

How Classical are Gaussian States in Loop Quantum Cosmology?

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Motivation

We do not typically observe fundamentally quantum features of spacetime. If loop quantum gravity is to provide an adequate picture of reality, it must have a semi-classical sector that recovers certain cosmological phenomena in appropriate limits. In particular, loop quantum cosmology must admit semi-classical states.

We will say a quantum state is *kinematically* semi-classical if it satisfies the following:¹

- 1 It should approximate a state having a definite coordinate x .
- 2 It should have small quantum fluctuations in x and p .

¹Martin-Dussaud, Quantum, 2021.

When are fluctuations “small”?

In ideal conditions, $\Delta^2 x \Delta^2 p = 0$.

From the uncertainty relation,² we know:

$$\Delta^2 x \Delta^2 p \geq \underbrace{\frac{1}{2} |\langle \{x, p\} \rangle - \langle x \rangle \langle p \rangle|^2 + \left| \frac{1}{2i} \langle [x, p] \rangle|^2}_{\text{Robertson-Schrödinger}} \geq \underbrace{\left| \frac{1}{2i} \langle [x, p] \rangle \right|^2}_{\text{Heisenberg}}$$

“Small” means the RS inequality is saturated.

²Robertson, Phys. Rev., 1929.

Classicality

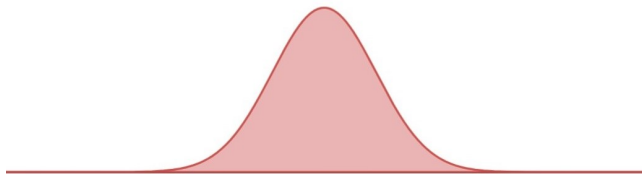
In traditional quantum mechanics, Gaussian wave functions satisfy these basic criteria:

$$|\psi\rangle = c \int dx e^{-(x-\mu)^2/2\sigma^2} |x\rangle,$$

Coordinates become sharply peaked as $\sigma \rightarrow 0$.

$$\Delta^2 x \Delta^2 p = \left| \frac{1}{2i} \langle [x, p] \rangle \right|^2$$

Fluctuations are minimal.



Question

Are Gaussian states semi-classical in LQC?

Foreshadowing

No! Not even at the kinematic level.
(Though they are asymptotically so.)

Consider an FLRW spacetime with a metric:

$$ds^2 = -dt^2 + a(t)^2[dx^2 + dy^2 + dz^2].$$

Then the spatial volume of a time-like slice $V = a^3(t)$ and the Hubble parameter $\beta = \dot{a}(t)/a(t)$ are canonical variables from which the EFEs may be derived using a Hamiltonian treatment via the Holst action.

LQC is the theory one obtains by applying loop quantization to V and β .

The Holst action has a divergent symplectic term. Thus, the loop quantization procedure must introduce volume regularization. One therefore considers LQC on a compact fiducial spacetime region with volume V_0 .³

³Ashtekar & Singh, *Class. Quant. Grav.*, 2011; Rovelli & Wilson-Ewing, *Phys. Rev. D*, 2014; Rovelli & Vidotto, CUP, 2015.

The “coordinate” observable in LQC is the volume operator \hat{V} . However, the loop quantization procedure does not allow for a unique operator $\hat{\beta}$.⁴ Instead, the conjugate momentum is chosen to be \hat{S}_λ , determined by the holonomy operator \hat{h}_λ :

$$\hat{V}|V\rangle = V|V\rangle, \quad \hat{h}_\lambda|V\rangle = \left| V + \frac{4\pi G\gamma\hbar\lambda}{V_0} \right\rangle, \quad \hat{S}_\lambda = \frac{1}{2i\lambda}(\hat{h}_\lambda - \hat{h}_\lambda^*).$$

Henceforth, $\alpha := 4\pi G\gamma\hbar$.

⁴Ashtekar, Bojowald, & Lewandowsky, Adv. Theor. Math. Phys., 2003.

Gaussian States in LQC

The kinematic volume Hilbert space is given by states of the form:

$$|\psi\rangle = \sum_{n \in \mathbb{Z}} f(n) |V_n\rangle, \quad \sum_{n \in \mathbb{Z}} |f(n)|^2 < \infty$$

for $V_n \in \mathbb{R}$, where $\langle V_n | V_m \rangle = \delta_{nm}$. The sign of V_n indicates the orientation of the spatial slice. Physical states in LQC must have parity symmetry:⁵

$$\psi(V) = \psi(-V)$$

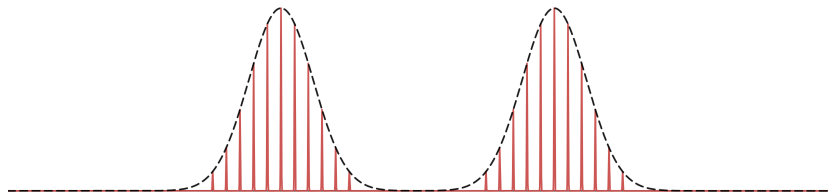
⁵Bentivegna & Pawłowski, Phys. Rev. D, 2008; Ashtekar & Wilson-Ewing, Phys. Rev. D, 2009; Ashtekar & Singh, Class. Quant. Grav., 2011.

