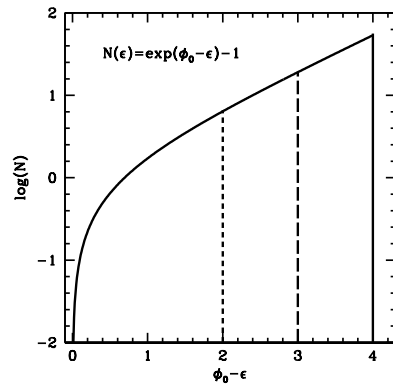
$$\ln W = \cancel{\ln N!} + \sum_j n_j \ln g_j - \sum_j \ln \Gamma(n_j + 1)$$

normalization irrelevant

To find the most likely state with a fixed total energy we define **entropy**:

$$S_D = \ln W - \beta \sum_j n_j E_j \Rightarrow S_D = - \int \ln \Gamma(n(E) + 1) dE - \beta \int n(E) E dE$$

$S_D$  should attain a maximum at equilibrium  $\frac{dS_D}{dn} = 0$   $\frac{d^2 S_D}{dn^2} < 0$



$$n(E) \propto e^{-\beta(E-E_0)} - 1$$

**DARKexp**



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$S_D$  should increase with time, as halo evolves

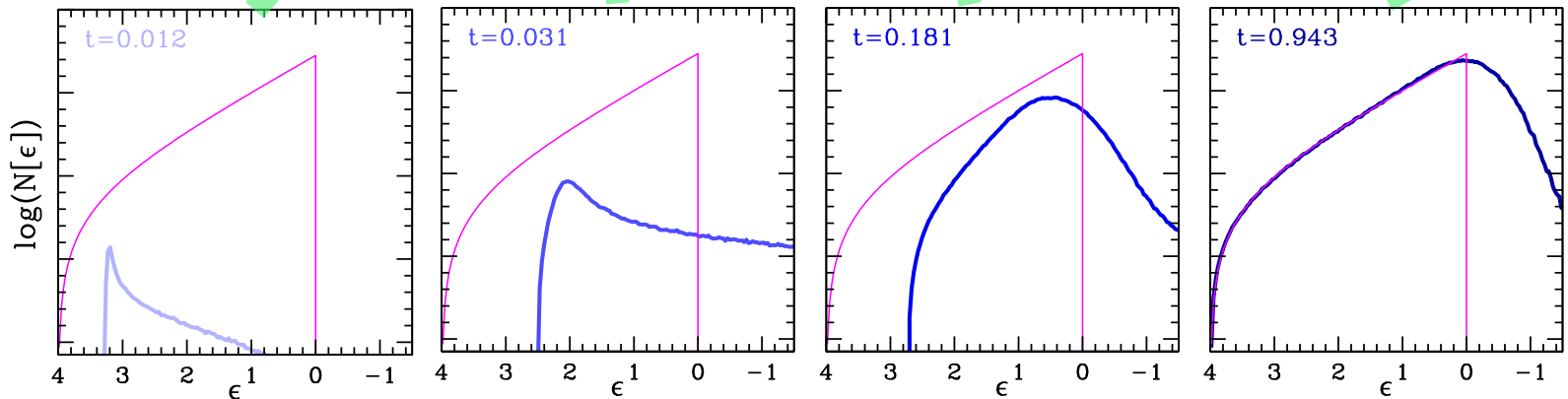
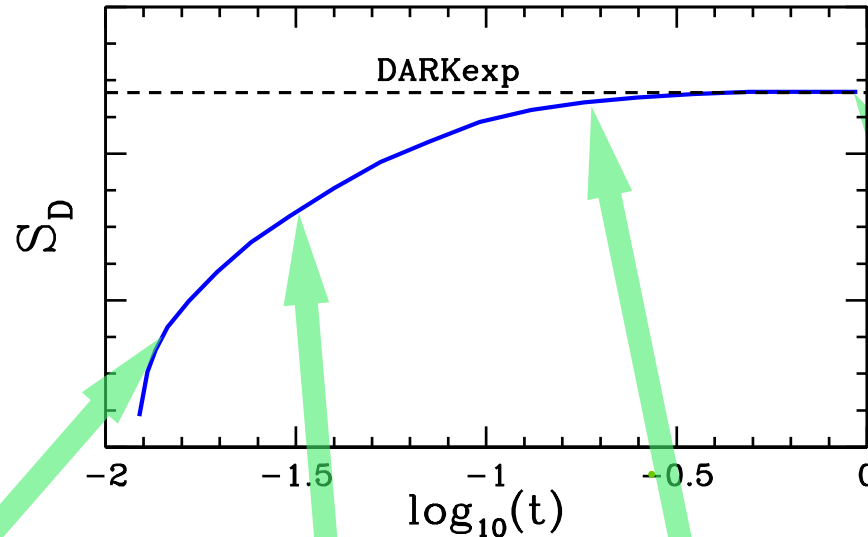
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**DARKexp**

