

### **Towards future multi-messenger detections of** core-collapse supernovae harbouring choked jets

Based on A. Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

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RUB



**RUHR-UNIVERSITÄT** BOCHUM

17th Marcel Grossmann Meeting 7-12 July 2024, Pescara (Italy)





### Outline

CHOKED JETS FROM CORE-COLLAPSE SUPERNOVAE

- Modelling and possible progenitors
- Why choked jets are important in the multi-messenger field
  - Ultraviolet (UV) and optical electromagnetic signatures from choked jets
  - Choked jets as potential sources of high-energy neutrinos

### **FUTURE PROSPECTS FOR MULTI-MESSENGER DETECTIONS**

- Description of methodology adopted
- Detection prospects of choked jet progenitors in UV band by the future UV mission ULTRASAT in relation to their visibility in optical (ZTF-like instruments)
- How neutrino observations by Cherenkov-based high-energy neutrino telescopes (e.g., IceCube, KM3NeT) can be used in association to UV and optical signals
- Results and discussion of multi-messenger implications



# **Core-collapse Supernovae (CCSNe)**



### **SNe classification**







# Choked jets in CCSNe





### A choked jet with a cocoon breakout







### **Observable electromagnetic signatures of GRB cocoons**

Even if the **jet is choked**, and  $\gamma$  **rays** emitted by the relativistic jet cannot escape, other observable electromagnetic (EM) signatures are:

- 1. Shock Breakout (SBO): bright X-ray or  $\gamma$ -ray flash (it lasts from a few seconds up to fractions of hour)
- 2. <u>Cooling phase of the expanding cocoon envelope</u>: UV and optical signal (timescale of days)



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### Choked jet in the multi-messenger field





### Choked jet in the multi-messenger field





### Choked jet in the multi-messenger field





### Aims of our work

**RSG** and **BSG** stars as progenitors of Type II SNe harbouring **choked relativistic jets** in their stellar envelopes



### **Ultraviolet Transient Astronomy Satellite**

An eye in the UV sky with the UItraviolet Transient Astronomy Satellite (ULTRASAT) Expected to be <u>launched in 2026</u>



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ULTRASAT Key Properties				
Property	Value	Comments		
Mission lifetime	3 years	Propellent for 6 years		
Drbit	GEO			
Fotal FoV	204 deg <sup>2</sup>	Unprecedented FoV		
Operation waveband	230-290 nm	Near-UV band		
Cadence	300 s	For the high cadence survey		
Mean effective PSF	8.3"	In central 170 deg <sup>2</sup> For T = 20,000 K blackbody source		
Mean limiting magnitude (in 900s, 5 <b>o</b> )	22.4 ABmag	In central 170 deg <sup>2</sup> For T = 20,000 K blackbody source		

Shvartzvald et al. 2023





### **Zwicky Transient Facility**

<u>Bellm et al. 2019</u>

### Palomar Observatory in Southern California, USA



Image credit: Palomar Observatory/Caltech

Specifications of the ZTF Observing System

### ZTF has been surveying the **Northern sky** since June 2018 every 2-3 nights in gband (370-560) nm and *r*-band (550-740) nm filters, while the *i*-band filter at (690-895) nm is used for partnership observations

Telescope	Palo
Location	33° 2
Camera field dimensions	7°.50
Camera field of view	55.0
Light-sensitive area	47.7
Fill factor	86.7
Filters	ZTF
Filter exchange time	$\sim 110$
Image quality	<i>g</i> =
Median Sensitivity	$m_g =$
$(30 \text{ s}, 5\sigma)$	_
	$m_g =$







### **Zwicky Transient Facility**

Palomar Observatory in Southern California, USA

### The Zwicky Tr been surveying **June 2018** eve

(370-560) nm and 7-band (550-740) nm nm is used for partnership observations

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Median Sensitivity	$m_g =$
$(30 \text{ s}, 5\sigma)$	
	$m_g =$





### Early UV/optical emission from choked jets embedded in CCSNe

Model from *Waxman et al., ApJ 667, 351 (2007), Rabinak & Waxman, ApJ 728, 63 (2011)* 