Gaussian States in LQC

Gaussian states centered around $|V| = \mu l_P$ are states of the following form:⁶

$$|\psi\rangle = \frac{c}{\sqrt{2}} \sum_{n \in \mathbb{Z}} \left(e^{-(nl_P - \mu l_P)^2 / 2\sigma^2} + e^{-(nl_P + \mu l_P)^2 / 2\sigma^2} \right) |nl_P\rangle$$

where l_P is the Planck length; henceforth $l_P = 1$.



⁶A generalization of Willis, PhD Thesis, 2004.

Fluctuations for Gaussian States

By construction, Gaussian states in LQC are “sharply-peaked” around $|V| = \mu$. Thus, in order for them to be semi-classical kinematic states, we only require that their fluctuations in \hat{V} and \hat{S}_λ are minimal.

Fluctuations for Gaussian States

The exact volume-fluctuations are thus:

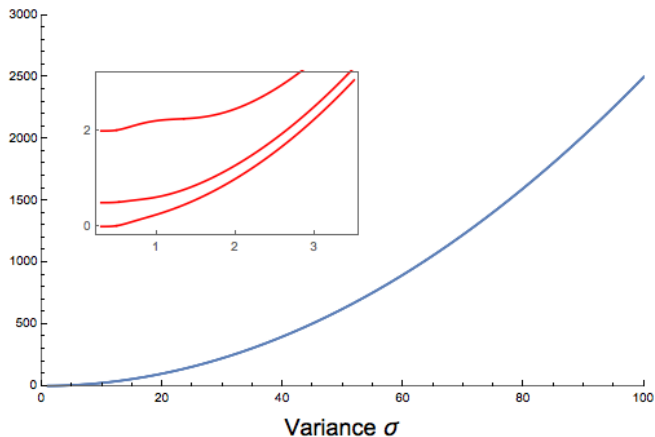
$$\begin{aligned} \Delta_V^2 = & \frac{\sqrt{-\ln v}}{4\sqrt{\pi} (1 + e^{\mu^2\pi^2/\ln v}) \vartheta_3(v^{1/\pi^2\sigma^4})} \\ & \times \left\{ \left[2\mu^2 - \frac{\ln v}{\pi^2} (1 + e^{\mu^2\pi^2/\ln v}) \right] \vartheta_3(v) \right. \\ & \left. - (1 + e^{\mu^2\pi^2/\ln v}) \left(\frac{2v(\ln v)^2}{\pi^2} \right) \left(\frac{d}{dv} \vartheta_3(v) \right) \right\} \end{aligned}$$

where $v := \exp\{-\pi^2\sigma^2\}$ and $\vartheta_3(q) = \vartheta_3(0, q)$ is the Jacobi ϑ -function.

Fluctuations for Gaussian States

$$\vartheta_3(z, q) = \sum_{n \in \mathbb{Z}} q^{n^2} e^{2niz}$$

Fluctuations for Gaussian States



Δ_V^2 as a function of σ for $\mu \in \{0, 1, 2\}$ with $\frac{\alpha\lambda}{V_0} = 1$.

Fluctuations for Gaussian States

The exact holonomy-flux fluctuations have two cases to consider. If $2\alpha\lambda/V_0 \notin \mathbb{Z}$, then:

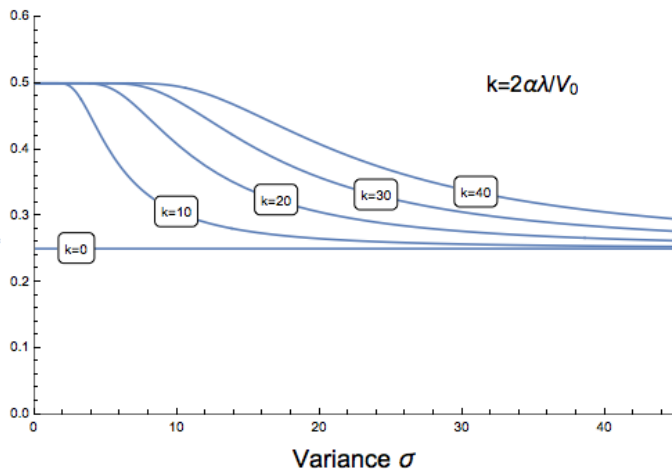
$$\Delta_{S_\lambda}^2 = \frac{1}{2\lambda^2}$$

