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## Distinguishing Signature of Kerr-MOG Black Hole and Naked Singularity via Lense-Thirring Effect

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We examine the geometrical difference between non-extremal black hole(NXBH), extremal black hole(XBH) and naked singularity(NS) via Lense-Thirring(LT) effect in spinning modified-gravity(MOG). For NXBH, we find that the LT frequency ( $\Omega_{LT}$ ) is proportional to the angular-momentum( $J = a \text{ cal}M$ ) parameter or spin parameter( $a$ ) i.e.  $\Omega_{LT} \propto J$  or  $\Omega_{LT} \propto a$ [where  $\text{cal}M = M(1 + \alpha)$  is ADM mass,  $\alpha$  is MOG parameter and  $M$  is Komar mass] and is inversely proportional to the cubic value of radial parameter i.e.  $\Omega_{LT} \propto \frac{1}{r^3}$ .

For XBH ( $a^2 = \frac{G_N^2 \text{cal}M^2}{1+\alpha}$ ), we find LT frequency is proportional to the angular-momentum parameter i.e.  $\Omega_{LT} \propto \frac{1}{\sqrt{1+\alpha}}$  and is inversely proportional to the cubic value of radial parameter i.e.  $\Omega_{LT} \propto \frac{1}{r^3}$ . While for NS, we find  $\Omega_{LT} \propto \frac{\text{cal}M^3}{J^3}$  and  $\Omega_{LT} \propto \frac{\left[ r - \left( \frac{\alpha}{1+\alpha} \right) \frac{G_N \text{cal}M}{2} \right]}{\sqrt{1 + \left( \frac{\alpha}{1+\alpha} \right) \frac{G_N^2 \text{cal}M^2}{a^2}}}$  in the limit  $\theta = 0$  and  $a = \frac{J}{\text{cal}M} \gg r$ . It depends both on the angular momentum parameter and MOG parameter  $\sim \alpha$ .

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