# DARKexp entropy $S_D$ increases with time

$$S_D = -\int \ln \Gamma(n(E)+1) dE - \beta \int n(E) E dE$$

evolution of  
one ESIM halo

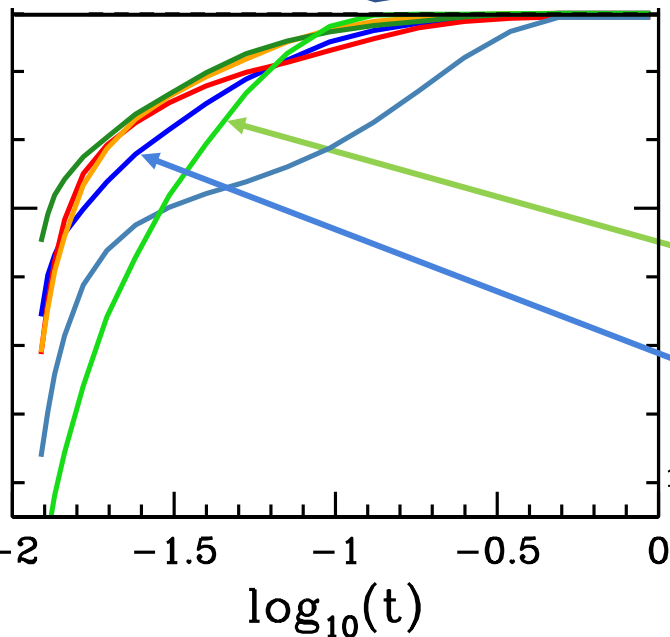


*ESIM  $n(E)$  shape approaches that of DARKexp*

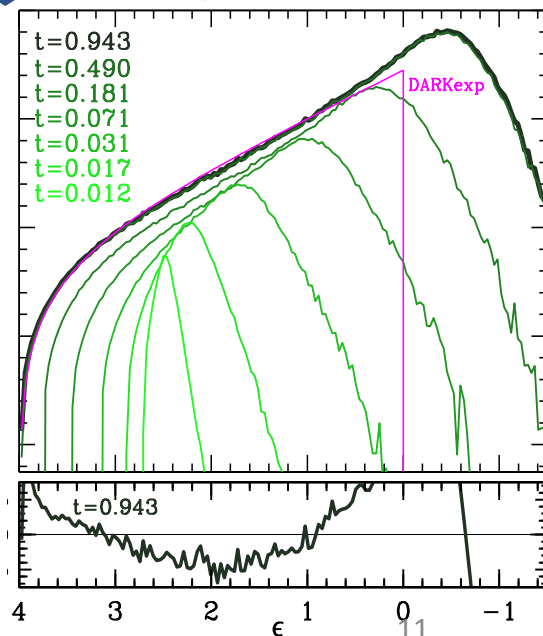
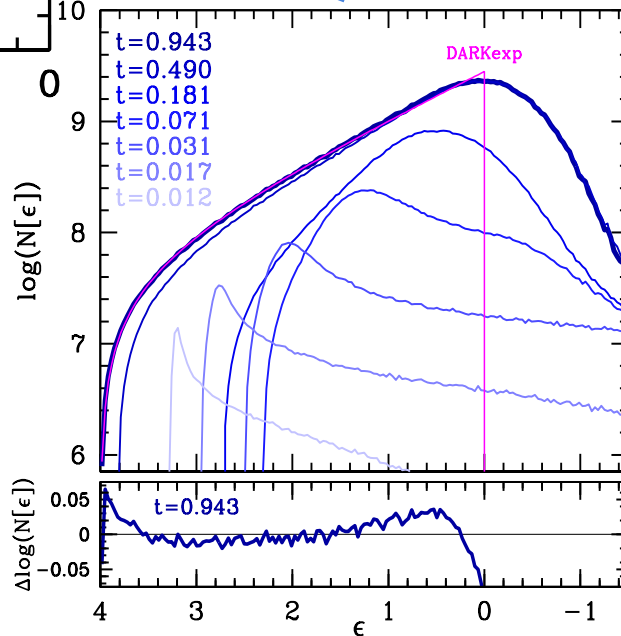
**Entropy  $S_D$  increases with time**