Photosphere  

$$r_{\rm ph}(t) = \begin{cases} 3.3 \times 10^{14} f_{\rho}^{-0.062} \frac{E_{51}^{0.41} k_{0.34}^{0.093}}{(M_{\rm ej}/M_{\odot})^{0.31}} t_{5}^{0.45} \\ 3.3 \times 10^{14} f_{\rho}^{-0.036} \frac{E_{51}^{0.39} k_{0.34}^{0.11}}{(M_{\rm ej}/M_{\odot})^{0.28}} t_{5}^{0.45} \\ T_{\rm ph}(t) = \begin{cases} 1.6 f_{\rho}^{-0.037} \frac{E_{51}^{0.027} R_{*,13}^{1/4}}{(M_{\rm ej}/M_{\odot})^{0.054} k_{0.34}^{0.28}} t_{5}^{-0.45} \\ 1.6 f_{\rho}^{-0.022} \frac{E_{51}^{0.016} R_{*,13}^{1/4}}{(M_{\rm ej}/M_{\odot})^{0.033} k_{0.34}^{0.27}} t_{5}^{-0.47} \end{cases}$$

 $R_* =$  progenitor radius  $f_{\rho}$  = factor related to the average ejecta density (it varies linearly with mass of progenitors M) k = opacity of the stellar envelope in cm<sup>2</sup> g<sup>-1</sup>  $E = 10^{51}E_{51}$  erg = released energy  $t = 10^5t_5$  s = time from the shock breakout n = 3/2 for convective envelopes (RSGs), and n = 3 for radiative envelopes (BSGs)

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 $^{.81}$  cm (n = 3/2) $^{.78}$  cm (n = 3)eV (n = 3/2)eV (n = 3)

# Model specific intensity observed in UV/optical $f_{\lambda}(\lambda, t) = \left(\frac{r_{\text{ph}}}{D(\tau)}\right)^2 \sigma T_{\text{ph}}^4 \frac{T_{\text{col}}}{L_{\lambda}} g_{BB}(x) e^{-\tau_{\lambda}}$

$$(\lambda, t) = \left(\frac{1}{D_L(z)}\right) \sigma T_{\text{ph}}^4 \frac{1}{hc} \cos^2 g_{BB}(x) e^{-\tau_{\lambda}}$$

$$g_{BB}(x) = \frac{15}{\pi^4} \frac{x^5}{e^x - 1}$$

$$x = \frac{hc}{\lambda} T_{\text{col}}$$

 $T_{\rm col} \simeq 1.2 \ T_{\rm ph}$  = color temperature, i.e., temperature at which a black body would emit radiation of the same color of a given source  $\tau_{\lambda}$  = extinction optical depth at a given wavelength  $\lambda$ 





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### Early UV/optical emission from choked jets embedded in CCSNe

### Model specific intensity observed in UV/optical



Black-Body radiation modified by extinction

$$f_{\lambda}(\lambda, t) = \left(\frac{r_{\rm ph}}{D_L(z)}\right)^2 \sigma T_{\rm ph}^4 \frac{T_{\rm col}}{hc} g_{BB}(x) e^{-\tau_{\lambda}}$$

$$\int g_{BB}(x) = \frac{15}{\pi^4} \frac{x^5}{e^x - 1}$$

$$x = hc/\lambda T_{\rm col}$$

 $T_{\rm col} \simeq 1.2 \ T_{\rm ph}$  = color temperature, i.e., temperature at which a black body would emit radiation of the same color of a given source

 $\tau_{\lambda}$  = extinction optical depth at a given wavelength  $\lambda$ 

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![](_page_19_Picture_5.jpeg)

![](_page_20_Figure_0.jpeg)

![](_page_20_Picture_5.jpeg)

![](_page_21_Figure_0.jpeg)

![](_page_21_Picture_5.jpeg)

![](_page_22_Figure_0.jpeg)

ZTF Bright Transient Survey catalogue available <u>online</u>

Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

![](_page_22_Picture_5.jpeg)

![](_page_23_Figure_0.jpeg)

![](_page_23_Picture_5.jpeg)

![](_page_24_Figure_0.jpeg)

![](_page_24_Picture_5.jpeg)

### Maximum detectable redshift as a function of time after SBO

By investigating the evolution of the cooling UV/optical signal after SBO for different redshift values, the following is obtained

![](_page_25_Figure_2.jpeg)

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	$z_{ m lim}$	$D_L$ [Mpc]	$t - t_{\rm SBO}$ [days
ULTRASAT	0.08	360	4
ZTF	0.06	270	10
	z <sub>lim</sub>	$D_L$ [Mpc]	$t - t_{\rm SBO}$ [days
ULTRASAT	0.023	100	1
ZTF	0.017	75	4

