Probing effective quantum gravity with fluid dynamics

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- 1. Introduction
- 2. The kinetic gas model
- 3. Gravitational field of a kinetic gas
- 4. Kinetic gas in modified Schwarzschild spacetime
- 5. Conclusion

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 - Replace metric spacetime geometry by Finsler geometry.
 - Similarly: replacing flat spacetime by curved spacetime led to GR.
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 - Similarly: replacing flat spacetime by curved spacetime led to GR.
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- Questions arising from new matter model:
 - ✓ How does a kinetic gas react to a gravitational field?
 - ? How does a kinetic gas create a gravitational field?

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\Rightarrow Here: effective quantum gravity phenomenology with gas dynamics near black holes.

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→ Only need to study (all) possible quantum corrections!

- Perfect fluid:
 - Most general energy-momentum tensor compatible with cosmological symmetry.
 - No shear stress, no friction.
 - Characterized by density ρ and pressure p.
 - Dust, dark matter: p = 0.
 - Radiation: $p = \frac{1}{3}\rho$.
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- Hyperfluid:
 - Additional coupling to affine connection generates hypermomentum.
 - Intrinsic property of matter, e.g., spin.

Why study matter beyond fluids?

- Dynamical friction:
 - Massive object passes distribution of light objects.
 - \Rightarrow Gravity of massive object changes positions of lighter objects.
 - ⇒ Perturbation of light objects asserts gravity on massive object.
 - Example: globular cluster passing through galaxy.

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- Dynamics of intergalactic medium:
 - Cosmic gas highways: gas in and near filaments
 - Crossing sheets in collapse and structure formation.

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Definition of kinetic gas

- Single-component gas:
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 - Particles follow piecewise geodesic curves.
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- Multi-component gas: multiple types of particles.

- Kinetic gas described by density in velocity space:
 - Consider space *O* of physical (unit, timelike, future pointing) four-velocities.
 - Consider density on physical velocity space.

One-particle distribution function

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For every hypersurface $\sigma \subset O$,

$$N[\sigma] = \int_{\sigma} \phi \Omega$$

of particle trajectories through σ .

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• For multi-component fluids: ϕ_i for each component *i*.

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of outbound trajectories - # of inbound trajectories.

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- \Rightarrow Collision density measured by $\mathcal{L}_{\mathbf{r}}\phi$.
- Collisionless fluid: trajectories have no endpoints, $\mathcal{L}_{\mathbf{r}}\phi = \mathbf{0}$.
- \Rightarrow Simple, first order equation of motion for collisionless fluid.
- $\Rightarrow \phi$ is constant along integral curves of **r**.

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Quantum gravity & fluid dynamics

Geodesic dust fluid:

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Geodesic dust fluid: Collisionless fluid: Interacting fluid: $\phi(\mathbf{x},\mathbf{y})\sim\delta(\mathbf{y}-\mathbf{u}(\mathbf{x}))\,.$ $\mathcal{L}_{\mathbf{r}}\phi \neq \mathbf{0}$. $\mathcal{L}_{\mathbf{r}}\phi = \mathbf{0}$. "Humppa"

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$$0 = \nabla u^{a} = u^{b} \partial_{b} u^{a} + u^{b} N^{a}{}_{b},$$
$$0 = \nabla_{\delta_{a}}(\rho u^{a}) = \partial_{a}(\rho u^{a}) + \frac{1}{2} \rho u^{a} g^{F bc} \left(\partial_{a} g^{F}_{bc} - N^{d}{}_{a} \bar{\partial}_{d} g^{F}_{bc} \right).$$

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- ⇒ Generalized (pressureless) Euler equations to Finsler geometry [MH 15].
- Metric limit $F^2(x, y) = |g_{ab}(x)y^ay^b|$ yields Euler equations:

$$u^b \nabla_b u^a = 0$$
, $\nabla_a (\rho u^a) = 0$.

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Action for a single point particle:

$$S=m\int_0^t (F\circ c_1)(au)\,d au$$
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Assume arc length parameter τ :

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$$S = mt$$
.

Quantum gravity & fluid dynamics

 $c_1(t)$

 $c_1(0)$

Action for *P* point particles:

$$S_{\text{gas}} = m \sum_{i=1}^{P} \int_{0}^{t} (F \circ c_i)(\tau) \, d au$$
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Assume arc length parameter τ :

$$S_{\text{gas}} = Pmt$$



• Hypersurface of starting points:

 $c_i(0) \in \sigma_0$.