**Shape of  $n(E)$  approaches DARKexp**

$$S_D = -\int \ln \Gamma(n(E)+1) dE - \beta \int n(E) E dE$$



**many ESIM halos  
with different  
initial conditions**



# Conclusions

## DARKexp:

Derived using principles of statistical mechanics

Derivation is straightforward

(with 1 conceptual difference from Maxwell-Boltzmann stats)

Energy distribution  $n(E)$  has no shape parameters,  
only a truncation near unbound energies

$$n(E) \propto e^{-\beta(E-E_0)} - 1$$

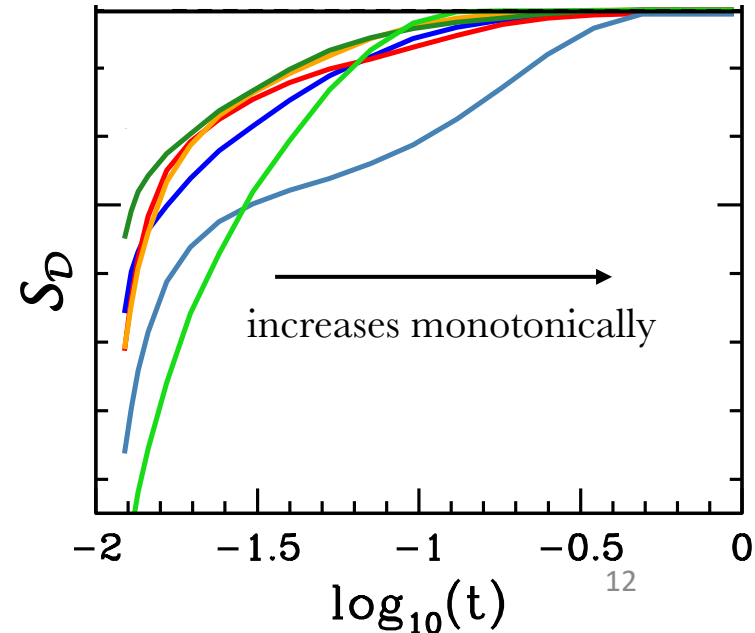
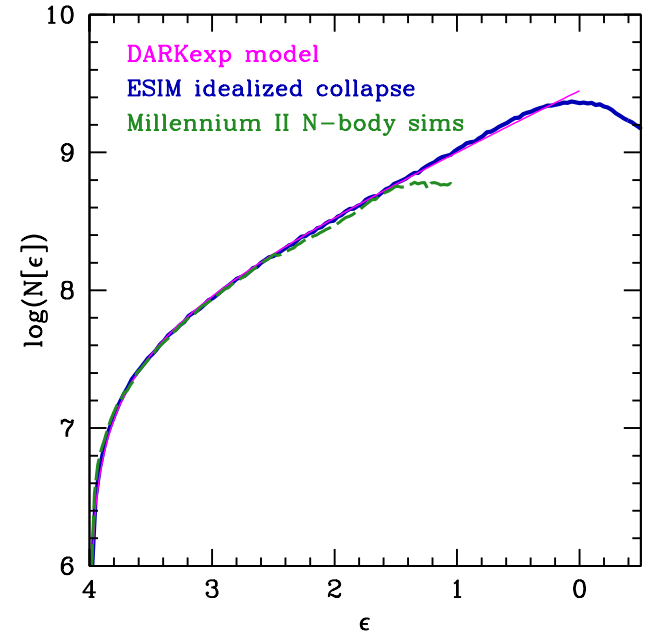
Reproduces  $n(E)$ , density & velocity dispersion profiles of  
N-body sims and simplified collapse of ESIM halos

## Entropy $S_D$

$$S_D = -\int \ln \Gamma(n(E)+1) dE - \beta \int n(E) E dE$$

increases monotonically during evolution  
of ESIM halos and attains a maximum  
at equilibrium