**Previous and later emissions can still exceed the** sensitivity of the detectors only for closer SNe

![](_page_25_Figure_7.jpeg)

$$\forall t \to \dot{N} = \int d\Omega \int_0^{z_{\rm lim}} \frac{dN(z)}{dz} dz = \int d\Omega \int_0^{z_{\rm lim}} \frac{R(z)}{1+1} dz$$

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 $\frac{dV(z)}{dz} \frac{dV(z)}{dz} dz$ 

![](_page_26_Picture_6.jpeg)

$$\forall t \to \dot{N} = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{dN(z)}{dz} dz = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{R(z)}{1+z_{\text{lim}}} dz$$
Telescope FoV

Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

 $\frac{dz}{dz} \frac{dV(z)}{dz} dz$ 

![](_page_27_Picture_5.jpeg)

$$\forall t \to \dot{N} = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{dN(z)}{dz} dz = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{R(z)}{1+1} dz$$
Telescope FoV

Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

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 $(z) \frac{dV(z)}{dz}$ +z dz

Differential comoving volume

![](_page_28_Picture_6.jpeg)

$$\forall t \to \dot{N} = \int d\Omega \int_0^{z_{\rm lim}} \frac{dN(z)}{dz} dz = \int d\Omega \int_0^{z_{\rm lim}} \frac{R(z)}{1+z_{\rm lim}} dz$$

**Telescope FoV** 

This factor accounts for Differential the cosmological time comoving volume dilation of the observed rate

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 $(z) \frac{dV(z)}{dz}$ dz

![](_page_29_Picture_8.jpeg)

$$\forall t \to \dot{N} = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{dN(z)}{dz} dz = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{R(z)}{1-z} dz$$

**Telescope FoV** 

This factor accounts for Differential comoving volume the cosmological time dilation of the observed rate

Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

![](_page_30_Picture_6.jpeg)

![](_page_30_Figure_7.jpeg)

![](_page_30_Picture_8.jpeg)

$$\forall t \to \dot{N} = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{dN(z)}{dz} dz = \int d\Omega \int_{0}^{z_{\text{lim}}} \frac{R(z)}{1-z} dz$$

**Telescope FoV** 

This factor accounts for the cosmological time dilation of the observed rate

Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

![](_page_31_Figure_7.jpeg)

![](_page_31_Picture_8.jpeg)

![](_page_32_Figure_0.jpeg)

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![](_page_32_Picture_5.jpeg)

![](_page_32_Picture_6.jpeg)

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![](_page_33_Figure_0.jpeg)

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![](_page_33_Picture_3.jpeg)

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![](_page_34_Figure_2.jpeg)

![](_page_34_Picture_6.jpeg)

## High-energy $\nu$ - UV - optical follow-ups

1. Neutrino alert from Cherenkovbased high-energy neutrino telescopes (e.g., IceCube, KM3NeT) 2. ULTRASAT(ZTF-like telescopes) could point at the suggested direction in the sky within ~4(10) days to search for possible EM counterparts maximising the reachable sky volume and hence the number of detectable sources

![](_page_35_Figure_3.jpeg)

![](_page_35_Picture_6.jpeg)

# High-energy $\nu$ - UV - optical follow-ups

1. Neutrino alert from Cherenkovbased high-energy neutrino telescopes (e.g., IceCube, KM3NeT)

2. ULTRASAT(ZTF-like telescopes) could To consider the **production of neutrinos** point at the suggested direction in the sky during the shock propagation inside the within ~5(11) days to search for possible stellar envelope, we need to enlarge EM counterparts maximising the reachable these time windows up to ~1 day sky volume and hence the number of  $10^{4}$ detectable sources Kister et al. 2013 **Red Supergiants**  $\begin{bmatrix} 10^3 \\ 8ec \end{bmatrix}$  SBO duration [sec] 10<sup>2</sup> 10<sup>2</sup> 10<sup>2</sup> Optical **Blue Supergiants 5** × 10<sup>50</sup> erg Wolf–Rayet Stars •  $3 \times 10^{51}$  erg  $10^{3}$  $10^{4}$ 10<sup>5</sup>  $10^{2}$ 10 5 days 11 days Shock propagation time [sec] Time 1 hour 1 day Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

![](_page_36_Figure_3.jpeg)

![](_page_36_Figure_5.jpeg)

![](_page_36_Picture_6.jpeg)

# High-energy $\nu$ - UV - optical follow-ups

1. **Neutrino alert** from Cherenkovbased high-energy neutrino telescopes (e.g., IceCube, KM3NeT) 2. ULTRASAT(ZTF-like telescopes) could point at the suggested direction in the sky within ~5(11) days to search for possible EM counterparts maximising the reachable sky volume and hence the number of detectable sources