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• Hypersurface of end points:

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Consider volume

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• Recall particle action integral:

$$S_{gas} = Pmt = m \int_{0}^{t} \left(\int_{\sigma_{\tau}} \phi \Omega \right) d\tau$$

$$= m \int_{V} \phi \Omega \wedge \omega$$

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Defined through 1-PDF ϕ

[MH, Pfeifer, Voicu '19].

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$$\sigma_t$$

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⇒ Forget particle trajectories!

$$S_{
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• Finsler Ricci scalar $R_0 = L^{-1} R^a{}_{ab} y^b$ from curvature of non-linear connection:

$$R^{a}{}_{bc}\bar{\partial}_{a} = (\delta_{b}N^{a}{}_{c} - \delta_{c}N^{a}{}_{b})\bar{\partial}_{a} = [\delta_{b}, \delta_{c}].$$

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! Unique action obtained from variational completion of Rutz equation [MH, Pfeifer, Voicu '18]. ⇒ Reduces to Einstein-Hilbert action for metric geometry.
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• Variation of the Finsler gravity action:

$$\delta_{F}S_{\text{grav}} = 2\int_{V} \left[\frac{1}{2}g^{Fab}\bar{\partial}_{a}\bar{\partial}_{b}(F^{2}R_{0}) - 3R_{0} - g^{Fab}(\nabla_{\delta_{a}}P_{b} - P_{a}P_{b} + \bar{\partial}_{a}(\nabla P_{b}))\right]\frac{\delta F}{F}\Sigma$$

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Landsberg tensor measures deviation from metric geometry:

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⇒ Gravitational field equations with kinetic gas matter [MH, Pfeifer, Voicu '19]:

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Physical implications

- There are no metric non-vacuum solutions to the field equations.
 - Field equations in case of a metric geometry $F^2 = g_{ab}(x)y^ay^b$:

$$3r_{ab}(x)y^ay^b - r(x)g_{ab}(x)y^ay^b = -\kappa^2\phi g_{ab}(x)y^ay^b$$
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Second derivative with respect to velocities y^a and y^b:

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- \Rightarrow 1-PDF ϕ must depend only on *x*, i.e., independent of velocities *y*.
- Juphysical velocity distribution: uniform over all (arbitrarily high) velocities!
- \Rightarrow Gravitational field of a kinetic gas always depends on the velocity of the observer.
 - For observers whose velocity exceeds that of any gas particles:

$$\frac{1}{2}g^{Fab}\bar{\partial}_{a}\bar{\partial}_{b}(F^{2}R_{0}) - 3R_{0} - g^{Fab}(\nabla_{\delta_{a}}P_{b} - P_{a}P_{b} + \bar{\partial}_{a}(\nabla P_{b})) \rightarrow 0$$

• Solution of the differential equation still depends on ϕ via boundary conditions. \Rightarrow Observers at velocities beyond gas velocities are still affected, but differently.

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Example: collisionless dust fluid φ(x, y) ~ ρ(x)δ_{Sx}(y, u(x)):

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• Next task: solve cosmological field equations with kinetic gas.

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- 2. The kinetic gas model
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 - Spherically symmetric spacetime.
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- *κ*-Poincaré modification of spacetime:
 - Interaction between particles and "quantum structure of spacetime".
 - Interaction depends on de Broglie wavelength (momentum).
 - → Distinguished time direction (vector field).
 - $\Rightarrow \kappa$ -Minkowski spacetime has modified symmetry algebra.
 - Black hole spacetime: assume spherically symmetric vector field.
 - \Rightarrow Vector field may only have time and radial components.
 - $\circ\,$ Modification depends on a parameter ℓ (Planck length).
 - $\circ~$ Spacetime approaches Schwarzschild for $\ell \rightarrow 0.$

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- \rightsquigarrow Consider gas $\phi \sim \delta(E)\delta(L)\delta(H)$ of identical energy, angular momentum, mass.

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Example: radial free fall from infinity

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 - Model many-particle systems defined by individual point mass trajectories.
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- More general applications of kinetic gas model:
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 - Weak field limit: Newtonian, post-Newtonian...
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 - Stellar streams?
 - Dynamics of heterogeneous systems: stars + gas in galaxies?