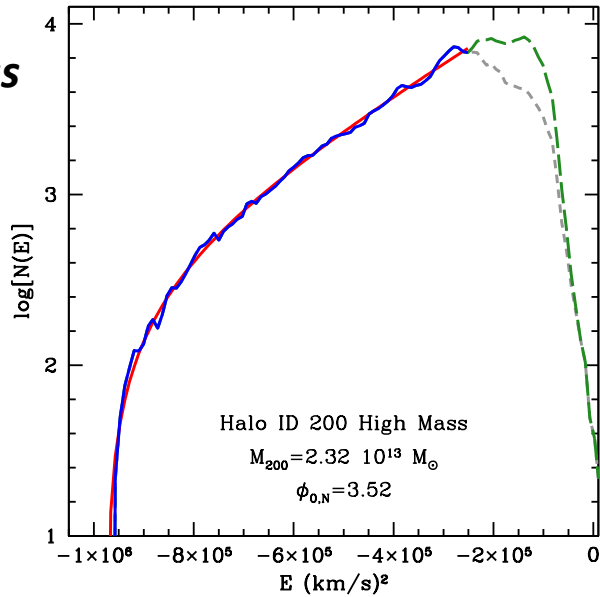
**→ self-consistent theory**



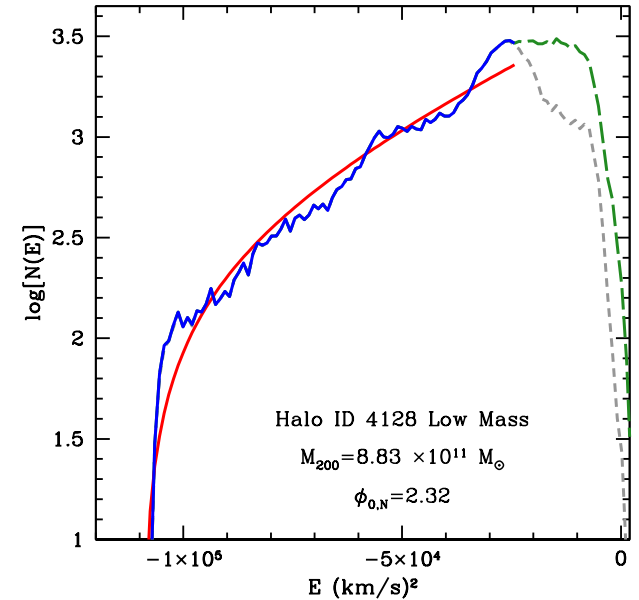
# EXTRA SLIDES

# Individual Millennium-II halos fitted with DARKexp

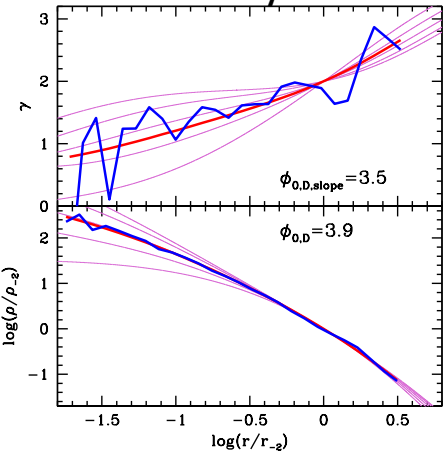
**High mass**



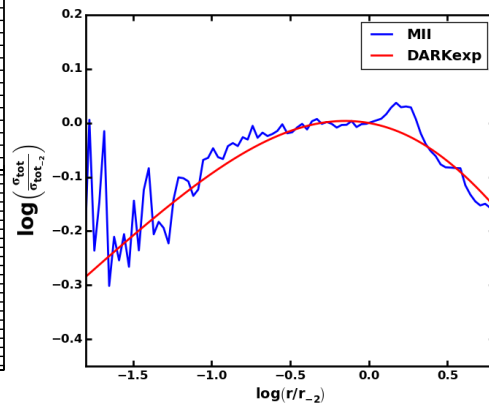
**Low mass**



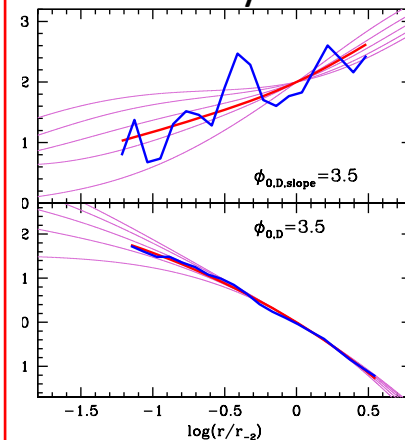
**density**



**velocity dispersion**



**density**



**velocity dispersion**