Fluctuations for Gaussian States

If $k = 2\alpha\lambda/V_0$ is an integer, then:

$$\begin{aligned} \Delta_{S_\lambda}^2 = & \frac{\alpha^2 4\pi\sigma}{k^2 V_0^2} - \frac{\alpha^2 \pi\sigma}{k^2 V_0^2 (1 + e^{-\mu^2/\sigma^2}) \vartheta_3(e^{-1/\sigma^2})} \times \\ & \left\{ 2e^{-2(\alpha\lambda/V_0)^2/\sigma^2} \left[\vartheta_3\left(\frac{-\alpha\lambda/V_0}{i\sigma^2}, e^{-1/\sigma^2}\right) + \vartheta_3\left(\frac{\alpha\lambda/V_0}{i\sigma^2}, e^{-1/\sigma^2}\right) \right] \right. \\ & + e^{-2(\alpha\lambda/V_0 - \mu)^2/\sigma^2} \left[\vartheta_3\left(\frac{-\alpha\lambda/V_0 + \mu}{i\sigma^2}, e^{-1/\sigma^2}\right) + \vartheta_3\left(\frac{\alpha\lambda/V_0 - \mu}{i\sigma^2}, e^{-1/\sigma^2}\right) \right] \\ & \left. + e^{-2(\alpha\lambda/V_0 + \mu)^2/\sigma^2} \left[\vartheta_3\left(\frac{-\alpha\lambda/V_0 - \mu}{i\sigma^2}, e^{-1/\sigma^2}\right) + \vartheta_3\left(\frac{\alpha\lambda/V_0 + \mu}{i\sigma^2}, e^{-1/\sigma^2}\right) \right] \right\} \end{aligned}$$

Fluctuations for Gaussian States

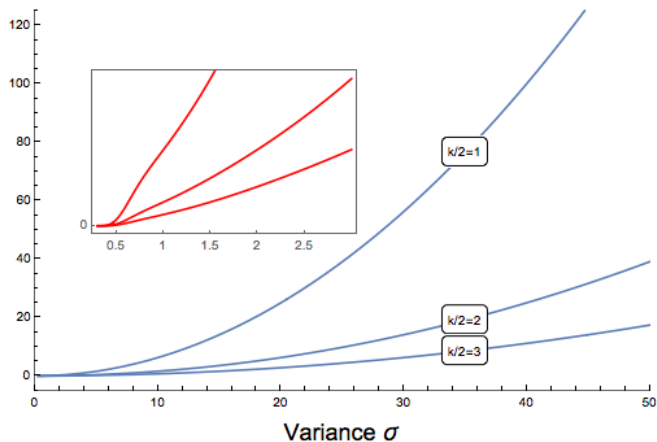


$\Delta_{S_\lambda}^2$ as a function of σ for different λ with $\mu = 0$.

Fluctuations for Gaussian States

The question now is whether or not $\Delta_V^2 \Delta_{S_\lambda}^2$ is equal to the RS lower-bound. There are three different cases that lead to different behaviours, namely, if $\alpha\lambda/V_0 \notin \mathbb{Z}$, or $\alpha\lambda/V_0 \in \mathbb{Z}$ and is even or odd. In all three cases, the fluctuations *never* saturate the inequality (for any values of σ , μ , and λ).

Fluctuations for Gaussian States



$\Delta_V^2 \Delta_{S_\lambda}^2 - L.B.$ as a function of σ for different k with $\mu = 0$ and $\alpha\lambda/V_0 \in \mathbb{Z}$ even.

Conclusion

Unlike in traditional quantum mechanics, Gaussian states in LQC do *not* minimize the uncertainty relation. Thus, they are not strictly semi-classical!

Follow-Up

Are the conditions under which they are *essentially* semi-classical?

Fluctuations for Gaussian States

Yes! One can check that in the $V_0 \rightarrow \infty$ limit, the fluctuations (and lower bound) vanish.⁷ However, V_0 must be fixed within a model of LQC; it is not a parameter that can be varied.

⁷Rovelli & Wilson-Ewing, Phys. Rev. D, 2014.

Summary

- 1 For a Gaussian state with a fixed width σ and a fixed parity-symmetric mean μ , one can always choose a model of LQC with a sufficiently large fiducial volume V_0 such that the fluctuations of this state become negligible.
- 2 Within a fixed model of LQC (with a given finite V_0), one can always find a Gaussian state with sufficiently large σ such that its fluctuations become arbitrarily large.
- 3 The relation between λ and V_0 play an important role in determining the phenomenology of a theory of LQC, namely, they determine whether or not the holonomy operator \hat{h}_λ is closed, and the different integer cases affect fluctuations.

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