![](_page_37_Figure_3.jpeg)

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With our results, we would like to stress the **importance of defining a proper strategy**, **focussed on Type II SNe** 

We encourage UV, optical, and neutrino telescopes to optimise both their alert sending and external follow-up programmes based on our results

Type IIP SNe (associated to RSGs) might

<u>contribute</u> to the production of high-energy neutrinos for ~60% of astrophysical diffuse flux between 10<sup>3</sup> and 10<sup>5</sup> GeV

11 days

Optical

Time

Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

![](_page_37_Picture_12.jpeg)

### How many choked jets producing neutrinos?

M. Fasano, S. Celli, D. Guetta, A. Capone, A. Zegarelli, I. Di Palma, JCAPO9 (2021) 044 M. M. Fasano, S. Celli, D. Guetta, A. Capone, A. Zegarelli, I. Di Palma, JCAPO9 (2021) 044

 $p\gamma$  interactions simulated inside a choked GRB jet through a detailed Monte Carlo code

![](_page_38_Figure_3.jpeg)

![](_page_38_Figure_6.jpeg)

![](_page_38_Picture_7.jpeg)

### How many choked jets producing neutrinos?

M. Fasano, S. Celli, D. Guetta, A. Capone, A. Zegarelli, I. Di Palma, JCAPO9 (2021) 044 M. M. Fasano, S. Celli, D. Guetta, A. Capone, A. Zegarelli, I. Di Palma, JCAPO9 (2021) 044

 $p\gamma$  interactions simulated inside a choked GRB jet through a detailed Monte Carlo code

![](_page_39_Figure_3.jpeg)

![](_page_39_Picture_8.jpeg)

### Multi-messenger prospects

1. Neutrino alert from Cherenkovbased high-energy neutrino telescopes (e.g., IceCube, KM3NeT)

2. ULTRASAT(ZTF-like telescopes) could point at the suggested direction in the sky within ~5(11) days to search for possible EM counterparts maximising the reachable sky volume and hence the number of

![](_page_40_Figure_3.jpeg)

### Take-home messages

Choked jets (failed GRBs) as appealing sources in the multi-messenger astronomy field -> possible contributors to the astrophysical diffuse neutrino flux

- of the cocoon emission with time and the extinction dependence with wavelength
- ~4(5) days, with a subsequent follow-up by instruments like ZTF about one week after
- - **TeV neutrinos** —> <u>need to run these multi-messenger analyses for several years</u>
- harbouring choked jets as contributors to the diffuse astrophysical high-energy neutrino flux

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We propose an optimised follow-up strategy between UV, ZTF-like, and neutrino telescopes, considering RSG and BSG as progenitors of CCSNe possibly harbouring choked jets, taking into account the evolution

The delay between neutrino produced at the SBO occurrence (during the jet propagation inside the stellar envelope) and ULTRASAT (future UV satellite with unprecedented FoV) observations should be of

Less than 1% of CCSNe from RSGs detectable in UV with ULTRASAT could host a choked jet and release

**EM** and  $\nu$  detections, if accompanied by additional photometric and spectroscopic follow-ups with compelling evidence for a relativistic jet launched by the source central engine, would suggest CCSNe

![](_page_41_Picture_13.jpeg)

# Thank you for your attention!

![](_page_43_Picture_1.jpeg)

### Effect of dust extinction

$$f_{\rm obs}(\lambda) = f_{\lambda}(\lambda, t) = \left(\frac{r_{\rm ph}}{D_L(z)}\right)^2 \sigma T_{\rm ph}^4 \frac{T_{\rm col}}{hc} g_{BB}(x) e^{-\tau_{\lambda}} = f_{\rm int}(\lambda)$$

$$f_{obs}(\lambda) = f_{int}(\lambda)e^{-\tau_{\lambda}} = f_{int}(\lambda)10^{-0.4A_{\lambda}}$$
Observed color excess
$$E_{B-V} = E_{B-V}^{observed} - E_{B-V}^{intrinsic}$$

$$A_{\lambda} = 1.086\tau_{\lambda} = k(\lambda)E_{B-V} = k(\lambda)\frac{A_{V}}{R_{V}}$$
Reddening curve
We adopt the averaged dust extinction model from
$$Cardelli \ et \ al. \ (1989)$$
for diffuse interstellar medium in the Milky Way with  $R_{V} = 3$ .

Zegarelli et al., submitted to A&A, arXiv:2403.16234 [astro.ph.HE]

![](_page_44_Picture_5.jpeg)

![](_page_44_Figure_6.jpeg)

### Signatures of SBO in SN light curves

First entire observations of the shock breakouts of Type II Plateau SNe were reported in 2008 by ultraviolet and optical observations by the GALEX satellite and supernova legacy survey (SNLS), named SNLS-04D2dc and SNLS-06D1jd

![](_page_45_Figure_2.jpeg)

![](_page_45_Figure_4.jpeg)

![](_page_45_Picture_5.jpeg)

![](_page_46_Figure_0.jpeg)

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The grasp is given for a 20,000 K black-body source spectrum (e.g., a hot transient).

ULTRASAT's grasp is an order of magnitude larger than that of current surveys, two orders of magnitudes larger than that of GALEX, the largest grasp UV mission to date, and com- parable to that of LSST, the largest grasp optical survey under construction.

![](_page_46_Picture_5.jpeg)

### Limiting magnitude of ULTRASAT

![](_page_47_Figure_1.jpeg)

![](_page_47_Picture_3.jpeg)

![](_page_48_Figure_0.jpeg)

### Effect of extinction

![](_page_49_Figure_1.jpeg)

![](_page_49_Picture_3.jpeg)

### **Previous predictions for ULTRASAT from literature**

<u>Ganot et al., ApJ 820, 57 (2016)</u>

![](_page_50_Figure_2.jpeg)

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Data points from <u>SNLS-04D2dc Type II detected by GALEX</u> (extinction  $A_{\text{NUV}} = 1.45$ )

$$L_{\rm NUV}(t) = 4\pi D_L^2(z) \int_{\lambda_{min}}^{\lambda_{max}} f_{\lambda}(\lambda, t) d\lambda$$

assume the reasonable parameters for RSG and BSG stars, i.e.,  $R = 3.5 \times 10^{13} \text{ cm} (= 500 \text{ R}_{Sun}), E = 2 \times 10^{51} \text{ erg}, M = 10 \text{M}_{Sun}$  $R = 3.5 \times 10^{12} \text{ cm} (= 50 \text{ R}_{Sun}), E = 1 \times 10^{51} \text{ erg}, M = 10 M_{Sun}$ and for WR, values from SN2008D

 $R = 10 \times 10^{11}$  cm,  $E = 6 \times 10^{51}$  erg,  $M = 7M_{Sun}$ (extinction  $A_{\text{NUV}} = 2.2$ )

![](_page_50_Picture_9.jpeg)

![](_page_50_Figure_12.jpeg)

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### Previous predictions for ULTRASAT from literature

<u>Ganot et al., ApJ 820, 57 (2016)</u>

Table 2 Predicted SN Explosion Detection Numbers by Various Surveys

Survey I	Band	Cadence	FOV (deg <sup>2</sup> )	Expected Number (SN yr <sup>-1</sup> )					
				RSG		BSG		W–R	
				<1 hr	<1 day	<1 hr	<1 day	<1 hr	<1 day
GALEX/PTF	NUV	3 day	600	0	30 <sup>a</sup>	0	0	0	0
ULTRASAT	NUV	900 s	210	4	85	1	8	0	3
ULTRASAT	NUV	3600 s	210	16	314	2	31	1	14
iPTF <sup>b</sup>	r	1 day	1000	0	7	0	2	0	1
ZTF <sup>c</sup>	g	0.5 hr	2100	0	10	0	2	0	1
LSST	g	0.5 hr	9.6	0	17	0	3	0	1

### Notes.

<sup>a</sup> For our GALEX/PTF experiment, we report the expected number within three days (not one day) to match its low actual cadence. As the survey ran for 2 m (1/6 yr), the expected number of SNe from RSG explosions for the actual experiment is 30/6 = 5 events. Assumed temporal efficiency of 25% (including loss due to daytime and average weather) and lunation-averaged depth of 20.6 mag. 25% temporal efficiency as above, average depth 20.4 mag, and 50% survey time spent in g-band. Assumed the following for the LSST deep drilling project: 1 LSST field observed at any given time, 25% temporal efficiency as above, and g = 24.2 mag lunationaveraged depth.

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These estimations were performed considering fiducial parameters for each type of source and <u>without extinction</u>

![](_page_51_Figure_9.jpeg)

![](_page_51_Picture_10.jpeg)